

Dual Reciprocity Boundary Element Method for Magma Ocean Simulations

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- ▶ Understanding how crystals settle in a magma ocean is critical for answering questions about Earth's early history
- Experiments have been performed that have given insight into this behavior
- ► The ability to simulate this behavior numerically is not available
- ▶ It is being proposed to use the Dual Reciprocity Method as the basis for a numerical solver to address this problem.

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Physical Motivation

- Earth's early history is marked by a giant impact with a Mars-sized object
- This led to a substantial amount of interior melting followed by rapid crystallization of this 'magma ocean'
- ▶ How this crystallization took place and crystals settled provide insight into the planet's rate of cooling

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Giant Impact



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- ▶ Must solve a coupled Stokes flow and heat equation
- Settling crystals create a free boundary condition which makes traditional methods, such as FEM, a challenge to use
- First order accurate asymptotic solutions exists for limited parameters

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Asymptotic Solution



Algorithm

- ▶ Boundary Element Method almost works
 - ▶ Naturally handles free boundaries
 - Approximation performed on only on the boundary
 - ▶ Cannot handle transient heat equation with source term
- ▶ Dual Reciprocity Method fits the problem
 - ▶ Naturally handles free boundaries
 - Approximation performed on only on the boundary
 - ▶ Can solve transient heat equation with source term

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Given a non-linear PDE

$$\mathcal{D}u = b \tag{1}$$

Rewrite it as a combination of a linear operator \mathcal{L} and non-linear residual operator \mathcal{D}'

$$\mathcal{L}u = b - \mathcal{D}'u =: b' \tag{2}$$

Define b', the combination of original source with the non-linear operator, to be the new source term to the linear PDE

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The Heat Equation

Consider the heat equation

$$\frac{\partial}{\partial t}u + \vec{v} \cdot \nabla u - \Delta u = b \tag{3}$$

Break it down into a linear operator and its residual

$$\mathcal{D} = \frac{\partial}{\partial t} + \vec{v} \cdot \nabla - \Delta \tag{4}$$

$$\mathcal{L} = -\Delta \tag{5}$$

$$\mathcal{D}' = \frac{\partial}{\partial t} + \vec{v} \cdot \nabla \tag{6}$$

$$b' = b - \frac{\partial}{\partial t}u - \vec{v} \cdot \nabla u \tag{7}$$

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Formulating Integral Equations

Form the Boundary Integral Equation using the linear PDE

$$-\Delta u = b' \qquad x \in \Omega \tag{8}$$

$$-\Delta u \cdot w = b' \cdot w \tag{9}$$

$$-\int_{\Omega} \Delta u \cdot w \, \mathrm{d}\Omega = \int_{\Omega} b' \cdot w \, \mathrm{d}\Omega \tag{10}$$

Perform integration by parts twice on the left hand side

$$\int_{\Omega} \Delta u \cdot w \, \mathrm{d}\Omega = \int_{\Gamma} w D u \cdot \vec{n} \, \mathrm{d}\Gamma - \int_{\Omega} D u \cdot D w \, \mathrm{d}\Omega$$
$$= \int_{\Gamma} (w D u - u D w) \cdot \vec{n} \, \mathrm{d}\Gamma + \int_{\Omega} u \Delta w \, \mathrm{d}\Omega$$

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Fundamental Solutions

Choose $w=u^*$ where u^* is the fundamental solution to the adjoint operator \mathcal{L}^*

$$\Delta u^* = -\delta(x - \xi) \tag{11}$$

then

$$u(\xi) = \int_{\Gamma} (u^* Du - u Du^*) \cdot \vec{n} \, d\Gamma + \int_{\Omega} b' u^* \, d\Omega$$
(12)
for $\xi \in \Omega$

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Approximation of Source

- ▶ Converted the linear term to the boundary
- ▶ Still need to convert source term to boundary
 - Given particular solution to $\mathcal{L}u' = b'$ is know this is possible
 - Particular solution is almost never known
- Dual Reciprocity approximates the source term by a linear combination of functions for which a particular solution is known

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Properties of Approximation

Approximate the source term b'

$$b' = \sum_{q=1}^{\mathcal{N}} f^q \alpha_q \tag{13}$$

Choose basis function f^q such that the particular solution u^q is known

$$\Delta u^q = f^q, \tag{14}$$

Then the original PDE becomes

$$-\Delta u = \sum_{q=1}^{\mathcal{N}} f^q \alpha_q \tag{15}$$

Substitution

Substituting for b' the integral equation becomes

$$u(\xi) = \int_{\Gamma} (u^* D u - u D u^*) \cdot \vec{n} \, \mathrm{d}\Gamma + \sum_{q=1}^{\mathcal{N}} \alpha_q \int_{\Omega} f^q u^* \mathrm{d}\Omega \qquad (16)$$

Substituting for f^p

$$u(\xi) = \int_{\Gamma} (u^* D u - u D u^*) \cdot \vec{n} \, \mathrm{d}\Gamma - \sum_{q=1}^{\mathcal{N}} \alpha_q \int_{\Omega} \Delta u^q u^* \mathrm{d}\Omega \qquad (17)$$

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The Second Reciprocity Equation

Using the second $\operatorname{Gau}\beta$ theorem

$$\int_{\Omega} (\Delta u^q u^* - \Delta u^* u^q) \, \mathrm{d}\Omega = \int_{\Gamma} (D u^q u^* - D u^* u^q) \cdot \vec{n} \, \mathrm{d}\Gamma \qquad (18)$$

Which, after remembering the choice of u^* , becomes

$$\int_{\Omega} \Delta u^q u^* \, \mathrm{d}\Omega = -u^q(\xi) + \int_{\Gamma} (Du^q u^* - Du^* u^q) \cdot \vec{n} \, \mathrm{d}\Gamma \qquad (19)$$

Substituting into the integral equation yields...

Dual Reciprocity

The Dual Reciprocity equation

$$u(\xi) = \int_{\Gamma} (u^* Du - u Du^*) \cdot \vec{n} \, d\Gamma + \sum_{q=1}^{\mathcal{N}} \left(u^q(\xi) - \int_{\Gamma} (Du^q u^* - Du^* u^q) \cdot \vec{n} \, d\Gamma \right) \alpha^q$$

for $\xi\in\Omega$

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Computation

- ▶ Equations are non-dimensionalized
- ▶ Discretized using the method of collocation $(\xi \to \Gamma)$
- ▶ Yields a linear system of N equations and N unknowns where N is the number of discretized elements
- ▶ All known and unknown values are on the boundary
- BEM formulation provides a method for approximating values in the region once unknown values are found

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- ▶ Linear systems involve dense and asymmetric matrices
 - ► *LU*-Decomposition with pivoting for stable backward and forward substitution
 - ► Fast Multipole Method demands extra upfront work but generates a sparse linear system (optional)
- ▶ Weakly and Strongly singular integrals
 - Appear from collocation $(\xi \to \Gamma)$
 - Iterative methods
 - Radial approximation techniques



Runtime

- Dominated by calculation of matrix and solving the associated linear system
 - $\mathcal{O}(N^2)$ with dense matrices
 - $\mathcal{O}(N \log N)$ with FMM
 - $\triangleright \mathcal{O}(N)$ upfront work
 - $\mathcal{O}(N \log N)$ compute entries of sparse matrix
 - $\triangleright \mathcal{O}(N)$ solve using sparse solver such as GMRES

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Platform

- ▶ Software: Fortran 90
 - Very fast execution
 - Object Orientated
 - Existing BEM Poisson and Stokes flow solver
- ▶ Hardware: Single core x86 architecture
 - ▶ Readily available
 - ▶ Local access to 32 node x86 cluster
 - Scalable from netbooks to supercomputers

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- ▶ Use initial conditions and parameters from literature
 - Modify these conditions and parameters as experimentation demands
- Parameters for asymptotic approximations for steady-state Stokes flow and heat problems

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Validation

- ▶ Test steady-state code
 - Validate new Poisson and Stokes portion
 - Test Poisson code in isoviscous environment
 - ▶ Test Stokes code in isothermal environment
 - Use asymptotic approximations for non-isoviscous and non-isothermal problem

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Validation

► Test Dual Reciprocity code

- Use test functions
 - Create test function v
 - Compute $v_0 = v(t = 0)$ and $\mathcal{D}v =: b$
 - Then v is an exact solution to $\mathcal{D}u = b$ with $u(t = 0) = v_0$

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- Phase I (present early November)
 - Merge Stokes flow and Poisson solver code
- ▶ Phase II (November December)
 - Test and validate steady-state code
 - Optimize code
- Phase III (December early February)
 - Add Dual Reciprocity code
 - Add Fast Multipole Method code (optional)
- Phase IV (February March)
 - Test and validate Dual Reciprocity code
 - Optimize code

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Deliverables

- ▶ Collection of source and compiled libraries
 - Validated and optimized
 - Easily modifiable and modular
 - Robust and reusable
 - ▶ Platform for further research and development
- Reports
 - ► AMSC664 Report
 - Geophysics paper detailing methods and results for physical simulations
- Presentations
 - ► AMSC664 Presentation
 - Presenting paper at 12th International Workshop on Modeling of Mantle Convection and Lithospheric Dynamics in Berlin, Germany, Summer 2011

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Thank You

Questions?

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