# Two Classical models of Quantum Dynamics 

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## Outline:

- Review of dynamics and localization for Random Schrödinger $H$ on $\ell_{2}\left(\mathbb{Z}^{d}\right)$

$$
H=-\Delta+\lambda V(j) ; \quad j \in \mathbb{Z}^{d}, \quad V(j) i i d
$$

- Motion of Classical particle on Manhattan lattice deflected by random obstructions. Quantum Network Model
- Edge Reinforced Random Walk - P. Diaconis SUSY Statistical mechanics M. Zirnbauer
- Classical dynamics: Kesten, Papanicolaou, Komorowski and Ryzhik

$$
H=\dot{x}^{2}+\lambda V(x), \quad x \in \mathbb{R}^{d}
$$

Localization and Dynamics for Random Schrödinger

$$
\begin{gathered}
H=-\Delta+\lambda V(j) ; \quad j \in \mathbb{Z}^{d}, \quad V(j) i i d . \\
\mathbb{E}(V(j))=0 ; \quad \mathbb{E}\left(V(j)^{2}\right)=1 \\
-\Delta e^{i j \cdot p}=\sum_{\alpha=1}^{d}\left(2-2 \cos p_{\alpha}\right) e^{i j \cdot p} \approx p^{2} e^{i j \cdot p} \\
i \frac{\partial}{\partial t} \psi=H \psi ; \quad \psi=\psi(t, j) ; \quad \psi(t=0, j)=\delta_{0}(j) \\
R^{2}(t) \equiv \mathbb{E} \sum_{j}|\psi(t, j)|^{2}|j|^{2} ; \quad \sum_{j}|\psi(t, j)|^{2}=1 ;
\end{gathered}
$$

$R^{2}(t) \approx C t^{2} ; \quad$ Ballistic ; if $\lambda=0$ or $V$ periodic

## Localization

$$
R^{2}(t) \leq \text { Const; } \quad \text { Uniform Localization }
$$

All eigenfunctions of $H$ decay rapidly about some point $c(\alpha) \in \mathbb{Z}^{d}$ :

$$
\left|\psi_{\alpha}(j)\right| \leq C e^{-|j-c(\alpha)| / \ell(E, \lambda)} ; \quad \ell(E, \lambda)=\text { localization length }
$$

Theorem In one dimension, $\mathbb{Z}^{1}$, all eigenstates are uniformly localized $\ell(E, \lambda) \approx \lambda^{-2}$, all $\lambda>0$

Furstenberg, Goldsheid, Molchanov, Pastur, Ledrappier, Margulis Conjecture A : In $\mathbb{Z}^{2}$, Localization for all $\lambda>0 ; \quad \ell(E, \lambda) \leq e^{\lambda^{-2}}$
Conjecture $A^{\prime}:$ In a 1 D strip of width $\mathrm{W}, \quad \ell(E, \lambda) \leq W \lambda^{-2}$

## Localization on $\mathbb{Z}^{d}$ for $d \geq 2$

If $\lambda \gg 1$ then

1) All eigenstates of H are exponentially localized.
2) $R(t)^{2} \leq$ Const for all d . Note this is false in the continuum!

If $|\lambda| \ll 1$, there is localization for $E \leq-\lambda^{2} c_{d}$
In 3D expect extended states for E above $-\lambda^{2} c_{d}^{\prime}$ - Anderson transition.

Anderson, Thouless, Fröhlich, Sp, Martinelli, Scoppola, Simon, Wolff, Aizenman, Molchanov ...

## Localization via Green's function

$$
G_{E+i \epsilon}(x, y) \equiv[H-E-i \epsilon]^{-1}(x, y), \quad x, y \in \mathbb{Z}^{d}, \quad \epsilon>0
$$

Fractional moment method:

$$
\mathbb{E}\left|G_{E+i 0}(x, y)\right|^{1 / 2} \leq e^{-|x-y| / \ell(E)} \Rightarrow \text { Localization near } \mathrm{E}
$$

Aizenman, Molchanov, Hundertmark, Friedrich, Schenker.

Conjecture - 3D quantum diffusion:

$$
\sum_{x} e^{i x \cdot p} \mathbb{E}\left|G_{E+i \epsilon}(x, 0)\right|^{2} \approx C\left[D(E, \epsilon) p^{2}+\epsilon\right]^{-1}
$$

Diffusion constant $=D(E, \epsilon)=D(E) ; \epsilon \downarrow 0$, time scale $=\epsilon^{-1}$

$$
\mathbb{E}\left|G_{E+i \epsilon}(x, x)\right|^{2} \leq \text { Const } ; \Rightarrow \text { Absolutely Continuous spectrum }
$$

## Quantum Dynamics in 3D

$$
H=-\Delta+\lambda V(j) ; \quad j \in \mathbb{Z}^{d}, \quad V(j) \text { iid }
$$

Conjecture - 3D quantum diffusion:

$$
R^{2}(t) \approx D t, \quad \text { for small } \lambda
$$

The wave spread ballistically until time scale $\lambda^{-2}$. After this time we expect quantum diffusion.

Erdős-Salmhofer-Yau get control for times $\leq \lambda^{-2-\delta}, \delta$ small.

## Manhattan Pinball

## Quantum Network Model with random scatterers

Motivation: Chalker's network model Integer Quantum Hall.
Particle moves on $\mathbb{Z}^{2}$ with alternating orientations for the streets.

Independently with Prob $0<p<1$, an obstruction is randomly placed at a vertex. Particle moves in the direction of orientation until it meets an obstruction. Then it turns in the direction of orientation.
(Beamond, Cardy, Owczarek; Gruzberg, Ludwig, Read)


Figure: Manhattan Lattice

Theorem If $p>1 / 2$ then all orbits are closed with probability 1 Localization. Proof by percolation.

Conjecture: All loops are closed for any $p>0$
Average loop length $\ell \approx e^{c p^{-2}} \gg 1$.
At intermediate distances paths should look like random walk: diffusion at scales $p^{-m}, m \geq 2$.

Problems: Prove localization for some $p<1 / 2$ Show that on a 1D cylinder of width $W, \quad \ell(W) \leq$ Const $p^{-2} W$. Note easy to prove: $\ell(W) \leq e^{C W}$

## Linearly Edge Reinforced Random Walk

History dependent walk $W_{n} \in \mathbb{Z}^{d}, n \in \mathbb{Z}^{+}$:
Walk takes nearest neighbor steps and favors edges $j, k, \in \mathbb{Z}^{d},|j-k|=1$, it has visited in the past.

Introduced by P. Diaconis in 1986 while wandering the streets of Paris. He liked to return to streets he had visited in the past.

## Why Linearly Edge Reinforced?

Related to Polya's Urn.
Partially exchangeable process - generalization of de Finetti
Not Markovian but is a superposition of Markov Processes
Equivalent to a random walk in a random environment - must average over environment: $\mathcal{E}$

## Definition of Reinforced Random Walk

Let $C_{j k}(n)=$ number of times the walk has crossed edge $j k$ up to time n and let $\beta>0$.

$$
\operatorname{Prob}\left\{W_{n+1}=k \mid W_{n}=j\right\}=\frac{1+C_{j k}(n) / \beta}{\mathcal{N}_{\beta}}, \quad|j-k|=1
$$

$\mathcal{N}$ is the normalization: $\mathcal{N}_{\beta}=\sum_{k^{\prime}}\left(1+C_{j k^{\prime}}(n) / \beta\right), \quad\left|j-k^{\prime}\right|=1$
$0<\beta \ll 1$, strong reinforcement, (high temperature) $\beta \gg 1$, weak reinforcement, (low temperature)

## Long time behavior of $W_{n}$, $n$, large?

Is ERRW recurrent? Localized?

## Localization:

$\operatorname{Prob}_{\beta}\left\{\left|W_{n}-W_{0}\right| \geq R\right\} \leq C e^{-R / \ell}, \quad \ell(\beta)=$ localization length

Is ERRW Transient? Diffusive?
Is there a Phase Transition as we vary the reinforcement $\beta$ ?

## P. Diaconis and D. Coppersmith (1986):

ERRW $\approx$ random walk in a random environment.
Environment: The rate at which an edge $j, j^{\prime}$ is crossed $w_{j, j^{\prime}}>0$ are correlated random variables. (conductances)

Explicit Joint Distribution of $w_{j, j^{\prime}}>0$ first appeared in an unpublished paper Diaconis and Coppersmith. It given by statistical mechanics.

The generator for the RW is a weighted Laplacian $L$

$$
v^{t} \cdot L v=\sum_{\left|j-j^{\prime}\right|=1} w_{j, j^{\prime}}\left(v_{j}-v_{j^{\prime}}\right)^{2}
$$

However, L is NOT uniformly elliptic.
Important: The distribution of $w_{j, j^{\prime}}$ depends on the starting point of the Walk.

## Relation to Random Schrödinger

Spectral properties of Random Schrödinger Equivalent dual model in statistical mechanics with Supersymmetric Hyperbolic symmetry - 1982 Efetov: U(1, 1|2) SUSY - Rigorous equivalence but very complicated.

In 1991 Martin Zirnbauer defined a simplified version of Efetov's dual model: SUSY Hyperbolic sigma model $H^{2 \mid 2}$

In any dimension, spin correlations of $H^{2 \mid 2}$ can be expressed as a random walk in a random environment.

$$
w_{j, j^{\prime}}=e^{t_{j}+t_{j^{\prime}}}, \quad \text { joint distribution } \equiv e^{-E_{S U S Y}\left(\beta,\left\{t_{j}\right\}\right)}
$$

The expectation is denoted by $\mathcal{E}$.

## The SUSY Hyperbolic Model - $H^{(2 \mid 2)}$

$$
\begin{gathered}
E_{\operatorname{SUSY} Y}\left(\left\{t_{j}\right\}\right)= \\
\beta \sum_{j \sim j^{\prime}} \cosh \left(t_{j}-t_{j^{\prime}}\right)-1 / 2 \log \operatorname{det} L_{\beta, \epsilon}(t)+\epsilon \sum \cosh t_{j}
\end{gathered}
$$

where $L$ is weighted Laplacian:

$$
\left[v ; L_{\beta, \epsilon}(t) v\right]=\beta \sum_{\left(j^{\prime} \sim j\right)} \mathrm{e}^{t_{j}+t_{j^{\prime}}}\left(v_{j}-v_{j^{\prime}}\right)^{2}+\epsilon \sum_{k \in \Lambda} \mathrm{e}^{t_{k}} v_{k}^{2}
$$

Spin-Spin correlation:

$$
<e^{t_{0}+t_{x}} L_{\beta, \epsilon}(t)^{-1}(0, x)>_{\operatorname{SUSY}}(\beta, \epsilon)
$$

## Some results in 1 and 2D

R. Pemantle analyzed ERRW on the Regular tree. Showed that it has sharp transition in $\beta$ from recurrent to transient.

Merkl and Rolles studied one dimensional strips of Width W and show the ERRW is localized with $\ell(W) \leq \beta W$
In 2D Merkl and Rolles prove the conductance $\mathcal{E} w_{j j^{\prime}}^{1 / 4}$ has a power law decay away from the origin - via a deformation argument.

Sabot and Zeng: In 2D ERRW is recurrent for all $\beta$.
Conjecture: In 2D the walk is exponentially localized for all $\beta$.

Theorem (Disertori-S-Zirnbauer '10) Phase transition:
For $\beta \gg 1$ and $d \geq 3$ the conductances $w_{j j^{\prime}}$ in the $H^{2 \mid 2}$ model are bounded above and below with high probability - Transient, quasi-diffusion

For $0<\beta \ll 1$, Conductance goes 0 , Recurrent, Localization
Kozma and later Sznitman pointed out similarities of $\mathrm{H}^{2 \mid 2}$ to the formulas for ERRW.

Sabot and Tarres ('12) Showed how to modify $\mathrm{H}^{2 \mid 2}$ model to get the law for ERRW.

## Phase Transition for ERRW in 3D

Theorem (Sabot-Tarres, Angel-Crawford-Kozma) For strong reinforcement, $0<\beta \ll 1$, ERRW is recurrent and we have Localization:

$$
\operatorname{Prob}\{|W(t)-W(0)| \geq R\} \leq C e^{-R / \ell(\beta)}, \quad \mathcal{E} W^{2}(t) \leq \text { Const }
$$

and $\ell(\beta)$ is the localization length.
Theorem (Disertori-Sabot-Tarres) For weak reinforcement, $\beta \gg 1$, and $d \geq 3, W(t)$ is Transient, quasi-diffusion.

## Outlook and Problems

A) For Random Schrödinger show that $R^{2}(t) \leq t^{2-\delta}$ strictly sub-ballistic.
B) For Manhattan model show that for long time scales $t \leq p^{-m}$ typical trajectories behave like a random walk for p is small. (For ESY, $m=1.1$ )
C) Is there a sharp transition in 3D?: Localization for $0<\beta<\beta_{c}$ and Diffusion for $\beta>\beta_{c}$
D)Multi-fractal transition is expected in 3D at $\beta_{c}$ for both Random Schrödinger and ERRW. Is MF present on Bethe lattice for ERRW?
E) In 2D, is ERRW is localized for weak reinforcement? For all $\beta>0, \quad \ell(\beta) \approx e^{C \beta}, \quad \mathcal{E} w_{j, j^{\prime}}^{1 / 4} \approx e^{-|j| / \ell(\beta)}$.
Note analogy to Aizenman-Molchanov fractional moment

