## Numerical Simulation of Dynamic Stall

AMSC 663-664 Proj ect
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## Introduction

- Aim
- To study the dynamic stalling of rotorcraft blade sections
- Dynamic Stall
- Major factor in helicopter performance
- Involves unsteady aerodynamics (time varying angle of attack and incoming flow velocity)
- Numerical Solution
- Navier Stokes equations w/ turbulence model


## Flow around Airfoil

- Airfoil: 2D cross-section of wing/rotor
- Exhibits basic characteristics of flow around wings and other lifting surfaces
- Provides an estimate of the lift and drag for actual flows
- Numerical Solution of flow around airfoils
- 2D Euler or Navier Stokes equations
- To study nature of flow over wings, away from edges
- Step towards building a 3D flow solver


## Inviscid Flows

- Governed by the Euler Equations
- Based on the conservation of mass, momentum and energy
- Neglects viscosity and heat conduction
- Relevance
- High speed flows, away from solid surfaces
- Negligible cross-derivatives of flow velocity
- Provides a good estimate of lift but not drag


## Viscous Flows

- Governed by the Navier Stokes equations
- Effect of viscous forces and heat conduction on the momentum and energy equations
- Required to compute drag and flow separation
- Drag = pressure drag + shear forces
- Flow separation at high angles of attack (pockets of reverse flows)
- Computationally more expensive
- Requires very fine mesh spacing near body surface to capture boundary layer


## Governing Equations

Euler Equations (Inviscid) $\quad \frac{\partial \mathbf{u}}{\partial t}+\frac{\partial f(\mathbf{u})}{\partial \mathrm{x}}+\frac{\partial \mathbf{g}(\mathbf{u})}{\partial \mathrm{y}}=0$
Navier Stokes Equations $\frac{\partial \mathbf{u}}{\partial \mathrm{t}}+\frac{\partial \mathbf{f}(\mathbf{u})}{\partial \mathrm{x}}+\frac{\partial \mathbf{g}(\mathbf{u})}{\partial \mathrm{y}}=\frac{\partial \boldsymbol{f}_{\mathbf{v}}(\mathbf{u})}{\partial \mathrm{x}}+\frac{\partial \mathbf{g}(\mathbf{u})}{\partial \mathrm{y}}$

$$
\begin{aligned}
& \tau_{\mathrm{xx}}=(\lambda+2 \mu) \frac{\partial \mathrm{u}}{\partial \mathrm{x}}+\lambda \frac{\partial \mathrm{v}}{\partial \mathrm{y}} ; \tau_{\mathrm{yy}}=\lambda \frac{\partial \mathrm{u}}{\partial \mathrm{x}}+(\lambda+2 \mu) \frac{\partial \mathrm{v}}{\partial \mathrm{y}} ; \tau_{\mathrm{xy}}=\tau_{\mathrm{yx}}=\mu\left(\frac{\partial \mathrm{u}}{\partial \mathrm{y}}+\frac{\partial \mathrm{v}}{\partial \mathrm{x}}\right)
\end{aligned}
$$

## Constitutive Relations

- Equation of State $E=\frac{p}{\gamma-1}+\frac{1}{2} \rho\left(u^{2}+v^{2}\right)$
- Ideal Gas Law

$$
\mathrm{p}=\rho \mathrm{RT}
$$

- Fluid Properties:
${ }^{\circ} \mu$
Fluid Viscosity Coefficient
- $\lambda$

Bulk Viscosity Coefficient

- k

Thermal Conductivity

- R

Universal Gas Constant

- $\gamma$

Ratio of Heat Capacities

## Finite Volume Formulation

- Based on the Integral Form of the governing equations

$$
\begin{aligned}
& \int_{V} \frac{\partial \mathbf{u}}{\partial \mathrm{t}} \mathrm{dV}+\int_{\partial \mathrm{V}} \mathbf{F} \cdot \hat{\mathbf{n}} \mathrm{dS}=0 \\
& \mathbf{F}(\mathbf{u})=\left[\mathbf{f}(\mathbf{u})-\mathbf{f}_{\mathbf{v}} \mathbf{( u )}\right) \hat{\mathbf{i}}+\left[\mathbf{g}(\mathbf{u})-\mathbf{g}_{\mathbf{v}}(\mathbf{u}) \hat{\mathbf{j}}\right.
\end{aligned}
$$

- Semi-discrete equation

$$
\frac{d \mathbf{u}_{i j}}{d t}+\sum_{\text {faces }} \mathbf{F} \cdot \hat{\mathbf{n}} d S=0 \Rightarrow \frac{d \mathbf{u}_{i j}}{d t}=\boldsymbol{\operatorname { R e s }}(i, j)
$$

## Inviscid Flow Computations

- Upwinded flux reconstruction at the interfaces
- Characteristics based upwinding (solution composed of acoustic and entropy waves)
- $1^{\text {st }}$ Order Roe Scheme
- $2{ }^{\text {nd }}, 3{ }^{\text {rd }}$ Order Essentially Non-Oscillatory schemes
- Explicit Time Marching
- Runge Kutta family of ODE solvers
- $2^{\text {nd }}$ and $3^{\text {rd }}$ order Total Variation Diminishing RK schemes


## Viscous Flow Computations

- Explicit time marching unsuitable
- Extremely small mesh spacing near the surface results in a very restrictive stability limit for time step size
- Implicit Time Stepping: Backward Euler
- Unconditionally stable
- Computationally more intensive per iteration (solution of a system of equations)
- Non-linear system of equations requires linearization in time for each time step
- Time step size limited by accuracy considerations


## Curvilinear Coordinates

- Transformation of the governing equations from Cartesian coordinates ( $x, y$ ) to curvilinear coordinates $(\xi, \eta)$


- Metrics of transformations numerically computed

$$
\mathrm{x} \xi, \mathrm{y} \xi, \mathrm{x}_{\eta}, \mathrm{y}_{\eta} \leftrightarrow \xi_{\mathrm{x}}, \xi_{\mathrm{y}}, \eta_{\mathrm{x}}, \eta_{\mathrm{y}}
$$

## Numerical Scheme

- Governing Equation in curvilinear coordinates

$$
\frac{\partial \hat{\mathbf{u}}}{\partial \mathrm{t}}+\frac{\partial \hat{\mathbf{f}}(\hat{\mathbf{u}})}{\partial \xi}+\frac{\partial \hat{\mathbf{g}}(\hat{\mathbf{u}})}{\partial \eta}=\frac{\partial \hat{\mathbf{f}}_{\mathbf{v}}(\hat{\mathbf{u}})}{\partial \xi}+\frac{\partial \hat{\mathbf{g}}_{\mathbf{v}}(\hat{\mathbf{u}})}{\partial \eta}
$$

- Linearization of flux with respect to time

$$
\mathrm{f}(\mathrm{u}(\mathrm{t}+\Delta \mathrm{t}))=\mathrm{f}(\mathrm{u}(\mathrm{t}))+\mathrm{A} \Delta \mathrm{u}+\mathrm{O}\left(\Delta \mathrm{t}^{2}\right) ; \mathrm{A}=\frac{\partial \mathrm{f}}{\partial \mathrm{u}}
$$

- Implicit formulation (h-Time step)

$$
\begin{align*}
& {\left[\mathrm{I}+\mathrm{h} \partial_{\xi} \mathrm{A}+\mathrm{h} \partial_{\eta} \mathrm{B}\right] \Delta \mathrm{u}=\mathrm{h}\left[\partial_{\xi} \hat{\mathrm{f}}+\partial_{\eta} \hat{\mathrm{g}}\right] ; \mathrm{A}=\frac{\partial \mathrm{f}}{\partial \mathrm{u}}, \mathrm{~B}=\frac{\partial \mathrm{g}}{\partial \mathrm{u}}} \\
& \Rightarrow\left[\mathrm{I}+\mathrm{h} \partial_{\xi} \mathrm{A}\right]\left[\mathrm{I}+\mathrm{h} \partial_{\eta} \mathrm{B}\right] \Delta \mathrm{u}=\mathrm{h}\left[\partial_{\xi} \hat{\mathrm{f}}+\partial_{\eta} \hat{\mathrm{g}}\right] \quad(\mathrm{ADI}) \tag{ADI}
\end{align*}
$$

## Numerical Scheme

- Right-hand side computed using finite difference/finite volume formulations
- Convective flux computed as before
- Dissipative terms computed using second order central differencing
- Left-hand side
- Banded penta-diagonal system
- ADI approximation: two banded tri-diagonal systems
- $\partial_{\xi}(\mathrm{A} \Delta \mathrm{u}), \partial_{\eta}(\mathrm{B} \Delta \mathrm{u})$ computed using $1^{\text {st }}$ order upwind finite differences


## Validation

- Cartesian Meshes
- 2D Riemann Problems
- Oblique Shock Reflection
- Couette Flow
(Inviscid)
(Inviscid)
(Viscous)
- Airfoil computations
- NACA0012 Subsonic
- NACAoo12 Transonic
- RAE2822 Transonic
- NACAoo12 High Angle of Attack
(Inviscid)
(Inviscid)
(Viscous)
(Viscous)


## 2D Riemann Problems


$\downarrow$ Results from Kurganov-Tadmor $L_{\downarrow}$


## Oblique Shock Reflection



## Oblique Shock Pressure Contours and Streamlines

Pressure
3.79303
3.48944
3.18585
2.88226
2.57867
2.27508
1.97149
1.6679
1.36431
1.06072

Exact Solution obtained through Oblique Shock relations

$3^{\text {rd }}$ order ENO $+$
$3^{\text {rd }}$ order TVD
Runge Kutta
Mach 2.9
Inflow

## Couette Flow

- Flow between a stationary plate and a moving plate ( $u=100 \mathrm{~m} / \mathrm{s}$ )
- Steady state is a linear velocity profile
- 2D Cartesian mesh
- Periodic boundary conditions along x boundaries
- No-slip conditions along y-boundaries


Initial, intermediate and steady state velocity profiles for Couette flow

## Airfoil Computations - Domain



C-Type Structured Mesh with outside boundary 20 chords away
Freestream boundary conditions on outer boundary

Magnified view of mesh around airfoil (unit chord)


Curved Wall Boundary Conditions at Airfoil Surface

## NACA0012 Subsonic (Inviscid)

- Coefficient of Pressure

$$
\mathrm{C}_{\mathrm{p}}=\frac{\mathrm{p}-\mathrm{p}_{\infty}}{\rho_{\infty} \mathrm{u}_{\infty}{ }^{2} / 2}
$$

- Results validated with TURNS code



Pressure and streamlines around the airfoil at Mach 0.63 and 2 degrees angle of attack

## NACA0012 Transonic (Inviscid)



Pressure and streamlines around the airfoil at Mach 0.85 and 1 degrees angle of attack

Results validated with TURNS code
Higher order schemes show better shock resolution than $1^{\text {st }}$ order


## RAE2822 Transonic (Viscous)

RAE 2822 Airfoil at Mach 0.75 and 2.8 degrees angle of attack

1st Order Roe's scheme with Euler Backward time stepping



## NACA0012 High Angle of Attack (Viscous)




NACA0012 Airfoil at Mach 0.3 and 13.5 degrees angle of attack

1st $^{\text {st }}$ Order Roe scheme with Euler Backward time stepping

## Conclusions

- Flow solver (2D) for inviscid and viscous flow around airfoils
- Presently validated for steady state computations
- Post-processing required to compute lift, drag and moment coefficients
- Efficient solver for banded tri-diagonal system needed
- Need to incorporate a turbulence model
- Future work - for computation of dynamic stall
- Increase time accuracy by sub-iterations
- Validation for unsteady cases (plunging airfoils, rotating airfoils)


## End!

