Multi-scale hyperbolic relaxation system, kinetic-fluid domain coupling, and singular limit analysis

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Set up of the problem

Jin-Xin hyperbolic relaxation system:

$$\begin{array}{rcl} \partial_t u^{\epsilon} + \partial_x v^{\epsilon} &=& 0, \\ \partial_t v^{\epsilon} + a^2 \partial_x u^{\epsilon} &=& \lambda_{\epsilon}(x) \big(f(u^{\epsilon}) - v^{\epsilon} \big), \end{array}$$

with a two-scale relaxation rate (0 < $\epsilon \ll$ 1):

$$\lambda_{\epsilon}(x) = \left\{ egin{array}{c} 1, \ x < 0, \ rac{1}{\epsilon}, \ x > 0. \end{array}
ight.$$

Domain decomposition method:

- overlap vis non-overlap
- data transmission and derivatives matching
- hamiltonian system vs dissipation system
- steady calculations vs dynamical calculations

Motivation:

- Analysis of Asymptotic Preserving schemes,
- Kinetic-fluid domain coupling, localized kinetic upscaling
- Related to Milne problem in kinetic theory,

Set up of the problem

Assume a sub-characteristic condition:

$$|f'(u)| < a,$$

with compactly supported initial data $u_0(x) \in L^{\infty}(\mathbb{R}) \cap BV(\mathbb{R})$. For simplest, assume f(0) = 0 and $v_0(x) = f(u_0(x))$. We assume either

$$0 < C_0 \leq f'(u), \quad ext{for all } u \in \mathbb{R},$$

or

$$f'(u) \leq -C_0 < 0$$
, for all $u \in \mathbb{R}$.

Under this assumption, no shock wave can stick to the interface.

A decoupled zero ϵ limit system

When $f'(u(t, 0^+)) < 0$,

$$\begin{cases} \partial_t u + \partial_x f(u) = 0, \quad x > 0, \ t > 0, \\ u(0, x) = u_0(x), \end{cases}$$

and set v(t,x) = f(u(t,x)).

$$\begin{cases} \partial_t u + \partial_x v = 0, & x < 0, \ t > 0, \\ \partial_t v + a^2 \partial_x u = f(u) - v, & v(t, 0^-) = v(t, 0^+) \\ u(0, x) = u_0(x), v(0, x) = v_0(x). \end{cases}$$

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When $f'(u(t, 0^+)) > 0$,

$$\begin{cases} \partial_t u + \partial_x v = 0, & x < 0, \ t > 0, \\ \partial_t v + a^2 \partial_x u = f(u) - v, \\ v(t, 0^-) = f(u(t, 0^-)), \\ u(0, x) = u_0(x), v(0, x) = v_0(x), \end{cases}$$

$$\begin{cases} \partial_t u + \partial_x f(u) = 0, & x > 0, \ t > 0, \\ u(t, 0^+) = u(t, 0^-), \\ u(0, x) = u_0(x), \end{cases}$$

and set v(t,x) = f(u(t,x)).

We will only analyze these two cases.

Coupled system for general flux *f* (characteristic)

$$\begin{cases} \partial_t u + \partial_x v = 0, \ x < 0, \ t > 0 \\ \partial_t v + a^2 \partial_x u = f(u) - v, \\ v(t, 0^-) = f(u(t, 0^+)), \\ u(0, x) = u_0(x), \ v(0, x) = v_0(x) \\ \partial_t u + \partial_x f(u) = 0, \ x > 0, \ t > 0 \\ v(t, x) = f(u(t, x)), \\ u(t, 0^+) \stackrel{BLN}{:=} u(t, 0^-), \\ u(0, x) = u_0(x) \end{cases}$$

Bardos-Lerous-Nédélec condition: for all k between $u(t, 0^-)$ and $u(t, 0^+)$.

$$\operatorname{sgn}(u(t,0^+) - u(t,0^-))(f(u(t,0^+)) - f(k)) \le 0,$$

If f' > 0, then BLN reduced to $u(t, 0^+) = u(t, 0^-)$ If f' < 0, then BLN is always hold. Hence, when $f' \neq 0$, this system is reduced to the previous decoupled system.

Connection profile for general flux f (characteristic)

The connection profile is given by

$$a^{2} \frac{d}{dy} \mathcal{U}(t, y) = f(\mathcal{U}(t, y)) - v(t, 0^{+}), \ y > 0, \quad \mathcal{U}(t, 0) = u(t, 0^{-})$$

and it is possible to connected to a middle value $u(t, 0^m)$ with $f(u(t, 0^m)) = f(u(t, 0^+))$ and there is a stationary shock at the end of profile connecting $u(t, 0^m)$ and $u(t, 0^+)$.

Example:

Take $f(u) = \frac{1}{2}u^2$, a = 1, $u(t, x) \equiv \begin{cases} 1 & \text{for } x \leq \delta \\ -1 & \text{for } x > \delta \end{cases}$, $\delta \geq 0$ $v(t, x) = f(u(t, x), U(t, y) \equiv 1 \text{ is an exact solution to the zero } \epsilon\text{-limit coupled system.}$ Vasseur's result on coupled kinetic and Burgers' eq Perthame-Tadmor kinetic model with a two-scale relaxation rate

$$\partial_t f^{\epsilon} + \xi \partial_x f^{\epsilon} = \lambda_{\epsilon}(x) \big(\mathcal{M} f^{\epsilon} - f^{\epsilon} \big)$$

$$u = \int_{\mathbb{R}} f d\xi$$
, $\mathcal{M}f := \mathcal{M}(u,\xi) := \mathbf{1}_{\{0 \le \xi \le u\}} - \mathbf{1}_{\{u \le \xi \le 0\}}$

compatibility condition

0

$$0 \leq \operatorname{Sign}(\xi) f(\xi) \leq 1$$

For $0 < u^+ < L$, $f_0 \in L^{\infty}(\mathbb{R}, [-L, L])$ Initial data case I:

$$f_0(x,\xi) \geq \left\{ egin{array}{ll} M(u^+,\xi), & ext{ for } x < 0, \ M(-u^+,\xi), & ext{ for } x > 0, \ \ \xi \in [-L,L] \end{array}
ight.$$

Initial data case II:

$$\begin{cases} \text{Supp } f_0 \in [-L, u^+ - \eta], & \text{for } x < 0, \\ f_0 = M(-u^+, \xi), & \text{for } x > 0, \ \xi \in [-L, L] \\ & < \square > < \square > < \square > < > > \end{cases}$$

The coupled zero ϵ limit system

$$\begin{cases} \partial_t f + \xi \partial_x f = \mathcal{M}f - f, \ x < 0, \ \xi \in \mathbb{R} \\ f(t, 0_-, \xi) = F(t, 0_+, \xi), \ \xi \le 0 \\ f(0, x, \xi) = f_0(x, \xi) \end{cases} \\\begin{cases} \xi \partial_y F = \mathcal{M}F - F, \ y > 0, \ \xi \in \mathbb{R} \\ F(t, 0_+, \xi) = f(t, 0_-, \xi), \ \xi \ge 0 \\ \int_{\mathbb{R}} \xi F(t, y, \xi) \ d\xi = \frac{1}{2}u(t, 0_+)^2, \ y > 0 \end{cases} \\\begin{cases} \partial_t u + \frac{1}{2}\partial_x u^2 = 0, \ x > 0, \ t > 0 \\ v(t, x) = f(u(t, x)), \\ u(t, 0^+) \stackrel{BLN}{:=} \sqrt{2 \int_{-\infty}^{\infty} \xi f(t, 0^-, \xi) \ d\xi}, \\ u(0, x) = u_0(x) \end{cases}$$

Vasseur's Results: Well-posedness and stability of the limit system; Convergence to the limit system for the two class of initial data.

Existence and stability for the ϵ system

For any $\epsilon > 0$ and T > 0, there is a unique weak solution $(u^{\epsilon}, v^{\epsilon})$ in $L^1((0, T), L^1(\mathbb{R}))$ and satisfies the following *a priori* estimates:

$$\begin{array}{ll} (i) & \| u^{\epsilon}(t,\cdot) \|_{L^{\infty}(\mathbb{R})} \leq C, & \| v^{\epsilon}(t,\cdot) \|_{L^{\infty}(\mathbb{R})} \leq C; \\ (ii) & TV(v^{\epsilon}(t,\cdot)) \leq C; \\ (iii) & TV_{\{x<0\}}(u^{\epsilon}(t,\cdot)) \leq C(1+T), \\ & TV_{\{x>0\}}(u^{\epsilon}(t,\cdot)) \leq C\left(1+\frac{T}{\epsilon}\right); \\ (iv) & \| u^{\epsilon} - f^{-1}(v^{\epsilon}) \|_{L^{1}((0,T);L^{2}(\mathbb{R}^{+}_{x}))} \leq C_{T} \sqrt{\epsilon}; \end{array}$$

for all time T > 0 while

Outline of the proof

- Smoothing the data $u_{0,\delta}(x) = \rho_{\delta} * u_0(x)$, $\lambda_{\epsilon,\delta}(x) = \rho_{\delta} * \lambda_{\epsilon}(x)$. We skip subscripts ϵ, δ below for clearlity in expression.
- Use the characteristic variables,

$$r_{\pm}(t,x) = a \ u(t,x) \pm v(t,x),$$

one has L^1 -contraction:

$$\begin{split} &\int_{|x| < M} \left(|r_{+} - \bar{r}_{+}| + |r_{-} - \bar{r}_{-}| \right) dx \\ &\leq \int_{|x| < M + at} \left(|(r_{+})_{0} - (\bar{r}_{+})_{0}| + |(r_{-})_{0} - (\bar{r}_{-})_{0}| \right) (x) dx. \end{split}$$

for any two classical solutions $(u^{\epsilon,\delta}, v^{\epsilon,\delta})$ and $(\bar{u}^{\epsilon,\delta}, \bar{v}^{\epsilon,\delta})$, vanishing outside the cone $\bigcup_{t\geq 0}[-(M+at), (M+at)]$.

Outline of the proof

• Uniform L^1 estimate for the time derivative $s_{\pm} = a \partial_t u \pm \partial_t v$

$$\begin{cases} \partial_t s_- - a \partial_x s_- = -\lambda_{\epsilon,\delta}(x) \mathcal{R}(s_-, s_+), \\ \partial_t s_+ + a \partial_x s_+ = \lambda_{\epsilon,\delta}(x) \mathcal{R}(s_-, s_+), \end{cases}$$

with

$$\mathcal{R}(s_{-},s_{+}) = f'(u) \frac{s_{-}+s_{+}}{2a} - \frac{s_{+}-s_{-}}{2},$$
 (8)

The sub-characteristic condition implies

$$\partial_{s_-}\mathcal{R}(s_-,s_+) > 0, \ \partial_{s_+}\mathcal{R}(s_-,s_+) < 0$$

$$\begin{split} \int_{|x| < M} \big(|s_{+}(t, x)| + |s_{-}(t, x)| \big) dx \\ &\leq \int_{|x| < M + at} \big(|s_{+}(0, x)| + |s_{-}(0, x)| \big) dx \end{split}$$

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Outline of the proof

- BV estimate of $v^{\epsilon,\delta}$ immediately follows from $\partial_x v = -\partial_t u$.
- Entropy dissipation estimate. For Chen-Levermore-Liu entropy pair $\Phi, \Psi,$

$$\partial_{vv}\Phi(u,v) \ge C, \Phi(0,0) = 0, \partial_{v}\Phi(u,f(u)) = 0, \forall u \in \mathbb{R},$$

$$\partial_{t}\Phi(u,v) + \partial_{x}\Psi(u,v) = \lambda_{\epsilon,\delta}(x)\partial_{v}\Phi(u,v)\Big(f(u) - v\Big).$$

$$\partial_{t}\Phi(u,v) + \partial_{x}\Psi(u,v) = -\lambda_{\epsilon,\delta}(x)\partial_{vv}\Phi(u,\theta_{1})f'(\theta_{2})^{2}|f^{-1}(v) - u|^{2}$$

for some θ_{1} in between v and $f(u)$ and θ_{2} in between u and
 $f^{-1}(v)$

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Strong convergence in $L^1((0, T) \times \mathbb{R})$

For any T > 0, $(u^{\epsilon}, v^{\epsilon})$ converges to (u, v) in $L^1((0, T) \times \mathbb{R})$.

- (u, v) is a weak solution in $(0, T) \times \mathbb{R}^{-}_{x}$.
- *u* is an entropy solution in $(0, T) \times \mathbb{R}^+_x$ and v = f(u) a.e.
- $u_t + v_x = 0$ in $\mathcal{D}'((0, T) \times \mathbb{R})$.
- $(u, v) \in BV((0, T) \times \mathbb{R})$. For a.e. t > 0, there exist left and right traces $u(t, 0^-)$, $u(t, 0^+)$, and $v(t, 0^-) = v(t, 0^+)$.
- L¹ contraction:

$$\begin{aligned} &\frac{1}{2a}||(r_{+}-\bar{r}_{+})(t,.)||_{L^{1}(\mathbb{R}^{-}_{x})}+ \frac{1}{2a}||(r_{-}-\bar{r}_{-})(t,\cdot)||_{L^{1}(\mathbb{R}^{-}_{x})}\\ &+||(u-\bar{u})(t,.)||_{L^{1}(\mathbb{R}^{+}_{x})}\leq \|u_{0}-\bar{u}_{0}\|_{L^{1}(\mathbb{R}_{x})},\end{aligned}$$

and the estimates reduced from ϵ -system.

Matched asymptotic analysis

• Blow-up: $y = \frac{x}{\epsilon}$, $(\mathcal{U}^{\epsilon,\delta}, \mathcal{V}^{\epsilon,\delta})(t, y) = (u^{\epsilon,\delta}, v^{\epsilon,\delta})(t, \epsilon y)$, $a^2 \partial_{\mathcal{V}} \mathcal{U}^{\epsilon,\delta} = \epsilon \lambda_{\epsilon,\delta}(\epsilon \mathbf{v}) (f(\mathcal{U}^{\epsilon,\delta}) - \mathcal{V}^{\epsilon,\delta}) - \epsilon \partial_t \mathcal{V}^{\epsilon,\delta}.$ • For any L > 0, $TV_{(-L,L)}(\mathcal{U}^{\epsilon,\delta}(t,\cdot)) \leq CL$ • Let $\delta \to 0$, then let $\epsilon \to 0$, $(\mathcal{U}^{\epsilon,\delta}(t,\cdot), \mathcal{V}^{\epsilon,\delta}(t,\cdot))$ converges to $(\mathcal{U}(t,\cdot),\mathcal{V}(t,\cdot))$ in $L^1_{loc}(\mathbb{R}_v)$ for any given t>0 and satisfy $\partial_{\mathbf{v}}\mathcal{V} = 0, \quad \mathbf{a}^2 \partial_{\mathbf{v}}\mathcal{U} = \mathbf{1}_{\{\mathbf{v}>0\}}(f(\mathcal{U}) - \mathcal{V}), \text{ in } \mathcal{D}'(-L,L)$ • $\mathcal{U}(t, \cdot), \mathcal{V}(t, \cdot)$ Lipschitz; admit $\mathcal{U}(t, \pm \infty), \mathcal{V}(t, \pm \infty);$ $\mathcal{V}(t, \mathbf{v}) = \mathcal{V}(t, +\infty) = \mathcal{V}(t, -\infty), \ \mathbf{v} \in \mathbb{R};$ $\mathcal{U}(t, y) = \mathcal{U}(t, -\infty), y < 0; f(\mathcal{U}(t, +\infty)) = \mathcal{V}(t, -\infty);$ $a^2 d_v \mathcal{U}(t, v) = f(\mathcal{U}(t, v)) - \mathcal{V}(t, +\infty), \quad v > 0,$

If $\mathcal{U}(t, \cdot)$ is not locally constant for y > 0, then it is strictly monotone for y > 0 and $f'(\mathcal{U}(t, +\infty)) < 0$.

Matched asymptotic analysis

• For a.e. t > 0, V and v perfectly match

$$\mathcal{V}(t,y)=v(t,0^-)=v(t,0^+), \ \ {
m for \ all} \ y\in\mathbb{R}.$$

 \mathcal{U} and u are linked according to

$$\mathcal{U}(t,y) = u(t,0^-), \quad y < 0.$$

Defining $R_{\pm}(t,y) = a\mathcal{U}(t,y) \pm \mathcal{V}(t,y)$, then

$$egin{aligned} &rac{1}{2}\Big(|R_+(t,y)-\ell|-|R_-(t,y)-h(\ell)|\Big)\ &\geq \mathrm{sgn}(u(t,0^+)-k)(f(u(t,0^+))-f(k)), \quad y>0, \end{aligned}$$

for any $\ell \in R$ s.t. $k = (\ell + h(\ell))/2a$ verifies $|k| < ||u_0||_{L^{\infty}(\mathbb{R})}$.

Matched asymptotic analysis

• For a.e. t > 0, the following Kružkov inequalities hold

$$\operatorname{sgn}(\mathcal{U}(t,+\infty)-k)(f(\mathcal{U}(t,+\infty))-f(k)) \\ \geq \operatorname{sgn}(u(t,0^+)-k)(f(u(t,0^+))-f(k)))$$

for all $k \in \lfloor \mathcal{U}(t, +\infty), u(t, 0^+) \rceil$. In particular, we have:

$$f(\mathcal{U}(t,+\infty)) = f(u(t,0^+)) = v(t,0^-).$$

hence

$$\mathcal{U}(t,+\infty)=u(t,0^+).$$

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Theorem:

- 1. There exists a unique solution $(u^{\epsilon}, v^{\epsilon})$ to the original two-scale problem.
- There exists a unique solution (u, v) ∈ L[∞] ∩ BV(ℝ⁺ × ℝ) for the zero-ε limit system and satisfying L¹-contraction,

$$\frac{1}{2a} \| r_{+}^{1}(t,\cdot) - r_{+}^{2}(t,\cdot) \|_{L^{1}(\mathbb{R}^{-}_{x})} + \frac{1}{2a} \| r_{-}^{1}(t,\cdot) - r_{-}^{2}(t,\cdot) \|_{L^{1}(\mathbb{R}^{-}_{x})} \\ + \| u^{1}(t,\cdot) - u^{2}(t,\cdot) \|_{L^{1}(\mathbb{R}^{+}_{x})} \leq \| u_{0}^{1} - u_{0}^{2} \|_{L^{1}(\mathbb{R}_{x})},$$

for any two solutions (u^1, v^1) and (u^2, v^2) . Here $r^i_{\pm} = (au^i \pm v^i)(t, x)$, i = 1, 2.

3. For any T > 0, $(u^{\epsilon}, v^{\epsilon})$ converges strongly to (u, v) in $L^{1}((0, T) \times \mathbb{R})$.

Numerical example 1

Let $f(u^{\epsilon}) = \frac{1}{4}(e^{-u^{\epsilon}} - 1)$, with initial condition $u^{\epsilon}(x, 0) = \sin(\pi x)^3$. In this case, f'(u) < 0, so there will be an interface layer at the interface x = 0.





Figure: T = 0.5, $\Delta x = 0.04$, $\Delta t = 0.02$.



Figure: T = 0.5, $\Delta x = 0.01$, $\Delta t = 0.005$.

$$u^{\epsilon}(x,0) = \left\{ egin{array}{ll} 1, & ext{if } -1 \leq x \leq 0.2; \ -1, & ext{if } 0.2 < x \leq 1. \end{array}
ight.$$



Figure: A shock from the right region passing through the interface.

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$$u^{\epsilon}(x,0) = \left\{ egin{array}{cc} -1, & ext{if } -1 \leq x \leq 0.2; \ 1, & ext{if } 0.2 < x \leq 1. \end{array}
ight.$$



Figure: Rarefaction wave

Numerical example 2

Now we consider the case f'(u) > 0. Let $f(u^{\epsilon}) = \frac{1}{4}(e^{u^{\epsilon}} - 1)$, initial condition $u^{\epsilon}(x, 0) = \sin(\pi x)^3$.





Figure: T = 0.6, $\Delta x = 0.04$, $\Delta t = 0.02$.



Figure: T = 0.5, $\Delta x = 0.01$, $\Delta t = 0.005$.

$$u^{\epsilon}(x,0) = \left\{ egin{array}{cc} -1, & ext{if } -1 \leq x \leq -0.2; \\ 1, & ext{if } -0.2 < x \leq 1. \end{array}
ight.$$



Figure: A contact discontinuity passing through the interface.

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