Waveform design and quantum detection matched filtering

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Outline and collaborators

- 1. CAZAC waveforms
- 2. Frames
- 3. Matched filtering and related problems
- 4. Quantum detection
- 5. Analytic methods to construct CAZAC waveforms

Collaborators: Matt Fickus (frame force), Andrew Kebo (quantum detection), Joseph Ryan and Jeff Donatelli (software).

CAZAC Waveforms

Constant Amplitude Zero Autocorrelation (CAZAC) Waveforms

A *K*-periodic waveform $u : \mathbb{Z}_K = \{0, 1, \dots, K-1\} \rightarrow \mathbb{C}$ is CAZAC if |u(m)| = 1, $m = 0, 1, \dots, K-1$, and the *autocorrelation*

$$A_u(m) = \frac{1}{K} \sum_{k=0}^{K-1} u(m+k)\overline{u}(k) \text{ is 0 for } m = 1, \dots, K-1.$$

The crosscorrelation of $u, v : \mathbb{Z}_K \to \mathbb{C}$ is

$$C_{u,v}(m) = \frac{1}{K} \sum_{k=0}^{K-1} u(m+k)\overline{\nu}(k)$$
 for $m = 0, 1, \dots, K-1$.

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- $u \text{ CAZAC} \iff \text{DFT of } u \text{ is CAZAC}.$
- User friendly software: http://www.math.umd.edu/~jjb/cazac

Rationale for CAZAC waveforms

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- CA allows transmission at peak power. (The system does not have to deal with the suprise of greater than expected amplitude.)
- Distortion amplitude variations can be detected using CA. (With CA amplitude variations during transmission due to additive noise can be theoretically eliminated at the receiver without distorting message.)
- A sharp unique peak in A_u is important because of distortion and interference in received waveforms, *e.g.*, in radar and communications–more later.

Examples of CAZAC Waveforms

$$\begin{split} &K = 75: u(x) = \\ &(1,1,1,1,1,1,e^{2\pi i \frac{1}{15}},e^{2\pi i \frac{2}{15}},e^{2\pi i \frac{1}{5}},e^{2\pi i \frac{4}{15}},e^{2\pi i \frac{1}{3}},e^{2\pi i \frac{7}{15}},e^{2\pi i \frac{3}{5}},\\ &e^{2\pi i \frac{11}{15}},e^{2\pi i \frac{13}{15}},1,e^{2\pi i \frac{1}{5}},e^{2\pi i \frac{2}{5}},e^{2\pi i \frac{3}{5}},e^{2\pi i \frac{4}{5}},1,e^{2\pi i \frac{4}{15}},e^{2\pi i \frac{8}{15}},e^{2\pi i \frac{4}{5}},\\ &e^{2\pi i \frac{16}{15}},e^{2\pi i \frac{1}{3}},e^{2\pi i \frac{2}{3}},e^{2\pi i},e^{2\pi i \frac{4}{3}},e^{2\pi i \frac{5}{3}},1,e^{2\pi i \frac{2}{5}},e^{2\pi i \frac{4}{5}},e^{2\pi i \frac{4}{5}},e^{2\pi i \frac{4}{5}},e^{2\pi i \frac{4}{5}},\\ &e^{2\pi i \frac{16}{15}},e^{2\pi i \frac{7}{15}},e^{2\pi i \frac{14}{15}},e^{2\pi i \frac{7}{5}},e^{2\pi i \frac{28}{15}},e^{2\pi i \frac{1}{3}},e^{2\pi i \frac{13}{15}},e^{2\pi i \frac{7}{5}},e^{2\pi i \frac{29}{15}},\\ &e^{2\pi i \frac{37}{15}},1,e^{2\pi i \frac{3}{5}},e^{2\pi i \frac{6}{5}},e^{2\pi i \frac{9}{5}},e^{2\pi i \frac{12}{5}},1,e^{2\pi i \frac{2}{3}},e^{2\pi i \frac{4}{3}},e^{2\pi i \cdot 2},e^{2\pi i \frac{8}{3}},\\ &e^{2\pi i \frac{3}{15}},e^{2\pi i \frac{16}{15}},e^{2\pi i \frac{38}{15}},e^{2\pi i \frac{49}{15}},1,e^{2\pi i \frac{4}{5}},e^{2\pi i \frac{4}{3}},e^{2\pi i \frac{12}{5}},e^{2\pi i \frac{16}{5}},\\ &1,e^{2\pi i \frac{13}{15}},e^{2\pi i \frac{26}{15}},e^{2\pi i \frac{13}{5}},e^{2\pi i \frac{13}{5}},e^{2\pi i \frac{13}{5}},e^{2\pi i \frac{49}{15}},1,e^{2\pi i \frac{49}{15}},e^{2\pi i \frac{19}{15}},e^{2\pi i \frac{47}{15}},e^{2\pi i \frac{47}{15}},e^{2\pi i \frac{61}{15}},\\ &1,e^{2\pi i \frac{13}{15}},e^{2\pi i \frac{26}{15}},e^{2\pi i \frac{13}{5}},e^{2\pi i \frac{52}{15}},e^{2\pi i \frac{13}{3}},e^{2\pi i \frac{19}{15}},e^{2\pi i \frac{19}{15}},e^{2\pi i \frac{47}{15}},e^{2\pi i \frac{61}{15}},\\ &1,e^{2\pi i \frac{13}{15}},e^{2\pi i \frac{26}{15}},e^{2\pi i \frac{13}{5}},e^{2\pi i \frac{52}{15}},e^{2\pi i \frac{19}{15}},e^{2\pi i \frac{19}{15}},e^{2\pi i \frac{47}{15}},e^{2\pi i \frac{61}{15}},\\ &1,e^{2\pi i \frac{13}{15}},e^{2\pi i \frac{26}{15}},e^{2\pi i \frac{13}{5}},e^{2\pi i \frac{52}{15}},e^{2\pi i \frac{19}{15}},e^{2\pi i \frac{19}{15}},e^{2\pi i \frac{47}{15}},e^{2\pi i \frac{61}{15}},\\ &1,e^{2\pi i \frac{13}{15}},e^{2\pi i \frac{26}{15}},e^{2\pi i \frac{13}{15}},e^{2\pi i \frac{52}{15}},e^{2\pi i \frac{19}{15}},e^{2\pi i \frac{19}{15}},e^{2\pi i \frac{47}{15}},e^{2\pi i \frac{61}{15}},\\ &1,e^{2\pi i \frac{13}{15}},e^{2\pi i \frac{13}{15}},e^{2\pi i \frac{13}{15}},e^{2\pi i \frac{13}{15}},e^{2\pi i \frac{61}{15}},\\ &1,e^{2\pi i \frac{13}{15}},e^{2\pi i \frac{13}{15}},e^{2\pi i \frac{13}{15}},e^{2\pi i \frac{13}{15}},e^{2\pi i \frac{13}{15}},e^{$$

Autocorrelation of CAZAC K = 75



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Finite ambiguity function

Given *K*-periodic waveform, $u : \mathbb{Z}_K \to \mathbb{C}$ let $e_m(n) = e^{\frac{-2\pi i m n}{K}}$.

• The ambiguity function of u, $A : \mathbb{Z}_K \times \mathbb{Z}_K \to K$ is defined as

$$A_u(j,k) = C_{u,ue_k}(j) = \frac{1}{K} \sum_{m=0}^{K-1} u(m+j)\overline{u(m)}e^{\frac{2\pi imk}{K}}$$

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• Analogue ambiguity function for $u \leftrightarrow U$, $||u||_2 = 1$, on \mathbb{R} :

$$A_u(t,\gamma) = \int_{\widehat{\mathbb{R}}} U(\omega - \frac{\gamma}{2}) \overline{U(\omega + \frac{\gamma}{2})} e^{2\pi i t(\omega + \frac{\gamma}{2})} d\omega$$
$$= \int u(s+t) \overline{u(s)} e^{2\pi i s \gamma} ds.$$

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- Standard Doppler frequency shift problem: Construct a statistic to determine unknown Doppler frequency shift.
 Do this for multiple frequencies.
- Provide rigorous justification for CAZAC simulations associated with the Doppler tolerance question and frequency shift problem.

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- There are unresolved "arithmetic" complexities which are affected by waveform structure and length.
- Noise analysis is ongoing.

Doppler Statistic



Doppler Statistic





 $\sum_{j=0}^{k-1} |C_{u,ue_k}(j)|^2 = 1$

Frames

Redundant signal representation

• Given $H = \mathbb{R}_d$ or $H = \mathbb{C}^d$, $N \ge d$. $\{x_n\}_{n=1}^N \subseteq H$ is a finite unit norm tight frame (FUN-TF) if each $||x_n|| = 1$ and, for each $x \in H$,

$$x = \frac{d}{N} \sum_{n=1}^{N} \langle x, x_n \rangle x_n.$$

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• $\{x_n\}_{n=1}^N \subseteq H$ is an *A-tight frame* if $\{x_n\}_{n=1}^N$ spans *H* and $A\|x\|^2 = \sum_{n=1}^N |\langle x, x_n \rangle|^2$ for each $x \in H$.

Recent applications of FUN-TFs

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- Multiple description coding [Goyal, Heath, Kovačević, Strohmer, Vetterli]

Properties and examples of FUN-TFs

Frames give redundant signal representation to compensate for hardware errors, to ensure numerical stability, and to minimize the effects of noise.

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Properties and examples of FUN-TFs

- Frames give redundant signal representation to compensate for hardware errors, to ensure numerical stability, and to minimize the effects of noise.
- Thus, if certain types of noises are known to exist, then the FUN-TFs are constructed using this information.
- Orthonormal bases, vertices of Platonic solids, kissing numbers (sphere packing and error correcting codes) are FUN-TFs.

DFT FUN-TFs

■ $N \times d$ submatrices of the $N \times N$ DFT matrix are FUN-TFs for \mathbb{C}^d . These play a major role in finite frame $\Sigma \Delta$ -quantization.

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$$N = 8, d = 5 \qquad \frac{1}{\sqrt{5}} \begin{bmatrix} * & * & \cdot & * & * & * & \cdot \\ * & * & \cdot & * & * & * & * & \cdot \\ * & * & \cdot & * & * & * & * & \cdot \\ * & * & \cdot & * & * & * & * & \cdot \\ * & * & \cdot & * & * & * & * & \cdot \\ * & * & \cdot & * & * & * & * & \cdot \\ * & * & \cdot & * & * & * & * & \cdot \end{bmatrix}$$
$$x_m = \frac{1}{5} (e^{2\pi i \frac{m}{8}}, e^{2\pi i \frac{m^2}{8}}, e^{2\pi i \frac{m^5}{8}}, e^{2\pi i \frac{m^5}{8}}, e^{2\pi i \frac{m^5}{8}}, e^{2\pi i \frac{m^5}{8}}, e^{2\pi i \frac{m^5}{8}})$$
$$m = 1, \dots, 8.$$

Sigma-Delta Super Audio CDs - but not all authorities are fans.

Frame force

The frame force: $F: S^{d-1} \times S^{d-1} \setminus D \to \mathbb{R}^d$ is defined as $F(a,b) = \langle a,b \rangle (a-b), S^{d-1}$ is the unit sphere in \mathbb{R}^d .

 \checkmark F is a (central) conservative force field.
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Total potential energy for the frame force of $\{x_n\}_{n=1}^N \subseteq S^{d-1}$:

$$P = \sum_{m=1}^{N} \sum_{n=1}^{N} |\langle x_m, x_n \rangle|^2$$

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- Compute these frames.

Multifunction vector-valued frame waveforms

Problem: Construct, code, and implement (user-friendly) N-periodic waveforms $(N \ge d)$

$$u: \mathbb{Z}_N \to S^{d-1} \subseteq \mathbb{R}^d \text{ (or } \mathbb{C}^d),$$

 $n \to u_n = (u_n(1), u_n(2), \dots, u_n(d)), n = 0, 1, \dots, N-1$
which are FUN-TFs (for redundant signal representation)
and CAZAC (zero or low correlation off dc), i.e.,

$$x = \frac{d}{N} \sum_{n=0}^{N-1} \langle x, u_n \rangle u_n \text{ and } A_u(m) = \frac{1}{N} \sum_{j=0}^{N-1} \langle u_{m+j}, u_j \rangle = 0,$$

 $m=1,\ldots N-1.$

The following are recent applications of FUN-TFs.

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- Grassmannian "min-max" waveforms [Calderbank, Conway, Sloane, et al., Kolesar, B]
- Grassmannian analysis gives another measure of the crosscorrelation. A FUN frame $\{u_n\}_{n=1}^N \subseteq H$ is *Grassmannian* if $\max_{k \neq l} |\langle u_k, u_l \rangle| = \inf \max_{k \neq l} |\langle x_k, x_l \rangle|$, where the infimum is over all FUN frames.

Matched Filtering

Processing



Natural problems associated with multifunction frame waveforms (1)

 Implement FUN-TF ΣΔ A/D converters to take advantage of proven improved error estimates for linear reconstruction over PCM and comparable to MSE-PCM. (MSE-PCM is based on Bennett's white noise assumption which is not always valid. With consistent reconstruction, and its added numerical complexity, MSE-PCM is comparable to FUN-TF MSE-ΣΔ.)

Natural problems associated with multifunction frame waveforms (2)

Distinguish multiple frequencies and times (ranges) in the ambiguity function,

$$A("t", "\gamma") = \int_{\hat{R}} U(\omega) (\sum \alpha_j \overline{U(\omega + \gamma_j)} e^{2\pi i t_j \omega}) d\omega,$$

by means of multifunction frame waveforms.

Natural problems associated with multifunction frame waveforms (3)

• Compute optimal 1-tight frame CAZAC waveforms, $\{e_n\}_{n=1}^N$, using quantum detection error:

$$P_e = \min_{\{e_n\}} (1 - \sum_{i=1}^N \rho_n |\langle u_n, e_n \rangle|^2), \quad \sum_{n=1}^N \rho_n = 1, \rho_n > 0,$$

where $\{u_n\}_{n=1}^N \subseteq S^{d-1}$ is given. This is a multifunction matched filtering.

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- If $v(t) = a\overline{u(t-t_0)}$ for some t_0 , then

$$C_{v,u}(t_0) = \sup_t |C_{v,u}(t)|.$$

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In digital case, CAZACs arise since travel time depends on crosscorrelation peak, and sharp peaks obviate distortion and interference in received waveform.

• QM formulates concept of measuring a dynamical quantity (e.g., position of an electron in \mathbb{R}^3) and the probability p that the outcome is in $U \subseteq \mathbb{R}^3$.

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- Positive operator-valued measure (POVM) gives rise to p.
- In $H = \mathbb{C}^d$, **POVMs** and 1-tight frames are equivalent.
- Given $\{u_n\}_{n=1}^N \subseteq S^{d-1}$. Compute/construct a 1-tight frame minimizer $\{e_n\}_{n=1}^N$ of quantum detection (QD) error P_e .

Outline of multifunction matched filtering algorithm

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- Transfer tight frames for \mathbb{C}^d to ONBs in \mathbb{C}^N (Naimark point of view and essential for computation).
- Show that the QD error is a potential energy function of frame force in C^N.
- Use the orthogonal group and the Euler-Lagrange equation for the potential P_e to *compute* equations of motion and a minimal energy solution $\{e_n\}_{n=1}^N$.

Quantum Detection

Positive-operator-valued measures

Let \mathcal{B} be a σ -algebra of sets of X. A *positive operator-valued measure* (POM) is a function $\Pi: \mathcal{B} \to \mathcal{L}(H)$ such that

- 1. $\forall U \in \mathcal{B}, \ \Pi(U)$ is a positive self-adjoint operator,
- 2. $\Pi(\emptyset) = 0$ (zero operator),
- 3. \forall disjoint $\{U_i\}_{i=1}^{\infty} \subset \mathcal{B}$ and $x, y \in H$,

$$\left\langle \Pi\left(\bigcup_{i=1}^{\infty}U_{i}\right)x,y\right\rangle = \sum_{i=1}^{\infty}\langle \Pi(U_{i})x,y\rangle,$$

- 4. $\Pi(X) = I$ (identity operator).
- A POM Π on \mathcal{B} has the property that given any fixed $x \in H$, $p_x(\cdot) = \langle x, \Pi(\cdot)x \rangle$ is a measure on \mathcal{B} . (Probability if ||x|| = 1).

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A POM Π on \mathcal{B} has the property that given any fixed $x \in H$, $p_x(\cdot) = \langle x, \Pi(\cdot)x \rangle$ is a measure on \mathcal{B} . (Probability if ||x|| = 1).

A dynamical quantity Q gives rise to a measurable space (X, \mathcal{B}) and POM. When measuring Q, $p_x(U)$ is the probability that the outcome of the measurement is in $U \in \mathcal{B}$.

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Suppose the state of the electron is given by $x \in H$ with unit norm. Then the probability that the electron is found to be in the region $U \in \mathcal{B}$ is given by

$$p(U) = \langle x, \Pi(U)x \rangle = \int_U |x(t)|^2 dt.$$

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Clear that Π satisfies conditions (1)-(3) for a POM. Since *F* is Parseval, we have condition (4) ($\Pi(X)x = \sum_{i \in X} \langle x, e_i \rangle e_i = x$). Thus Π defines a POM.

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- Conversely, let (X, \mathcal{B}) be a measurable space with corresponding POM Π for a *d*-dimensional Hilbert space *H*. If *X* is countable then there exists a subset $K \subseteq \mathbb{Z}$, a Parseval frame $\{e_i\}_{i \in K}$, and a disjoint partition $\{B_j\}_{j \in X}$ of *K* such that for all $j \in X$ and $y \in H$,

$$\Pi(j)y = \sum_{i \in B_j} \langle y, e_i \rangle e_i.$$

Quantum detection for finite frames

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- Our goal is to determine what state the system is in by performing a "good" measurement. That is, we want to construct a POM with outcomes $X = \mathbb{Z}_N$ such that if the state of the system is x_i for some $1 \le i \le N$, then

$$p_{x_i}(j) = \langle x_i, \Pi(j) x_i \rangle \approx \begin{cases} 1 & \text{ if } i = j \\ 0 & \text{ if } i \neq j \end{cases}$$

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Since $\langle x_i, \Pi(i)x_i \rangle$ is the probability of a successful detection of the state x_i , then the probability of a detection error is given by

$$P_e = 1 - \sum_{i=1}^{N} \rho_i \langle x_i, \Pi(i) x_i \rangle.$$

Quantum detection problem

If we construct our POM using Parseval frames, the error becomes

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Quantum detection problem: Given a unit normed set $\{x_i\}_{i=1}^N \subset H$ and positive weights $\{\rho_i\}_{i=1}^N$ that sum to 1. Construct a Parseval frame $\{e_i\}_{i=1}^N$ that minimizes

$$P_e = 1 - \sum_{i=1}^{N} \rho_i |\langle x_i, e_i \rangle|^2$$

over all N-element Parseval frames. ($\{e_i\}_{i=1}^N$ exists by a compactness argument.)

Naimark theorem

Naimark Theorem Let *H* be a *d*-dimensional Hilbert space and let $\{e_i\}_{i=1}^N \subset H, N \ge d$, be a Parseval frame for *H*. Then there exists an *N*-dimensional Hilbert space *H'* and an orthonormal basis $\{e'_i\}_{i=1}^N \subset H'$ such that *H* is a subspace of *H'* and

$$\forall i = 1, \dots, N, \ \mathcal{P}_H e'_i = e_i,$$

where \mathcal{P}_H is the orthogonal projection $H' \to H$.

Given $\{x_i\}_{i=1}^N \subset H$ and a Parseval frame $\{e_i\}_{i=1}^N \subset H$. If $\{e'_i\}_{i=1}^N$ is its corresonding orthonormal basis for H', then, for all i = 1, ..., N, $\langle x_i, e_i \rangle = \langle x_i, e'_i \rangle$.

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- Minimizing P_e over all N-element Parseval frames for H is equivalent to minimizing P_e over all N-element orthonormal bases for H'.
- Thus we simplify the problem by minimizing P_e over all *N*-element orthonormal sets in H'.

Quantum detection error as a potential

Treat the error term as a potential.

$$P = P_e = \sum_{i=1}^{N} \rho_i (1 - |\langle x_i, e'_i \rangle|^2) = \sum_{i=1}^{N} P_i.$$

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$$P_i = \rho_i (1 - |\langle x_i, e'_i \rangle|^2).$$

For $H' = \mathbb{R}^N$, we have the relation,

$$||e'_i - x_i||^2 = 2 - 2\langle x_i, e'_i \rangle$$

where we have used the fact that $||e'_i|| = ||x_i|| = 1$. We can rewrite the potential P_i as

$$P_i = \rho_i \left(1 - \left[1 - \frac{1}{2} \| x_i - e'_i \|^2 \right]^2 \right).$$

A central force corresponds to quantum detection error

Given P_i , define the function $p_i: \mathbb{R} \to \mathbb{R}$ by

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Thus P_i is a potential corresponding to a central force in the following way:

$$-xf_i(x) = p'_i(x) = 2\rho_i \left(1 - \frac{1}{2}x^2\right) x$$
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Hence, the force $F_i = -\nabla P_i$ is

$$F_i(x_i, e'_i) = f_i(||x_i - e'_i||)(x_i - e'_i) = -2\rho_i \langle x_i, e'_i \rangle (x_i - e'_i),$$

a multiple of the frame force! The total force is given by

$$F = \sum_{i=1}^{N} F_i$$

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- The equilibrium position of the points $\{e'_i\}_{i=1}^N$ is the position where all the forces produce no net motion. In this situation, the potential *P* is minimized.
- For the remainder, let $\{e'_i\}_{i=1}^N$ be an ONB for \mathbb{R}^N that minimizes P. Recall that $\{e'_i\}_{i=1}^N$ exists by compactness. The *quantum detection problem* is to construct or compute $\{e'_i\}_{i=1}^N$.

A parameterization of O(N)

Consider the orthogonal group

 $O(N) = \{ \Theta \in GL(N, \mathbb{R}) : \Theta^{\tau} \Theta = I \}.$

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Since O(N) is an N(N-1)/2-dimensional smooth manifold, we can locally parameterize O(N) by N(N-1)/2 variables, i.e., $\Theta = \Theta(q_1, \ldots, q_{N(N-1)/2})$ for each $\Theta \in O(N)$.

Hence, for all $\theta \in O(N)$ there is a surjective diffeomorphism b_{θ}



for relatively compact neighborhoods $\mathcal{U}_{\theta} \subseteq O(N)$ and $\mathcal{U} \subseteq \mathbb{R}^{N(N-1)/2}$, $\theta \in \mathcal{U}_{\theta}$.

A parameterization of ONBs

• Let $\{w_i\}_{i=1}^N$ be the standard ONB for $H' = \mathbb{R}^N$: $w_i = (0, \dots, 0, \underbrace{1}_{i=1}, 0, \dots, 0)$.

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Since any two orthonormal sets are related by an orthogonal transformation, we can smoothly parameterize an orthonormal set $\{e_i\}_{i=1}^N$ with N elements by N(N-1)/2 variables, i.e.,

$$\{e_i(q_1,\ldots,q_{N(N-1)/2})\}_{i=1}^N = \{\Theta(q_1,\ldots,q_{N(N-1)/2})w_i\}_{i=1}^N \subset H'.$$

where for all $\Psi \in O(N)$, $W_i(\Psi) = \Psi w_i$.

$$e_i(\vec{q}) = e_i(q_1, \dots, q_{N(N-1)/2}) = W_i \circ b_{\theta}^{-1}(\vec{q}) = (b_{\theta}^{-1}(\vec{q}))w_i \in \mathbb{R}^N.$$

Lagrangian dynamics on ${\cal O}(N)$

We now convert the frame force F acting on the orthonormal set $\{e_i\}_{i=1}^N$ into a set of equations governing the motion of the parameterization points $\vec{q}(t) = (q_1(t), \dots, q_{N(N-1)/2}(t))$, see (1). We define the Lagrangian L and total energy E defined for $\vec{q}(t)$ by:

$$L = T - P_e, \quad E = T + P_e,$$

where

$$T = \frac{1}{2} \sum_{j=1}^{N(N-1)/2} \left(\frac{d}{dt} q_j(t)\right)^2.$$

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Using the Euler-Lagrange equations for the potential P_e

$$\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{q}_j}\right) - \frac{\partial L}{\partial q_j} = 0,$$

we obtain the equations of motion

(1)
$$\frac{d^2}{dt^2}q_j(t) = -2\sum_{i=1}^N \rho_i \langle x_i, e_i(\vec{q}(t)) \rangle \left\langle x_i, \frac{\partial e_i}{\partial q_j}(\vec{q}(t)) \right\rangle.$$

Point of view

Choose $\vec{q}' \in \mathbb{R}^{N(N-1)/2}$ such that $e_i(\vec{q}') = e'_i \in \mathbb{R}^N$ for all i = 1, ..., N.

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Define $\tilde{q} : \mathbb{R} \to \mathbb{R}^{N(N-1)/2}$ such that $\tilde{q}(t) = \vec{q}'$ (a constant function).

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Remark The definition of \tilde{q} and equation (1) introduce *t* into play for solving the quantum detection problem.



It can be shown that

Proof Theorem Denote by $\vec{q}(t) = (q_1(t), \dots, q_{N(N-1)/2}(t))$ a solution of the equations of motion that minimizes the energy *E* and denote by \mathcal{P}_H the orthogonal projection from *H'* into *H*. Then $\vec{q}(t)$ is a constant solution and the set of vectors

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- **Important** For the set of O(N).
- So we need only consider parameterizing SO(N).
- **P** Theorem A minimum energy solution, a minimizer of P_e , satisfies the expression

$$\sum_{i=1}^{N} \rho_i \langle x_i, e_i \rangle \left\langle x_i, \frac{\partial e_i}{\partial q_j} \right\rangle = 0.$$

Numerical problems

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With the parameterization of SO(N), the error P_e is a smooth function of the variables $(q_1, \ldots, q_{N(N-1)/2})$, that is,

$$P_e(q_1,\ldots,q_{N(N-1)/2}) = 1 - \sum_{i=1}^N \rho_i |\langle x_i, e_i(q_1,\ldots,q_{N(N-1)/2}) \rangle|^2.$$

A conjugate gradient method can be used to find the minimum values of P_e .

Analytical methods

Problem: Let $p = \{p_k\}_{k \in \mathbb{Z}}$ be positive definite, *i.e.*, for any finite set *F* ⊆ Z and any $\{c_j\}_{j \in F} \subseteq \mathbb{C}$:

$$\sum_{j,k\in F} c_j \bar{c}_k p(j-k) \ge 0$$

Suppose p = 0 on a given $F \subseteq \mathbb{Z}$. When can we construct unimodular $u : \mathbb{Z} \to \mathbb{C}$ such that:

$$p(k) = \lim_{N \to \infty} \frac{1}{2N+1} \sum_{|j| \le N} u(j+k) \overline{u(j)}?$$

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• This is the same problem for \mathbb{Z} that we have been addressing for \mathbb{Z}_N in the one-dimensional CAZAC case.
Iterative Generalized Harmonic Analysis (GHA of Wiener)

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Finite approximation and software as with algebraic CAZACs.

