

# Long time behavior of solutions to the 2D Keller-Segel equation

Yao Yao  
Georgia Tech

based on joint works with  
José Carrillo, Sabine Hittmeir, Bruno Volzone  
and Katy Craig, Inwon Kim

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# Keller-Segel equation with degenerate diffusion

- In  $\mathbb{R}^d$  with  $d \geq 2$ , the K-S equation with degenerate diffusion is

$$\rho_t = \Delta \rho^m + \nabla \cdot (\rho \nabla (\mathcal{N} * \rho)),$$

where  $\mathcal{N}$  is the Newtonian potential in  $\mathbb{R}^d$ . Here we assume  $m > 1$ , which models the anti-overcrowding effect. (Boi-Capasso-Morale '00, Topaz-Bertozzi-Lewis '06)

- The behavior of solutions depends on  $m$ , with  $m_c = 2 - \frac{2}{d}$  being the critical power:
  - For  $m > m_c$ , for any  $\rho_0 \in L^1 \cap L^\infty(\mathbb{R}^d)$ , the solution exists globally in time, and the  $L^\infty$  norm stays uniformly bounded in time. (Sugiyama '06)
  - For  $m < m_c$ , there might be a finite-time blow-up for initial data with arbitrarily small mass. (Sugiyama '06)
  - For  $m = m_c$ , the behavior of solution depends on its mass, and there is a critical mass  $M_c$ . (Blanchet-Carrillo-Laurençot '09)

- From now on, we focus on the “subcritical” case  $m > 2 - \frac{2}{d}$ , where the solutions are known to exist globally in time.
- Question: long time behavior of solutions?
- If  $\rho$  is a solution the Keller-Segel equation, then  $E_m[\rho]$  is non-increasing in time:

$$E_m[\rho] = \frac{1}{m-1} \int_{\mathbb{R}^d} \rho^m dx + \frac{1}{2} \int_{\mathbb{R}^d} \rho(\mathcal{N} * \rho) dx.$$

- Let  $\rho_A$  be the global minimizer of  $E_m$  among all densities with mass  $A$ . Then  $\rho_A$  must be a stationary solution.
- The following are known about the global minimizer  $\rho_A$ :
  - Existence (Lions '84)
  - Radial symmetry (by Riesz's rearrangement inequality)
  - Uniqueness + compact support (Lieb-Yau '87)

## Question

If  $\rho_0$  has mass  $A$ , is it always true that  $\rho(\cdot, t)$  converges to (a translation of)  $\rho_A$  as  $t \rightarrow \infty$ ?

- The answer is yes ONLY IF we have a positive answer to the following question:

## Question

Is  $\rho_A$  the unique stationary solution with mass  $A$  (up to a translation)?

- For Newtonian potential, it is known that radially symmetric stationary solution (with a fixed mass) is unique (Lieb-Yau '87), so the above question is equivalent with

## Question

Is every stationary solution radially symmetric (up to a translation)?

# Stationary solutions for aggregation equation with nonlinear diffusion

Consider the equation with a general attracting kernel  $\mathcal{K}$ :

$$\rho_t = \Delta \rho^m + \nabla \cdot (\rho \nabla (\mathcal{K} * \rho)),$$

where  $\mathcal{K}$  is radial and is strictly increasing in  $|x|$ . Thus any stationary solution  $\rho_s$  satisfies

$$\frac{m}{m-1} \rho_s^{m-1} + \mathcal{K} * \rho_s = C_i$$

in each connected component of  $\{\rho_s > 0\}$ . ( $C_i$  can differ in different components).

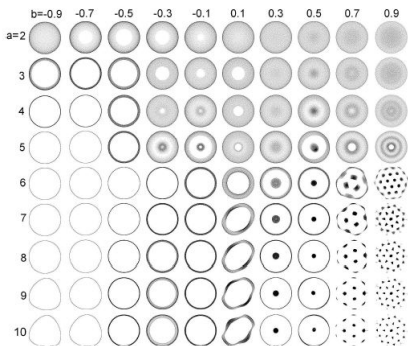
**Theorem (Carrillo-Hittmeir-Volzone-Y., '16)**

*Let  $\rho_s \in L^1_+(\mathbb{R}^d) \cap L^\infty(\mathbb{R}^d)$  be a stationary solution in the above sense. Then  $\rho_s$  must be radially decreasing up to a translation.*

# Contrast with the attractive-repulsive kernel

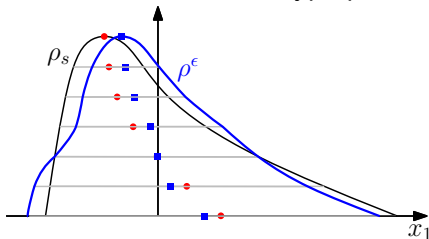
If  $K$  is repulsive in short-range and attracting in long-range, then stationary solutions to  $\rho_t = \nabla \cdot (\rho \nabla (\mathcal{K} * \rho))$  can have many non-radial patterns.

For example, when  $\mathcal{K}'(r) = \tanh((1 - r)a) + b$  with parameters  $a, b$ , below are the patterns of stationary solutions for some  $a, b$ :  
(Kolokolnikov-Sun-Uminsky-Bertozzi, '11)



# Sketch of the proof

- If a stationary solution  $\rho_s$  that is NOT radially decreasing after any translation, we perturb it using its **continuous Steiner symmetrization** about some hyperplane  $H$ :



- Since  $\int \rho_s^m = \int (\rho^\epsilon)^m$ , and interaction energy decreases in the first order for a short time (need some work to check this!),

$$E_m[\rho^\epsilon] - E_m[\rho_s] < -c\epsilon \quad \text{for all sufficiently small } \epsilon > 0,$$

where  $c > 0$  depending on  $\rho_s$  and  $\mathcal{K}$ .

- But no contradiction yet – the perturbation is only Hölder continuous near the zero levelset.

# Working towards a contradiction

- Let us now slow down the “velocity” at low density  $h < h_0$  as  $v(h) = (h/h_0)^{m-1}$ , and call the perturbation  $\mu^\epsilon$ . We still have

$$E_m[\mu^\epsilon] - E_m[\rho_s] < -c\epsilon \quad \text{for all sufficiently small } \epsilon > 0,$$

and using a priori regularity estimates on  $\rho_s$  gives

$$|\mu^\epsilon(x) - \rho_s(x)| \leq C\epsilon|\rho_s(x)| \quad \text{for all sufficiently small } \epsilon > 0.$$

- Combining the above pointwise estimate with the assumption that  $\rho_s$  is stationary, we have  $|E_m[\mu^\epsilon] - E_m[\rho_s]| < C\epsilon^2$ , contradicting the first inequality if  $\epsilon > 0$  is sufficiently small. So there cannot be such a  $\rho_s$ !



# Convergence of dynamical solution in 2D

- For any  $t_n \rightarrow \infty$ , weak lower semicontinuity of the entropy dissipation (Bian-Liu '13) gives that  $\|\rho(\cdot, t_{n_k}) - \rho_\infty\|_{L^1} \rightarrow 0$  for some stationary solution  $\rho_\infty$  along a subsequence  $t_{n_k} \rightarrow \infty$ .
- $\rho_\infty$  has the same center of mass as  $\rho_0$ : center of mass is preserved during evolution.
- $\rho_\infty$  also has the same mass as  $\rho_0$ : the second moment  $\int |x|^2 \rho(x, t) dx$  is uniformly bounded in time. (This argument works in 2D only!)

## Theorem (Carrillo-Hittmeir-Volzone-Y., '16)

For any  $\rho_0 \in L^\infty(\mathbb{R}^2) \cap L^1((1 + |x|^2)dx)$ , we have

$$\lim_{t \rightarrow \infty} \|\rho(\cdot, t) - \rho_s\|_{L^q} = 0 \text{ for any } 1 \leq q < \infty,$$

where  $\rho_s$  is the (unique) stationary solution with the same mass and same center of mass as  $\rho_0$ .

- But we are unable to obtain any convergence rates.

# The “ $m = \infty$ ” limit (joint with K.Craig and I.Kim)

If we take the “ $m \rightarrow \infty$ ” limit in  $\rho_t = \Delta \rho^m + \nabla \cdot (\rho \nabla (\rho * \mathcal{N}))$ ,  $\rho(\cdot, t)$  should intuitively evolve like congested penguins:



# Gradient flow with density constraint

- Consider the gradient flow (in the 2-Wasserstein metric space) of the functional

$$E_\infty[\rho] := \begin{cases} \frac{1}{2} \int_{\mathbb{R}^d} \rho(\mathcal{N} * \rho) dx & \text{if } \|\rho\|_\infty \leq 1, \\ +\infty & \text{otherwise.} \end{cases}$$

- Such problem has been studied when  $\mathcal{N} * \rho$  is replaced by a fixed potential  $\Phi$ . (Maury-Roudneff-Chupin-Santambrogio '10, Alexander-Kim-Yao '14)

## Question

Is the gradient flow of  $E_\infty$  well defined, and is it unique?

- Thanks to the  $L^\infty$  constraint, the energy  $E_\infty$  has certain convexity properties along generalized geodesics (Carrillo-Lisini-Mainini '14), hence the continuous gradient flow to  $E_\infty$  is well defined (Craig '15).

## Question

If  $\rho_0 = 1_{\Omega_0}$  (i.e. it is of “patch type”), then what PDE determines its boundary evolution?

- A heuristic argument suggests that the actual velocity field at time  $t$  should be  $-\nabla(\rho * \mathcal{N}) - \nabla u$ , with  $u$  given by

$$\begin{cases} \Delta u = -1 \text{ in } \Omega(t), \\ u = 0 \text{ on } \partial\Omega(t). \end{cases}$$

- So the free boundary  $\partial\Omega(t)$  should evolve like a Hele-Shaw type equation, with free boundary velocity given by above.

## Theorem (Craig-Kim-Y., '16)

*Let  $\rho_\infty(\cdot, t)$  be the gradient flow of  $E_\infty$  with  $\rho_0 = 1_{\Omega_0}$ . Then we have  $\rho_\infty(t) = 1_{\Omega(t)}$  a.e. for all time, where  $\Omega(t)$  is a viscosity solution of the above free boundary problem.*

# Convergence towards a disk

## Question

Long-time behavior of solutions?

## Theorem (Craig-Kim-Y., '16)

For  $d = 2$  and  $\rho_0 = 1_{\Omega_0}$ , we have  $\lim_{t \rightarrow \infty} \|\rho(t) - 1_B\|_{L^q} = 0$  for any  $1 \leq q < \infty$ , where  $B$  is a disk with area 1 whose center coincides with the center of mass of  $\rho_0$ . Moreover, we have

$$0 \leq E_\infty[\rho(t)] - E_\infty[1_B] \leq C(M_2[\rho_0])t^{-1/6} \text{ for all } t \geq 0.$$

- Reason: the evolution of second moment (in 2D) is given by

$$\frac{d}{dt} M_2[\rho(t)] = -\frac{1}{2\pi} + 4 \int_{\Omega(t)} u(x) dx,$$

where  $\Delta u = -1$  in  $\Omega(t)$  and  $u = 0$  on  $\partial\Omega(t)$ .

- Using similar ideas as [Talenti '76](#), we have  $\text{RHS} \leq 0$ , where the equality is achieved if and only if  $\Omega$  is a disk.

# Idea of convergence proof

- Using a stability result of isoperimetric inequality (Fusco-Maggi-Pratelli '08), we have

$$-\frac{1}{4\pi} + 2 \int_{\Omega(t)} u(x) dx \lesssim -A(\Omega(t))^3,$$

where

$$A(E) := \inf \left\{ \frac{|E \Delta (x_0 + B)|}{|E|} : x_0 \in \mathbb{R}^2, B \text{ is a disk with } |B| = |E| \right\}.$$

- This result enables us to get some (non-optimal) convergence rate for  $d = 2$ .
- For  $d \geq 3$ , we are unable to obtain any compactness result of  $\rho(t)$  as  $t \rightarrow \infty$ , therefore it is unknown whether  $\Omega(t)$  converges to a ball as  $t \rightarrow \infty$ .



Thank you for your attention!