

**L27** Today:

Topic 1: Wavelet Transform

~~Topic 2: Uncertainty Inequality~~

# The Continuous Wavelet Transform

Recall from examples of Windowed Fourier Transform (WFT)

$$f(x) = e^{2\pi i \omega_1 x} + e^{2\pi i \omega_2 x} \quad (\text{or, } f(x) = \sin(2\pi \omega_1 x) + \sin(2\pi \omega_2 x))$$

and a window  $g = e^{-\pi x^2}$

⇒ Spectrogram,  $|V_g f(t, \omega)|$  exhibits two <sup>separated</sup> ridges (for  $\omega > 0$ )

one for  $f_1(\omega) = e^{2\pi i \omega_1 x}$  and another for  $f_2(\omega) = e^{2\pi i \omega_2 x}$

if and only if  $|\omega_1 - \omega_2|$  is large enough.

On the other hand, if we perform the WFT with a different window  $g$  then the minimum separation  $|\omega_1 - \omega_2|$  may be different.

A different transform was introduced:

Definition. Fix a function  $\psi: \mathbb{R} \rightarrow \mathbb{R}$  called wavelet. The continuous wavelet transform (CWT) of a function (signal)  $f: \mathbb{R} \rightarrow \mathbb{C}$  with respect to  $\psi$ , denoted  $W_\psi f$  is the function,

$$W_\psi f: (\mathbb{R} \setminus \{0\}) \times \mathbb{R} \rightarrow \mathbb{C}, \quad W_\psi f(a, b) = \frac{1}{\sqrt{|a|}} \int_{-\infty}^{\infty} f(x) \cdot \psi\left(\frac{x-b}{a}\right) dx$$

$$W_{\psi} f(a, b) = \frac{1}{\sqrt{|a|}} \int_{-\infty}^{\infty} f(x) \cdot \psi\left(\frac{x-b}{a}\right) dx. \quad \rightarrow \text{CWT}$$

Compare to:

$$V_g f(t, \omega) = \int_{-\infty}^{\infty} e^{-2\pi i \omega x} f(x) g(x-t) dx. \quad \rightarrow \text{WFT}$$

If  $\psi_{a,b}(x) = \frac{1}{\sqrt{|a|}} \psi\left(\frac{x-b}{a}\right)$ ,  $g_{\omega,t}(x) = e^{2\pi i \omega x} g(x-t)$

Then:  $W_{\psi} f(a, b) = \langle f, \psi_{a,b} \rangle$ ,  $V_g f(t, \omega) = \langle f, g_{\omega,t} \rangle$ .

You can check:

$$\int_{-\infty}^{\infty} |\psi_{a,b}(x)|^2 dx = \frac{1}{|a|} \int_{-\infty}^{\infty} |\psi\left(\frac{x-b}{a}\right)|^2 dx = \int_{-\infty}^{\infty} |\psi(y)|^2 dy = \|\psi\|_2^2$$

$$\int_{-\infty}^{\infty} |g_{\omega,t}(x)|^2 dx = \int_{-\infty}^{\infty} |g(x-t)|^2 dx = \int_{-\infty}^{\infty} |g(x)|^2 dx = \|g\|_2^2$$

For the CWT:  $a \rightarrow$  Scale. (scale shift)  $\rightarrow$  dilation.  
 $b \rightarrow$  time. (time shift)  $\rightarrow$  translation.

$$G = \{ (a, b) : a, b \in \mathbb{R}, a \neq 0 \} \rightarrow (U(a, b) \psi)(x) = \frac{1}{\sqrt{|a|}} \psi\left(\frac{x-b}{a}\right)$$

$U$ : unitary operator

$\rightarrow$  Form a group,  $(G, *) \rightarrow$  NOT A COMMUTATIVE GROUP!

$$(a, b) = (a_1, b_1) * (a_2, b_2) \leftrightarrow U(a, b) = U(a_1, b_1) \cdot U(a_2, b_2)$$

$\Rightarrow$  The Representation Theory of the  $ax + b$  group.

$$\left( \underbrace{\cup(a_1, b_1) \cup (a_2, b_2)}_{\varphi} \right) (x) = \left( \cup(a_1, b_1) \varphi \right) (x) = \frac{1}{\sqrt{|a_1|}} \varphi\left(\frac{x-b_1}{a_1}\right) =$$

$$= \frac{1}{\sqrt{|a_1|}} \frac{1}{\sqrt{|a_2|}} \varphi\left(\frac{\frac{x-b_1}{a_1} - b_2}{a_2}\right) = \frac{1}{\sqrt{|a_1 a_2|}} \varphi\left(\frac{x-b_1 - a_1 b_2}{a_1 a_2}\right) =$$

$$= \frac{1}{\sqrt{|a|}} \varphi\left(\frac{x-b}{a}\right).$$

Hence:  $a = a_1 a_2$   
 $b = b_1 + a_1 b_2$

$$\underbrace{(a_1, b_1) * (a_2, b_2) = (a_1 a_2, b_1 + a_1 b_2)}$$

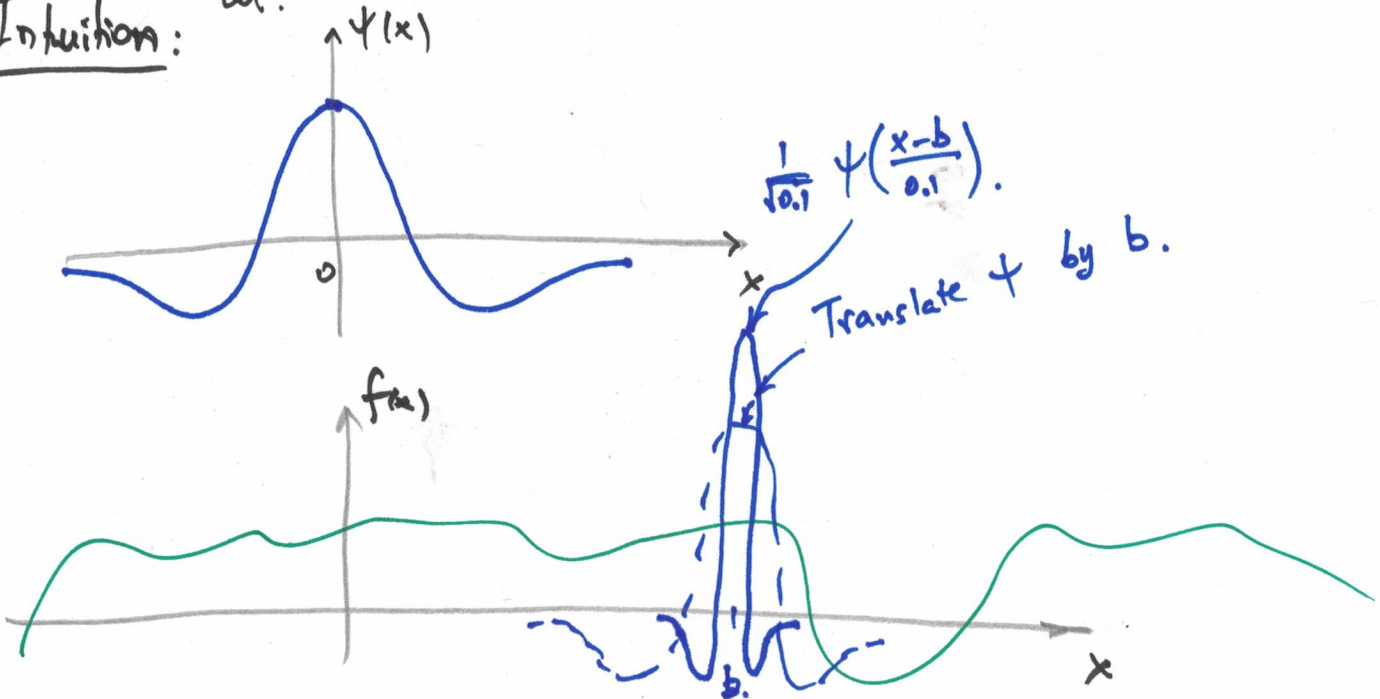
$$(a, b) * (1, 0) = (a \cdot 1, b + a \cdot 0) = (a, b)$$

$$(1, 0) * (a, b) = (1 \cdot a, 0 + 1 \cdot b) = (a, b)$$

$$\Rightarrow (a, b) * (1, 0) = (1, 0) + (a, b) = (a, b).$$

Intuition:

Let:



# How to invert the CWT:

Theorem. Assume  $\psi: \mathbb{R} \rightarrow \mathbb{R}$  the wavelet satisfies:

(1)  $\int_{-\infty}^{\infty} |\psi(x)|^2 dx < \infty, \psi \in L^2(\mathbb{R})$

(2)  $C_\psi := \int_{-\infty}^{\infty} \frac{|\hat{\psi}(\omega)|^2}{|\omega|} d\omega < \infty.$  ...  $\xrightarrow{\text{if } \hat{\psi} \text{ is continuous.}} \hat{\psi}(0) = 0 \rightarrow \int_{-\infty}^{\infty} \psi(x) dx = 0$   
"wavelet"

Then:

(1) If  $f_1, f_2 \in L^2(\mathbb{R})$ , let  $W_\psi f_1, W_\psi f_2$  denote their CWT respectively. Then:

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} W_\psi f_1(a,b) \cdot \overline{W_\psi f_2(a,b)} \frac{1}{a^2} db da = C_\psi \int_{-\infty}^{\infty} f_1(x) \overline{f_2(x)} dx$$

or,

$$\left\langle W_\psi f_1, W_\psi f_2 \right\rangle_{L(\mathbb{R}^2, \frac{1}{a^2} da db)} = C_\psi \cdot \langle f_1, f_2 \rangle$$

Generalized (Plancherel / Parseval Identity).

(2) Let  $f \in L^2(\mathbb{R})$  and  $F(a,b) = W_\psi f(a,b)$ . Then:

$$f(x) = \frac{1}{C_\psi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} F(a,b) \cdot \psi_{a,b}(x) \frac{1}{a^2} da db = \frac{1}{C_\psi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} F(a,b) \frac{1}{\sqrt{|a|}} \psi\left(\frac{x-b}{a}\right) \frac{da db}{a}$$

in  $L^2$ -sense

[T. Folland]

Why:

If  $f_1 = f_2 = f$ . Want: 
$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} |W_{\psi} f(a,b)|^2 \frac{da db}{a^2} = c_{\psi} \int_{-\infty}^{\infty} |f(x)|^2 dx$$

Generalized  
(Plancherel identity)

$$W_{\psi} f(a,b) = \frac{1}{\sqrt{|a|}} \int_{-\infty}^{\infty} f(x) \psi\left(\frac{x-b}{a}\right) dx = \langle f, \psi_{a,b} \rangle = \langle \hat{f}, \hat{\psi}_{a,b} \rangle.$$

Plancherel-Parseval.

$$\hat{\psi}_{a,b}(\omega) = \int_{-\infty}^{\infty} e^{-2\pi i \omega x} \frac{1}{\sqrt{|a|}} \psi\left(\frac{x-b}{a}\right) dx = \int_{-\infty}^{\infty} e^{-2\pi i \omega (ay+b)} \psi(y) \frac{1}{\sqrt{|a|}} |a| dy =$$

$$y = \frac{x-b}{a}; \quad x = ay+b$$

$$dx = |a| dy$$

$$= \sqrt{|a|} e^{-2\pi i b \omega} \int_{-\infty}^{\infty} e^{-2\pi i a \omega y} \psi(y) dy = \sqrt{|a|} e^{-2\pi i b \omega} \hat{\psi}(a\omega).$$

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} |W_{\psi} f(a,b)|^2 \frac{da db}{a^2} = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} |a| \left| \int_{-\infty}^{\infty} \hat{f}(\omega) e^{2\pi i b \omega} \hat{\psi}(a\omega) d\omega \right|^2 \frac{da db}{a^2} =$$

$$= \int_{-\infty}^{\infty} \frac{1}{|a|} \left[ \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{2\pi i b \omega} \underbrace{\left( \hat{f}(\omega) \overline{\hat{\psi}(a\omega)} \right)}_{g(\omega)} d\omega \right]^2 da = \int_{-\infty}^{\infty} \frac{1}{|a|} \int_{-\infty}^{\infty} |\hat{f}(\omega)|^2 |\hat{\psi}(a\omega)|^2 d\omega da$$

$\|g\|_2^2$

$$= \int_{-\infty}^{\infty} |\hat{f}(w)|^2 \left( \int_{-\infty}^{\infty} \frac{1}{|a|} |\hat{\psi}(aw)|^2 da \right) dw$$

But:

$$\int_{-\infty}^{\infty} \frac{1}{|a|} |\hat{\psi}(aw)|^2 da = \int_{-\infty}^{\infty} \frac{1}{\frac{|s|}{|w|}} |\hat{\psi}(s)|^2 \frac{1}{|w|} ds = \int_{-\infty}^{\infty} \frac{1}{|s|} |\hat{\psi}(s)|^2 ds = c_{\psi}$$

$s = a \cdot w$ ,  ~~$ds = a da$~~ ,  $a = \frac{1}{w} s$   
 ~~$da = \frac{1}{w} ds$~~   $da = \frac{1}{|w|} ds$ .

Thus:

$$\iint_{-\infty}^{\infty} |W_{\psi} f(a,b)|^2 \frac{da db}{a^2} = c_{\psi} \int_{-\infty}^{\infty} |\hat{f}(w)|^2 dw = c_{\psi} \int_{-\infty}^{\infty} |f(x)|^2 dx$$

By similar arguments as for WFT (polarization identity & computing the adjoint)  
 → obtain (1), (2).

~~$$f(x) = \frac{1}{c_{\psi}} \iint_{-\infty}^{\infty} W_{\psi} f(a,b) \psi_{a,b}(x) \frac{da db}{a^2} = \left\langle f, \frac{1}{c_{\psi}} \iint_{-\infty}^{\infty} \psi_{a,b}(x) \psi_{a,b}(\cdot) \frac{da db}{a^2} \right\rangle$$

$\downarrow$   
 $\langle f, \psi_{a,b} \rangle$~~

~~$k_x(\cdot) = k(x, \cdot)$  : kernel.~~

→  ~~$f(x) = \langle f, k_x \rangle \Rightarrow$  Reproducing Kernel Hilbert space (RKHS)~~

