

RECURRENT GEODESICS IN FLAT LORENTZ 3-MANIFOLDS

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ABSTRACT. We introduce the notion of recurrent geodesic rays in a complete flat Lorentz 3-manifold. We completely classify the dynamical behavior of geodesics in cyclic quotients, and apply this classification to more general quotients. In particular we show that for any pair of closed geodesics γ_1, γ_2 there exists a unique geodesic which spirals around γ_1 in forward time and γ_2 in backwards time.

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1. INTRODUCTION

In this paper, we investigate dynamical properties of geodesics in flat Lorentz 3-manifolds M . We assume M is geodesically complete, that is, M is the quotient $\mathbb{A}^{2,1}/\Gamma$ of 3-dimensional Minkowski spacetime $\mathbb{A}^{2,1}$ by a discrete group Γ of affine isometries acting properly on $\mathbb{A}^{2,1}$.

The first section develops preliminaries on Minkowski space and its isometries. Section 3 is devoted to the particular case of *cylinders*: quotients of spacetime by a cyclic hyperbolic group $\langle \gamma \rangle$. Such basic examples already display rich and interesting behavior. By means of a

Date: August 2, 2001.

The authors gratefully acknowledge partial support from National Science Foundation grant DMS-9803518.

$\langle \gamma \rangle$ -invariant function, we show that recurrent geodesic rays (which are the images of nonproper maps) must lie in one of two codimension-one submanifolds. Further, since not every geodesic ray in those submanifolds is recurrent, we give a characterization (in terms of their directing vectors) of recurrent geodesic rays in cylinders. The notion of a geodesic ray spiralling around a closed geodesic is introduced.

The most interesting examples are *Margulis spacetimes*, when Γ is a free purely hyperbolic discrete subgroup of the isometry group of $\mathbb{A}^{2,1}$, that is a *Schottky group*. Drumm [2, 3] showed that every noncocompact discrete subgroup of $\mathrm{SO}(2, 1)$ admits proper affine deformations. Nonclosed birecurrent geodesics can only be found when the rank of the fundamental group is greater than one.

The main results of this paper may also be found in [1].

2. GEOMETRY OF MINKOWSKI SPACETIME

2.1. Minkowski 2+1-Spacetime. Let $\mathbb{R}^{2,1}$ denote a three-dimensional real vector space equipped with the standard symmetric bilinear form of signature $(2, 1)$:

$$\mathbb{B}(\mathbf{x}, \mathbf{y}) = x_1y_1 + x_2y_2 - x_3y_3,$$

where $\mathbf{x} = (x_1, x_2, x_3)$ and $\mathbf{y} = (y_1, y_2, y_3)$. A vector $\mathbf{v} \in \mathbb{R}^{2,1}$ is *spacelike* (resp. *timelike*, *lightlike*) if $\mathbb{B}(\mathbf{v}, \mathbf{v}) > 0$ (resp. $\mathbb{B}(\mathbf{v}, \mathbf{v}) < 0$, $\mathbb{B}(\mathbf{v}, \mathbf{v}) = 0$). (Lightlike vectors are also called null.)

Denote by $\mathbb{A}^{2,1}$ the affine space modeled on $\mathbb{R}^{2,1}$: for every $p \in \mathbb{A}^{2,1}$, the tangent space

$$\mathbb{A}_p^{2,1} = \{q - p : q \in \mathbb{A}^{2,1}\}$$

is endowed with the bilinear form $\mathbb{B}(\cdot, \cdot)$. Clearly, $\mathbb{A}_p^{2,1} \cong \mathbb{R}^{2,1}$.

Any line in $\mathbb{A}^{2,1}$ can be described as $p + \mathbb{R}\mathbf{v}$, where $p \in \mathbb{A}^{2,1}$ and $\mathbf{v} \in \mathbb{R}^{2,1}$. Two lines $p + \mathbb{R}\mathbf{v}$, $q + \mathbb{R}\mathbf{w}$ are *parallel* if $\mathbf{v} = k\mathbf{w}$ for some $k \neq 0$. The line $p + \mathbb{R}\mathbf{v}$ is called *spacelike*, *timelike* or *lightlike* according to the causal character of \mathbf{v} .

The set of non-spacelike vectors, with the origin removed, has two connected components. A choice of component is a *time orientation* on $\mathbb{R}^{2,1}$. We will adopt the standard time orientation: a non-spacelike vector $\mathbf{v} = (v_1, v_2, v_3)$ is *future-pointing* if $v_3 > 0$ and *past-pointing* otherwise.

The *Lorentz-orthogonal plane* of $\mathbf{v} \in \mathbb{R}^{2,1}$ at p is the set of all vectors based at p which are *Lorentz-orthogonal* to \mathbf{v} :

$$\begin{aligned} p + \mathbf{v}^\perp &= \{q \in \mathbb{A}^{2,1} : \mathbb{B}(q - p, \mathbf{v}) = 0\} \\ &= \{p + \mathbf{x} : \mathbb{B}(\mathbf{x}, \mathbf{v}) = 0, \mathbf{x} \in \mathbb{R}^{2,1}\}. \end{aligned}$$

2.2. Isometries. An *affine isometry* of $\mathbb{A}^{2,1}$ is an affine transformation γ whose linear part preserves the bilinear form $\mathbb{B}(\cdot, \cdot)$. Thus the linear part of an affine isometry of $\mathbb{A}^{2,1}$ lies in $O(2, 1)$. The isometry group of $\mathbb{A}^{2,1}$ is denoted $\text{Isom}(\mathbb{A}^{2,1})$. The connected component of the identity of $O(2, 1)$, denoted $\text{SO}(2, 1)^0$, consists of those linear isometries which preserve orientation and time orientation.

An affine isometry is said to be *hyperbolic* if its linear part is hyperbolic: that is, it is an element of $\text{SO}(2, 1)^0$ which has three real distinct eigenvalues.

If $g \in \text{SO}(2, 1)^0$ is hyperbolic, then its eigenvalues are $\lambda < 1 < \lambda^{-1}$, for some positive $\lambda \in \mathbb{R}$. The λ - and λ^{-1} -eigendirections are lightlike and the 1-eigendirection is spacelike.

Let $g \in \text{SO}(2, 1)^0$ be a hyperbolic isometry with smallest eigenvalue $\lambda < 1$. Set $\mathbf{x}^+(g)$, $\mathbf{x}^-(g)$ to be future-pointing eigenvectors of λ^{-1} , λ , respectively, normalized so that $\|\mathbf{x}^\pm(g)\| = 1$, where $\|\cdot\|$ denotes Euclidean length.

Choose $\mathbf{x}^0(g)$ to be the unique spacelike 1-eigenvector, satisfying $\mathbb{B}(\mathbf{x}^0(g), \mathbf{x}^0(g)) = 1$, such that the *null frame*

$$\{\mathbf{x}^0(g), \mathbf{x}^-(g), \mathbf{x}^+(g)\}$$

is a positively oriented basis. For hyperbolic $\gamma \in \text{Isom}(\mathbb{A}^{2,1})$ with linear part g , set:

$$\{\mathbf{x}^0(\gamma), \mathbf{x}^-(\gamma), \mathbf{x}^+(\gamma)\} = \{\mathbf{x}^0(g), \mathbf{x}^-(g), \mathbf{x}^+(g)\}.$$

The following facts are well known:

Lemma 2.1. *Let $\gamma \in \text{Isom}(\mathbb{A}^{2,1})$ be hyperbolic. Then there exists a line $C_\gamma \subset \mathbb{A}^{2,1}$, parallel to $\mathbf{x}^0(\gamma)$, which is invariant under the action of γ . Moreover, γ acts by translation on C_γ . Furthermore, C_γ is the unique γ -invariant line if and only if γ acts freely on $\mathbb{A}^{2,1}$.*

Proof. Since γ is affine, it acts on N , the space of lines parallel to $\mathbf{x}^0(\gamma)$. Observe that N is isomorphic to the Lorentz-orthogonal plane $\mathbf{x}^0(\gamma)^\perp$, which, in turn, is isomorphic to two-dimensional Minkowski space. The eigenvalues of the induced action of γ are each different from 1, therefore this action has a fixed point.

Let $C_\gamma \subset \mathbb{A}^{2,1}$ be the line parallel to $\mathbf{x}^0(\gamma)$ corresponding to this fixed point in N ; clearly, C_γ is invariant under the action of γ : if $p \in C_\gamma$,

$$\gamma(p) = p + \alpha \mathbf{x}^0(\gamma), \tag{1}$$

where $\alpha \in \mathbb{R}$. Any other point q on C_γ can be written as $p + t\mathbf{x}^0(\gamma)$, for some $t \in \mathbb{R}$. Since $\mathbf{x}^0(\gamma)$ is fixed by γ :

$$\gamma(q) = \gamma(p + t\mathbf{x}^0(\gamma)) = \gamma(p) + t\mathbf{x}^0(\gamma) = q + \alpha \mathbf{x}^0(\gamma).$$

Thus γ acts by translation on C_γ .

Finally, γ fixes a point p if and only if the line $p + \mathbb{R}\mathbf{x}^0(\gamma)$ is pointwise fixed, so C_γ is the unique γ -invariant line if and only if γ acts freely. \square

Definition 2.2. For a hyperbolic isometry $\gamma \in \text{Isom}(\mathbb{A}^{2,1})$, set $\alpha(\gamma)$ to be the parameter α in Equation (1). Then α , a signed measure of Lorentzian length, is called the *Margulis invariant*.

This function was introduced by Margulis [4, 5] to show that an affine group whose linear part is a Schottky group could act properly discontinuously on $\mathbb{A}^{2,1}$. Since $\alpha(\gamma)$ is the translation factor along C_γ , the proof of Lemma 2.1 implies:

Lemma 2.3. *Let $\gamma \in \text{Isom}(\mathbb{A}^{2,1})$ be a hyperbolic isometry. Then γ acts freely on $\mathbb{A}^{2,1}$ if and only if $\alpha(\gamma) \neq 0$.*

\square

2.3. Stable and unstable planes.

Definition 2.4. Let $\gamma \in \text{Isom}(\mathbb{A}^{2,1})$ be hyperbolic and let $p \in C_\gamma$. The planes

$$\begin{aligned} E_\gamma^+ &= p + \langle \mathbf{x}^0(\gamma), \mathbf{x}^+(\gamma) \rangle \\ E_\gamma^- &= p + \langle \mathbf{x}^0(\gamma), \mathbf{x}^-(\gamma) \rangle, \end{aligned}$$

are respectively called the *weak-unstable plane* and the *weak-stable plane* of γ .

Note that $E_\gamma^+ = E_{\gamma^{-1}}^-$, $E_\gamma^+ \cap E_\gamma^- = C_\gamma$, and E_γ^\pm is γ -invariant.

Consider the orbit of a point q in $\mathbb{A}^{2,1}$, under the action of γ . We can write

$$q = p + A^+\mathbf{x}^+(g) + A^-\mathbf{x}^-(g),$$

where $p \in C_\gamma$ and $A^\pm \in \mathbb{R}$. Thus for every n ,

$$\gamma^n(q) = p + n\alpha(\gamma)\mathbf{x}^0(g) + A^+\lambda^{-n}\mathbf{x}^+(g) + A^-\lambda^n\mathbf{x}^-(g).$$

If $A^+ = 0$, the orbit converges towards C_γ . When $A^+ \neq 0$, the sequence approaches E^+ , but eventually leaves every compact set intersecting the weak-unstable plane.

2.4. Geodesics. Let $M = \mathbb{A}^{2,1}/\Gamma$, where $\Gamma < \text{Isom}(\mathbb{A}^{2,1})$ acts properly on $\mathbb{A}^{2,1}$. Then M is a complete Lorentz manifold and its fundamental group is isomorphic to Γ . Let $\pi : \mathbb{A}^{2,1} \rightarrow M$ denote the quotient projection.

A *geodesic* in M is a nonconstant affine map $l : \mathbb{R}^+ \rightarrow M$, that is, the composition $\pi \circ \tilde{l}$ where $\tilde{l} : \mathbb{R} \rightarrow \mathbb{A}^{2,1}$ is a nonconstant affine map. The *reverse* of a geodesic l is the geodesic $-l$ defined by

$$-l(t) := l(-t).$$

A *geodesic ray* in M is a nonconstant affine map $l : \mathbb{R}^+ \rightarrow M$, that is, the composition $\pi \circ \tilde{l}$ where $\tilde{l} : \mathbb{R}^+ \rightarrow \mathbb{A}^{2,1}$ is a nonconstant affine map. The *forward ray* of a geodesic l is the restriction of l to \mathbb{R}^+ and the *backward ray* of l is the forward ray of $-l$.

A geodesic ray l is *parallel* to a line, vector or ray, if its direction vector \mathbf{v} (defined by $l(t) = q + t\mathbf{v}$) is parallel to the line, vector or ray.

A geodesic l is *periodic* if for some $T > 0$, and all $t > 0$,

$$l(t + T) = l(t).$$

We will often identify a periodic geodesic with its image, which is an immersed S^1 in M . The *cylinder associated to l* is the covering space $M_l \rightarrow M$ induced by the equivalence class of l in $\pi_1(M, l(0))$.

The following conditions are equivalent:

- l is a proper map: for every compact $K \subset M$, the inverse image $l^{-1}(K) \subset \mathbb{R}^+$ is compact;
- For every increasing sequence $t_k \rightarrow +\infty$, the sequence $l(t_k)$ has no accumulation point;
- l is not periodic and the image $l(\mathbb{R}^+ \cup \{0\})$ is closed.

A geodesic ray l is *recurrent* if the mapping $l : \mathbb{R}^+ \rightarrow M$ is not proper. Equivalently, its image $l(\mathbb{R}^+)$ has compact closure in M .

For example, a closed geodesic is recurrent.

A geodesic ray r *spirals around a periodic geodesic l* if for every neighborhood N of l , there exists $T = T(N) > 0$ such that $r(t) \in N$ for $t > T$. In particular, such a geodesic is recurrent.

A geodesic l is *birecurrent* if both its forward ray and its backward ray recur. A geodesic l *bispirals* if both its forward ray and backward ray spiral around closed geodesics.

3. CYLINDERS

We now classify geodesics in quotients $\mathbb{A}^{2,1}/\langle\gamma\rangle$, where γ is a hyperbolic affine isometry of $\mathbb{A}^{2,1}$. the cyclic group it generates. In this section M will denote the quotient $\mathbb{A}^{2,1}/\langle\gamma\rangle$.

The *stable surface* $M^- = \pi(E_\gamma^-)$ and the *unstable surface* $M^+ = \pi(E_\gamma^+)$ are each diffeomorphic to an annulus to which M deformation retracts. Similarly M^+, M^- each deformation retract to the *core geodesic* $M^0 = M^+ \cap M^-$, which is the unique closed geodesic in M .

Lemma 3.1. *A recurrent geodesic ray $\mathbb{R}^+ \rightarrow M$ lies in either M^+ or M^- .*

Proof. Let p be an arbitrary point on the invariant line C_γ . For every point $q \in \mathbb{A}^{2,1}$, write

$$q = p + A^0 \mathbf{x}^0(\gamma) + A^- \mathbf{x}^-(\gamma) + A^+ \mathbf{x}^+(\gamma)$$

where $A^-, A^+, A^0 \in \mathbb{R}$ and $\{\mathbf{x}^-(\gamma), \mathbf{x}^+(\gamma), \mathbf{x}^0(\gamma)\}$ is the null frame associated to γ . Then

$$\gamma(q) = p + (A^0 + \alpha(\gamma))\mathbf{x}^0(\gamma) + \lambda A^- \mathbf{x}^-(\gamma) + \lambda^{-1} A^+ \mathbf{x}^+(\gamma) + .$$

and

$$\mathbb{B}(\gamma(q) - p, \mathbf{x}^\pm(\gamma)) = \mathbb{B}(\gamma(q - p), \mathbf{x}^\pm(\gamma)) = \lambda^{\pm 1} \mathbb{B}(q - p, \mathbf{x}^\pm(\gamma)).$$

Thus the quadratic function $\tilde{f} : \mathbb{A}^{2,1} \rightarrow \mathbb{R}$ given by

$$\tilde{f}(q) = \mathbb{B}(q - p, \mathbf{x}^-(\gamma))\mathbb{B}(q - p, \mathbf{x}^+(\gamma)) \quad (2)$$

is independent of the choice of $q \in C_\gamma$, and is $\langle \gamma \rangle$ -invariant. Define $f : M \rightarrow \mathbb{R}$ as $f := \pi \circ \tilde{f}$.

Suppose that $l : \mathbb{R}^+ \rightarrow M$ is a recurrent geodesic ray. Write $l = \pi \circ \tilde{l}$ where

$$\begin{aligned} \tilde{l} : \mathbb{R}^+ &\rightarrow \mathbb{A}^{2,1} \\ t &\mapsto q + tv \end{aligned} \quad (3)$$

is the corresponding geodesic ray in $\mathbb{A}^{2,1}$. Then

$$(f \circ l)(t) = (\tilde{f} \circ \tilde{l})(t) = a + bt + ct^2 \quad (4)$$

where p is an arbitrary point on C_γ and

$$\begin{aligned} a &= f(l(0)) \\ b &= \mathbb{B}(q - p, \mathbf{x}^+(\gamma))\mathbb{B}(v, \mathbf{x}^-(\gamma)) + \mathbb{B}(q - p, \mathbf{x}^-(\gamma))\mathbb{B}(v, \mathbf{x}^+(\gamma)) \\ c &= \mathbb{B}(v, \mathbf{x}^+(\gamma))\mathbb{B}(v, \mathbf{x}^-(\gamma)). \end{aligned}$$

Unless $c = 0$, the function $f \circ l : \mathbb{R}^+ \rightarrow \mathbb{R}$ in (4) is quadratic and tends to $\pm\infty$ as $t \rightarrow +\infty$. If $c = 0$ but $b \neq 0$, then $f \circ l$ is a nonconstant affine function, also tending to $\pm\infty$ as $t \rightarrow +\infty$. In either case $f \circ l$ defines a proper map $\mathbb{R}^+ \rightarrow \mathbb{R}$, contradicting recurrence of l . Thus $f \circ l$ is constant, that is $b = c = 0$.

Since $c = 0$, at least one of $\mathbb{B}(v, \mathbf{x}^+(\gamma)), \mathbb{B}(v, \mathbf{x}^-(\gamma))$ must vanish. Thus \tilde{l} is parallel to either E^- or E^+ respectively. If it is parallel to both, then v is parallel to $\mathbf{x}^0(\gamma)$. We postpone the discussion of this case to the end of the proof.

Suppose that \tilde{l} is parallel to E^+ , but not E^- . Thus $\mathbb{B}(\mathbf{v}, \mathbf{x}^+(\gamma)) = 0$ but $\mathbb{B}(\mathbf{v}, \mathbf{x}^-(\gamma)) \neq 0$. Then $b = 0$ implies that $\mathbb{B}(q - p, \mathbf{x}^+(\gamma)) = 0$. Therefore $\tilde{l} : \mathbb{R}^+ \longrightarrow E^+$.

In the same fashion, if \tilde{l} is parallel to E^- , but not E^+ , then $\tilde{l}(\mathbb{R}^+) \subset E^-$.

Finally, suppose that \tilde{l} is parallel to $\mathbf{x}^0(\gamma)$, but lies on neither E^- nor E^+ . Let

$$U = \mathbb{A}^{2,1} - (E^- \cup E^+).$$

Consider the quotient space N of $\mathbb{A}^{2,1}$ by the one-dimensional foliation parallel to C_γ . The restriction of the quotient map $\Pi : \mathbb{A}^{2,1} \longrightarrow N$ to U is a $\langle \gamma \rangle$ -equivariant mapping with respect to a *proper* action of $\langle \gamma \rangle \cong \mathbb{Z}$ on $U' = \Pi(U)$. Specifically, U' is the complement of two intersecting lines in the plane N and $\langle \gamma \rangle$ acts by hyperbolic linear maps with discrete orbits.

Thus the composition

$$\begin{aligned} \langle \gamma \rangle \times \{q\} &\longrightarrow \mathbb{A}^{2,1} \xrightarrow{\Pi} N \\ (\gamma^n, q) &\longmapsto \gamma^n q \end{aligned}$$

is proper. Consequently the mapping

$$\begin{aligned} \langle \gamma \rangle \times \mathbb{R}^+ &\longrightarrow \mathbb{A}^{2,1} \\ (\gamma^n, t) &\longmapsto \gamma^n(q + t\mathbf{v}) \end{aligned}$$

is proper. Therefore $l : \mathbb{R}^+ \longrightarrow M$ is proper, a contradiction. \square

Lemma 3.2. *Suppose l is a geodesic ray as defined in (3). Suppose one of the following conditions hold:*

(i) $l \subset M^+$ and

$$\mathbb{B}(\alpha(\gamma)\mathbf{x}^0(\gamma), \mathbf{v}) > 0,$$

(ii) $l \subset M^-$ and

$$\mathbb{B}(\alpha(\gamma)\mathbf{x}^0(\gamma), \mathbf{v}) < 0.$$

Then for every $p \in C_\gamma$, $\pi(p) \subset M^0$ is an accumulation point of l .

Proof. Write $l = \pi \circ \tilde{l}$ as above. We first show (i); the proof of (ii) is completely analogous. Suppose then that $l \subset M^+$, so that \tilde{l} lies in the weak-unstable plane E^+ . Let \tilde{V} be a neighborhood of p such that $\tilde{V} \cap E^+$ is bounded by $p \pm k^0\mathbf{x}^0(\gamma) \pm k^+\mathbf{x}^+(\gamma)$. Then:

$$\gamma^n \tilde{V} \cap E^+ = p + (\alpha(\gamma) \pm k^0)\mathbf{x}^0(\gamma) \pm k^+ \lambda^{-n} \mathbf{x}^+(\gamma).$$

Thus the γ^n -translates of \tilde{V} are dilated in the $\mathbf{x}^+(\gamma)$ -direction at the rate of λ^{-n} , whereas the $\mathbf{x}^+(\gamma)$ -coefficient of \tilde{l} grows linearly. Thus for

large enough n , every $\gamma^n(\tilde{V})$ intersects \tilde{l} and p is an accumulation point of l . \square

The restriction of γ to E^+ is represented by the affine map

$$\gamma^+ : \begin{bmatrix} x \\ y \end{bmatrix} \mapsto \begin{bmatrix} 1 & 0 \\ 0 & \lambda^{-1} \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} + \begin{bmatrix} \alpha \\ 0 \end{bmatrix}$$

where $0 < \lambda < 1$ and $\alpha \neq 0$.

Apply the coordinate change

$$\begin{bmatrix} x \\ y \end{bmatrix} \longrightarrow \begin{bmatrix} x \\ \eta \end{bmatrix}$$

where

$$\begin{aligned} \eta(x, y) &= \lambda^{x/\alpha} y \\ y(x, \eta) &= \lambda^{-x/\alpha} \eta. \end{aligned}$$

The action of $\langle \gamma \rangle$ in (x, η) -coordinates is given by horizontal translation by $n\alpha$:

$$(\gamma^+)^n : \begin{bmatrix} x \\ \eta \end{bmatrix} \mapsto \begin{bmatrix} x + n\alpha \\ \eta \end{bmatrix}.$$

This defines a γ -invariant diffeomorphism of the stable surface

$$\begin{aligned} \xi : M^+ &\longrightarrow \mathbb{R}/\alpha\mathbb{Z} \times \mathbb{R} \\ \begin{bmatrix} x \\ y \end{bmatrix} &\mapsto \begin{bmatrix} x \pmod{\alpha\mathbb{Z}} \\ \eta(x, y) \end{bmatrix}. \end{aligned}$$

A geodesic in the unstable plane falls into one of two categories, depending on whether it is parallel to the eigenvector $\mathbf{x}^+(\gamma)$ or not. If the geodesic is parallel to $\mathbf{x}^+(\gamma)$, then we call it a *vertical geodesic*.

Theorem 3.3. *A geodesic ray l (as defined by (3)) in M is recurrent if and only if one of the following holds:*

- (i) $l \subset M^+$ and $\mathbb{B}(\alpha(\gamma)\mathbf{x}^0(\gamma), \mathbf{v}) > 0$
- (ii) $l \subset M^-$ and $\mathbb{B}(\alpha(\gamma)\mathbf{x}^0(\gamma), \mathbf{v}) < 0$.

Proof. We shall show (i). Since in (x, η) -coordinates, γ acts by translation along the x -axis, a geodesic ray is proper if and only if its η -coordinate is unbounded. Let $\beta : \mathbb{R}^+ \rightarrow M^+$ be a geodesic in the unstable surface and choose a lift $\tilde{\beta} : \mathbb{R}^+ \rightarrow E^-$.

If $\tilde{\beta}$ is vertical, it admits the following representation:

$$\tilde{\beta}(t) = \begin{bmatrix} x_0 \\ y_0 \end{bmatrix} + t \begin{bmatrix} 0 \\ c \end{bmatrix}.$$

In the (x, η) -coordinate system:

$$\xi \circ \tilde{\beta}(t) = \begin{bmatrix} x_0 \pmod{\alpha\mathbb{Z}} \\ \eta_0 + c't \end{bmatrix}$$

where

$$\begin{aligned} \eta_0 &= \lambda^{x_0/\alpha} y_0 \\ c' &= \lambda^{x_0/\alpha} c. \end{aligned}$$

As $t \rightarrow +\infty$, the η -coordinate grows without bound and thus the geodesic ray $\xi \circ \tilde{\beta}(t)$ does not recur.

If β is not vertical, then

$$\tilde{\beta}(t) = \begin{bmatrix} x_0 \\ y_0 \end{bmatrix} + t \begin{bmatrix} 1 \\ m \end{bmatrix}$$

where $m \in \mathbb{R}$ is the *slope* of $\tilde{\beta}$. Then

$$\eta \circ \tilde{\beta}(t) = e^{\mu t} (y'_0 + tm')$$

where

$$\begin{aligned} y'_0 &= \lambda^{x_0/\alpha} y_0 \\ m' &= \lambda^{x_0/\alpha} m \\ \mu &= \log(\lambda)/\alpha \end{aligned}$$

The η -coordinate is bounded exactly when $\mu < 0$, that is, when $\alpha < 0$ (since $0 < \lambda < 1$).

Otherwise $e^{\mu t} > 1$ and β does not recur. □

(This theorem follows from a similar result in [1], using the fact that the weak-stable plane of a hyperbolic isometry is the weak-unstable plane of its inverse.)

Figure 1 illustrates the proof of the theorem, by showing the orbit of a recurrent geodesic in (x, η) -coordinates. The geodesic $\beta \subset \mathcal{K}^+$ may cross the closed geodesic M^0 . In that case β crosses M^0 transversely, and then spirals back towards M^0 , intersecting itself infinitely many times.

If $m = 0$, then β never intersects itself, but approaches M^0 asymptotically from one side. (See Figure 2.)

Corollary 3.4. The only birecurrent geodesic in M is the closed geodesic M^0 .

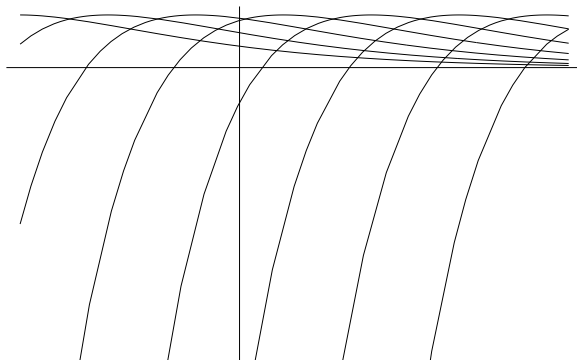


FIGURE 1. The orbit of a recurrent geodesic on the unstable surface in (x, η) -coordinates: the horizontal line is M^0 , which all recurrent geodesics in M^+ approach asymptotically.

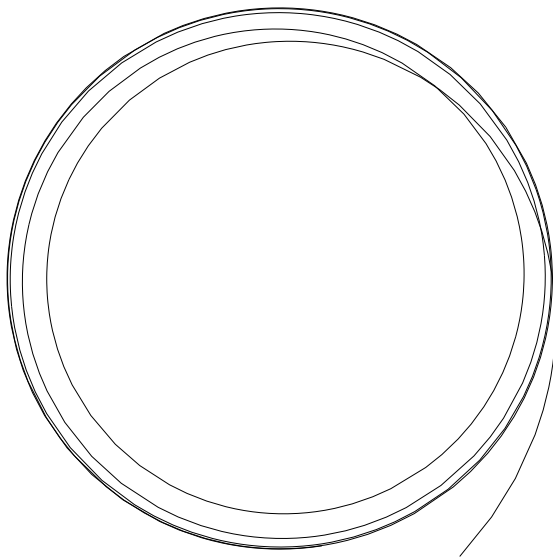


FIGURE 2. The behavior of a ray spiralling towards a closed geodesic.

4. RECURRENT GEODESICS IN FLAT LORENTZ 3-MANIFOLDS

Let M be a Margulis spacetime; recall that its fundamental group is purely hyperbolic. Let l_1, l_2 be (oriented) closed geodesics in M . Then there exists a birecurrent geodesic l whose forward ray spirals towards l_1 and whose backward ray spirals towards l_2 . Such geodesics correspond to equivalence classes of arcs a whose endpoints lie on l_1

and l_2 , under the equivalence relation defined by homotopies keeping the endpoints on l_1 and l_2 respectively (a homotopy relative to l_1 and l_2).

Choose a basepoint $x \in M$ and the corresponding universal covering space $\tilde{M} \rightarrow M$ and developing map $\text{dev} : \tilde{M} \rightarrow \mathbb{A}^{2,1}$. Join l_1, l_2 to x by arcs a_1, a_2 respectively such that the composition $a_1 a_2^{-1}$ is homotopic to a by a homotopy relative to l_1 and l_2 . Let γ_i be the holonomy of the based loops corresponding to l_i , $i = 1, 2$. Thus l_i lifts to C_{γ_i} in the universal cover.

For $\gamma \in \Gamma$, denote its associated stable and unstable surfaces in M by M_γ^\pm . Every geodesic in M is the projection of a geodesic in the cylinder M_{C_γ} , since it is a covering space. In particular, Theorem 3.3 implies:

Proposition 4.1. Let l be a geodesic which lies in $M_{\gamma_i}^+$. Then the forward ray of l spirals around l_i .

Theorem 4.2. If a geodesic $l \subset M$ lies in either M_γ^+ or M_γ^- , where $\gamma \in \Gamma$, then l is recurrent.

□

We obtain birecurrent geodesics in the following manner. Consider:

$$l = M_{\gamma_1}^+ \cap M_{\gamma_2}^+.$$

Write its inverse image \tilde{l} in $\mathbb{A}^{2,1}$ as $\tilde{l}(t) = p + tv$. Then:

$$\mathbb{B}(v, \alpha(\gamma_1)x^0(\gamma_1))\mathbb{B}(v, \alpha(\gamma_2)x^0(\gamma_2)) < 0.$$

Each condition of Theorem 3.3 applies, one in each direction of \tilde{l} , implying that l is birecurrent.

We obtain a similar result for the intersection of stable surfaces, since $E_\gamma^- = E_{\gamma^{-1}}^+$.

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