

TIME-ASYMPTOTIC CONVERGENCE RATES TOWARDS THE DISCRETE EVOLUTIONARY STABLE DISTRIBUTION

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ABSTRACT. This paper is concerned with the discrete dynamics of an integro–differential model that describes the evolution of a population structured with respect to a continuous trait. Various time–asymptotic convergence rates towards the discrete evolutionary stable distribution (ESD) are established. For some special ESD satisfying a strict sign condition, the exponential convergence rates are obtained for both semi-discrete and fully discrete schemes. Towards the generic ESD, the algebraic convergence rate we find is consistent with the known result for the continuous model.

1. INTRODUCTION

This paper is a continuation of the work [12] investigating an entropy satisfying finite volume method for a direct competitive selection model. The model is given by

$$(1.1a) \quad \partial_t f(t, x) = \left(a(x) - \int_X b(x, y) f(t, y) dy \right) f(t, x), \text{ for } t > 0, x \in X,$$

$$(1.1b) \quad f(0, x) = f_0(x), \quad x \in X.$$

This is an integro–differential equation that describes the evolution of a population of density $f(t, x)$ structured with respect to a continuous trait x . The space of traits X can be fairly general, even though for simplicity we will take a subset of \mathbb{R}^d . In this model, a is the reproduction rate for an individual alone (without competition with other individuals); and $b > 0$ corresponds to the interaction between individuals which we assume here to be only competitive. The total reproduction rate of each individual is thus determined by its trait and the environment through the selective pressure $a(x) - \int_X b(x, y) f(t, y) dy$, leading therefore to selection.

Existence and stability of regular or measure valued solutions for (1.1) are known, provided that the coefficients have enough regularity (see [7, 11, 19]). Together with variants, it has been investigated much in the literature; see e.g., [1, 3, 9, 10, 17]. In addition, equation (1.1) (with an additive mutation term) can be derived from stochastic models of finite population (see [4, 5, 8]), and there is a vast literature (see e.g., [2, 6, 13, 14, 15, 16, 18]) on the study of the combining effects of both selection and small mutation on the population dynamics.

The model without mutation is interesting from the point of view of large time behavior. Natural questions appear, such as does the population really converge to an equilibrium? Is this equilibrium an evolutionarily stable strategy or distribution (ESS or

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ESD)? Does this limit depend on the initial population distribution? A definite answer to these questions has been provided in [11] under additional assumptions on b ,

$$(1.2) \quad \forall g \in L^1(X) \setminus \{0\}, \quad \int \int b(x, y) g(x) g(y) dx dy > 0,$$

and b is assumed to satisfy some symmetry, for instance,

$$(1.3) \quad \forall x, y \in X, \quad b(x, y) = b(y, x),$$

so that solutions of (1.1) converge to the then unique ESD at rate $O(\log t/t)$ for some proper initial data.

The finite volume scheme investigated in [12] is shown to produce numerical solutions with satisfying long-time selection dynamics. This is achieved through a proper discretization, so that the numerical solution

$$f_\alpha^n \sim \frac{1}{h^d} \int_{I_\alpha} f(n\Delta t, x) dx,$$

approximates $f(t_n, x)$ over the cell I_α indexed by $\alpha \in \Lambda \subset \mathbb{Z}^d$, and the discrete relative entropy

$$F^n = \sum_\alpha \left(\tilde{f}_\alpha \log \left(\frac{\tilde{f}_\alpha}{f_\alpha^n} \right) + f_\alpha^n - \tilde{f}_\alpha \right) h^d,$$

satisfies the entropy dissipation inequality

$$F^{n+1} - F^n \leq -\frac{1}{2} \Delta t \|f^n - \tilde{f}\|_b^2.$$

This implies the convergence of f^n toward the discrete ESD. One may wonder what would be the time-asymptotic convergence rate, such an inquiry motivates therefore the present work.

1.1. Assumptions and main results. For simplicity, we restrict ourselves to only one dimensional setting for $X = [-1, 1]$, though the results and the proofs can easily be generalized for any dimension. We partition X into sub cells $I_j = [x_{j-1/2}, x_{j+1/2}]$, $j = 1 \cdots N$, for a uniform mesh $h = 2/N$ so that $x_{j-1/2} = x_{1/2} + (j-1)h$ with $x_{1/2} = -1$, $x_{N+1/2} = 1$. We consider the following semi-discrete scheme

$$(1.4) \quad \frac{d}{dt} f_j = f_j \left(\bar{a}_j - h \sum_{i=1}^N \bar{b}_{ji} f_i \right), \quad j = 1, \dots, N,$$

where

$$(1.5) \quad \bar{a}_j = \frac{1}{h} \int_{I_j} a(x) dx, \quad \bar{b}_{ji} = \frac{1}{h^2} \int_{I_i} \int_{I_j} b(x, y) dx dy,$$

and the numerical solution $f_j(t)$ approximates the cell average of the exact solution f ,

$$\bar{f}_j(t) = \frac{1}{h} \int_{I_j} f(t, x) dx.$$

Set

$$s_j[f] = \bar{a}_j - h \sum_{i=1}^N \bar{b}_{ji} f_i,$$

then the nonlinear dynamical system (1.4) admits many steady states satisfying

$$\tilde{f}_j s_j[\tilde{f}] = 0, \quad j = 1, \dots, N.$$

Of special interest is the discrete ESD $\tilde{f} = \{\tilde{f}_j\}$, which is defined as

$$(1.6a) \quad \forall j \in \{1 \leq i \leq N, \tilde{f}_i \neq 0\}, \quad s_j[\tilde{f}] = 0,$$

$$(1.6b) \quad \forall j \in \{1 \leq i \leq N, \tilde{f}_i = 0\}, \quad s_j[\tilde{f}] \leq 0,$$

and is conjectured to be the limit of the solution provided the initial data $f_j(0) > 0$ for all $j = 1, 2, \dots, N$.

From the basic assumptions at the continuous level one may derive similar assumptions at the discrete level:

$$(1.7a) \quad |\bar{a}_j| \leq \|a\|_{L^\infty}, \quad \{1 \leq j \leq N, \bar{a}_j > 0\} \neq \emptyset;$$

$$(1.7b) \quad 0 < b_f \leq \bar{b}_{ji} \leq \|b\|_{L^\infty} \text{ and } \bar{b}_{ji} = \bar{b}_{ij}, \text{ for } 1 \leq i, j \leq N;$$

$$(1.7c) \quad \sum_{j=1}^N \sum_{i=1}^N \bar{b}_{ji} g_i g_j > 0 \text{ for any } g_j \text{ such that } \sum_{j=1}^N |g_j|^2 \neq 0.$$

It is shown that these assumptions ensure the existence and uniqueness of the ESD, as defined by (1.6), in [12].

Given the positivity assumption, \bar{b}_{ij} induces a discrete weighted norm denoted by $\|\cdot\|_b$

$$(1.8) \quad \|g\|_b = \left(h^2 \sum_{i,j=1}^N \bar{b}_{ji} g_j g_i \right)^{\frac{1}{2}}.$$

We also use the discrete l^p norm

$$\|g\|_p = \left(\sum_{j=1}^N |g_j|^p h \right)^{1/2}.$$

Those norms are related through

$$(1.9) \quad \sqrt{h\lambda_{\min}} \|g\|_2 \leq \|g\|_b \leq \sqrt{h\lambda_{\max}} \|g\|_2,$$

where $\lambda_{\min}(\lambda_{\max})$ denotes the smallest (largest) eigenvalue of $B = (\bar{b}_{ji})_{N \times N}$ and $\|B\|_2 = \lambda_{\max}$.

Lemma 1.1. *One has*

$$2b_f \leq h\lambda_{\max} \leq 2\|b\|_{L^\infty}, \quad \lambda_{\max} \geq \lambda_{\min} > 0.$$

Proof. By the positivity of the matrix B , the upper bound can be obtained through the trace of B ,

$$\lambda_{\max} \leq \text{Tr}(B) = \sum_{i=1}^N \bar{b}_{ii} \leq N\|b\|_{L^\infty} = \frac{2}{h}\|b\|_{L^\infty}.$$

As for the lower bound, we use

$$\lambda_{\max} = h^{-1} \sup_{\{\|g\|_2=1\}} \|g\|_b^2 \geq \frac{h}{2} \sum_{i,j} \bar{b}_{ij} \geq \frac{h}{2} N^2 b_f = 2h^{-1} b_f,$$

by choosing $g = (1/\sqrt{2}, \dots, 1/\sqrt{2})$.

Finally as B is strictly positive then $\lambda_{\min} > 0$. \square

Remark 1.1. *The size of $h\lambda_{\max}$ is bounded from above and below, but $h\lambda_{\min}$ can be much smaller as the mesh size vanishes.*

We call the ESD a strict ESD if it also satisfies the following strict sign condition,

$$(1.10) \quad s_j[\tilde{f}] < 0 \quad \text{for } j \in \{i : \tilde{f}_i = 0\}.$$

We shall prove that the strict ESD is both linearly and nonlinearly stable, with perturbations decaying to zero exponentially in time. To precisely state the main results, we use the following notation,

$$(1.11) \quad I = \{j \mid \tilde{f}_j = 0 \text{ and } s_j < 0\}, \quad I^c = \{j, 1 \leq j \leq N\} - I,$$

and

$$s = \min_{j \in I} (-s_j[\tilde{f}]) > 0, \quad f_m = \min_{j \in I^c} \tilde{f}_j > 0.$$

In the sequel we also use

$$\mu = hf_m\lambda_{\min}, \quad r = \min\{s, \mu\}$$

to quantify the exponential decay of the perturbations.

The result for the semi-discrete scheme is summarized in the following.

Theorem 1.1. *Assume (1.7) holds. Let $f_j(t)$ be the solution to the semi-discrete scheme (1.4), associated with the strict ESD, then there exists $\delta^* > 0$ such that for any $\delta \in (0, \delta^*)$ if*

$$\|f(0) - \tilde{f}\|_2 \leq \delta,$$

then

$$\|f(t) - \tilde{f}\|_p \leq C(1+t)^\xi e^{-rt}, \quad \xi = 1_{\{s=\mu\}},$$

where $1 \leq p \leq 2$,

$$\delta^* = \frac{\alpha^2 \min\{1, \sqrt{f_m}\}}{\sqrt{2} \max\{1, \alpha\}}, \quad \alpha = \sqrt{\frac{r}{\|b\|_{L^\infty}} + \frac{\|\tilde{f}\|_1}{2}} - \sqrt{\frac{\|\tilde{f}\|_1}{2}},$$

and C may depend on the parameters and the norms of the initial data but not explicitly on N or h .

Remark 1.2. *While none of the constants in the previous result depend explicitly on the mesh size, most of them depend on it implicitly. For instance s is not in general bounded from below uniformly in h and in most cases one can actually prove that $s \rightarrow 0$ as $h \rightarrow 0$. This is because at the limit \tilde{f} should be an ESD for the continuous model, therefore $s[\tilde{f}]$ is a smooth function of x . The extension of I is now the set of x where the measure \tilde{f} vanishes, that is the complement of the support of \tilde{f} . But the function $s[\tilde{f}]$ vanishes on the support of \tilde{f} , which is a closed set, and therefore cannot be bounded from below on the complement. The same argument applies to f_m .*

As a consequence the exponential convergence is not uniform in h and actually degenerates as $h \rightarrow 0$. The same is true for all the exponential convergence results presented here. Only the algebraic rate, Th. 1.3, is uniform in h .

For the fully discrete scheme

$$(1.12) \quad \frac{f_j^{n+1} - f_j^n}{\Delta t} = f_j^{n+1} \left(\bar{a}_j - h \sum_{i=1}^N \bar{b}_{ji} f_i^n \right),$$

the exponential convergence rate toward the strict ESD can still be obtained under some restriction on the time step.

Theorem 1.2. *Assume (1.7) holds. Let f_j^n be the numerical solution to (1.12), associated with the strict ESD, $\tilde{f} = \{\tilde{f}_j\}$. If Δt satisfies*

$$\Delta t \leq \frac{\mu}{2\|b\|_{L^\infty}^2 \|\tilde{f}\|_1^2},$$

then there exists $\delta^* > 0$ such that for any $\delta \in (0, \delta^*)$ if

$$\|f^0 - \tilde{f}\|_2 \leq \delta,$$

then for $1 \leq p \leq 2$,

$$\|f^n - \tilde{f}\|_p \leq C(1 + n\Delta t)^\xi \max\{\sqrt{K_s}, K_*\}^n, \quad \xi = 1_{\{\sqrt{K_s} = K_*\}},$$

where

$$K_s = \frac{1}{\sqrt{1 + 2s\Delta t}}, \quad K_* = \frac{1}{\sqrt{1 + 2\mu\Delta t}} \left(1 + \frac{2\Delta t^2 \|b\|_{L^\infty}^2 \|\tilde{f}\|_1^2}{\sqrt{1 + 2\mu\Delta t}} \right) < 1,$$

and C may depend on the parameters but not explicitly on N or h .

Remark 1.3. *In the above two theorems, r or μ may well approach zero when N tends to ∞ . Hence, exponential convergence for the continuous model cannot be deduced from these results and is in fact not expected.*

One objective of this work is to establish an algebraic convergence rate but with parameters uniform in N or h thus extending the rates known at the limit.

Theorem 1.3. *Assume (1.7) holds. Let f_j^n be the numerical solution generated from scheme (1.12) with positive initial data $f_j^0 > 0$ for all $j = 1, \dots, N$, with $\tilde{f} = \{\tilde{f}_j\}$ as its associated ESD. If*

$$F^0 := \sum_{j=1}^N \left(\tilde{f}_j \log \left(\frac{\tilde{f}_j}{f_j^0} \right) + f_j^0 - \tilde{f}_j \right) h < +\infty,$$

then

$$\|f^n - \tilde{f}\|_b^2 \leq \frac{2F^0}{n\Delta t},$$

provided that

$$\Delta t \leq \min \left\{ \frac{\lambda_{\min}}{\lambda_{\max}[2(\|a\|_{L^\infty} + \|b\|_{L^\infty} \|\tilde{f}\|_1) + 2\lambda_{\max}S(F^0) + \lambda_{\min}S(F^0)]}, \frac{h}{\|b\|_{L^\infty} S(F^0)} \right\},$$

where S is an explicit non-decreasing, positive function, which we specify in the proof.

Remark 1.4. *When $0 < b_f \leq \bar{b}_{ji}$ is satisfied, the same convergence rate can be obtained under a different time step restriction; see Theorem 3.2.*

Several techniques are introduced and developed in the proofs of these results.

In the proofs of Theorem 1.1 and 1.2 on the exponential convergence, we start with a symmetrization of the system with weight depending on the strict ESD, and then obtain exponential decay of the perturbations using a Lyapunov functional approach, subject to a parameter tuned to allow for the largest possible initial perturbations. Finally the optimal convergence rate is obtained by a refined estimate. In the proof of Theorem 1.3 on the algebraic convergence, we first establish the dissipation inequality of relative entropy, and further show the decreasing property of the dissipation rate, these together ensure the algebraic convergence rate towards the generic ESD.

The rest of this paper is organized as follows. In section 2, we present linear and nonlinear asymptotic stability of the strict ESD for the semi-discrete scheme. Section 3 is devoted to the fully discrete scheme, including the exponential convergence towards the strict ESD, and the algebraic convergence towards the general ESD.

2. EXPONENTIAL CONVERGENCE TOWARDS THE ESD FOR THE SEMI-DISCRETE SCHEME

In this section, we show that the strict ESD that satisfies the sign condition (1.10) is both linearly and nonlinearly stable, with perturbations decaying exponentially in time.

2.1. Linear stability. We first investigate the linear stability of the strict ESD satisfying (1.10). To do so, we consider the linearized equation

$$(2.1) \quad \frac{d}{dt}g_j = s_j g_j - \tilde{f}_j h \sum_{i=1}^N \bar{b}_{ji} g_i, \quad j = 1, \dots, N.$$

For the strict ESD, we define the weighted l^2 -norm by

$$\|g\|_{\tilde{f}} = \left(\sum_{j \in I} g_j^2 h + \sum_{j \in I^c} \frac{g_j^2}{\tilde{f}_j} h \right)^{1/2}.$$

Theorem 2.1. *Assume (1.7) holds. Let $\tilde{f} = \{\tilde{f}_j\}$ be the strict ESD satisfying (1.10), and $g_j(t)$ be the solution to the linearized scheme (2.1) subject to initial data $g_j(0)$. If $\|g(0)\|_{\tilde{f}} < \infty$, then*

$$(2.2) \quad \|g(t)\|_{\tilde{f}} \leq C(1+t)^\xi e^{-rt}, \quad \xi = 1_{\{\mu=s\}},$$

for some C depending on μ , s , $\|b\|_{L^\infty}$, $\|\tilde{f}\|_1$ and $\|g(0)\|_{\tilde{f}}$.

Proof. For the strict ESD considered, and $j \in I$, one has $\frac{d}{dt}g_j = s_j g_j$ and so

$$(2.3) \quad g_j(t) = g_j(0)e^{s_j t},$$

hence by the definition of s

$$\left(\sum_{j \in I} g_j^2 h \right)^{1/2} \leq \left(\sum_{j \in I} |g_j(0)|^2 h \right)^{1/2} e^{-st}.$$

For $j \in I^c$, $s_j = 0$, and we have

$$\frac{dg_j}{dt} = -\tilde{f}_j h \sum_{i \in I^c} \bar{b}_{ji} g_i - \tilde{f}_j h \sum_{i \in I} \bar{b}_{ji} g_i.$$

Therefore

$$\begin{aligned}
\frac{d}{dt} \left(\sum_{j \in I^c} \frac{g_j^2}{f_j} h \right) &= -2h^2 \sum_{j, i \in I^c} \bar{b}_{ji} g_i g_j - 2h^2 \sum_{j \in I^c} g_j \sum_{i \in I} \bar{b}_{ji} g_i \\
&\leq -2h^2 f_m \lambda_{\min} \sum_{j \in I^c} \frac{g_j^2}{f_j} + 2h^2 e^{-st} \sum_{j \in I^c} \left(|g_j| \sum_{i \in I} \bar{b}_{ji} |g_i(0)| \right) \\
&\leq -2\mu \sum_{j \in I^c} \frac{g_j^2}{f_j} h + 2h^{\frac{3}{2}} e^{-st} \left(\sum_{j \in I^c} \frac{g_j^2}{f_j} h \right)^{1/2} \left[\sum_{j \in I^c} \tilde{f}_j \left(\sum_{i \in I} \bar{b}_{ji} |g_i(0)| \right)^2 \right]^{1/2} \\
&\leq -2\mu \sum_{j \in I^c} \frac{g_j^2}{f_j} h + 2C_1 e^{-st} \left(\sum_{j \in I^c} \frac{g_j^2}{f_j} h \right)^{1/2},
\end{aligned}$$

where $C_1 = \sqrt{2} \|b\|_{L^\infty} |\tilde{f}|_1^{\frac{1}{2}} \|g(0)\|_{L^2(I)}$. Calling $A(t) = \left(\sum_{j \in I^c} \frac{g_j^2}{f_j} h \right)^{1/2}$, we have

$$\frac{dA}{dt} \leq -\mu A + C_1 e^{-st},$$

which upon integration gives

$$A \leq \begin{cases} \left(A(0) - \frac{C_1}{\mu-s} \right) e^{-\mu t} + \frac{C_1}{\mu-s} e^{-st}, & \mu - s \neq 0, \\ (A(0) + C_1 t) e^{-st}, & \mu - s = 0. \end{cases}$$

Therefore, one has

$$A \leq C(1+t)^\xi e^{-rt},$$

with $\xi = 1_{\{\mu=s\}}$, and

$$C = A(0) + C_1 |\mu - s|^{-1} 1_{\{\mu \neq s\}} + C_1 1_{\{\mu=s\}}.$$

Then

$$\begin{aligned}
\|g(t)\|_{\tilde{f}} &\leq \left(\sum_{j \in I} |g_j(0)|^2 h \right)^{1/2} e^{-st} + A(t) \\
&\leq (\|g(0)\|_{\tilde{f}} + C_1 |\mu - s|^{-1} 1_{\{\mu \neq s\}} + C_1 1_{\{\mu=s\}}) (1+t)^\xi e^{-rt}.
\end{aligned}$$

This ensures claimed estimate (2.2). \square

2.2. Nonlinear stability. We now turn to the nonlinear stability of the ESD under assumption (1.10).

Theorem 2.2. *Assume (1.7) holds. Let $f_j(t)$ be the solution to (1.4), associated with the strict ESD $\tilde{f} = \{\tilde{f}_j\}$, then there exists $\delta^* > 0$ such that for any $\delta \in (0, \delta^*)$ if*

$$\|f(0) - \tilde{f}\|_{\tilde{f}} \leq \delta,$$

then

$$(2.4) \quad \|f(t) - \tilde{f}\|_{\tilde{f}} \leq C(1+t)^\xi e^{-rt}, \quad \xi = 1_{\{\mu=s\}},$$

for some C depending on $\mu, s, \|b\|_\infty, \|\tilde{f}\|_1$ and the norm of the initial data.

Proof. 1. Symmetrization with weight depending on \tilde{f} .

For the strict ESD considered, we substitute $f_j = \tilde{f}_j + g_j$ into (1.4) so that

$$\frac{d}{dt} g_j = s_j g_j - \tilde{f}_j h \sum_{i=1}^N \bar{b}_{ji} g_i - g_j h \sum_{i=1}^N \bar{b}_{ji} g_i, \quad j = 1, \dots, N.$$

For $j \in I$, $\frac{d}{dt}g_j = s_j g_j - g_j h \sum_{i=1}^N \bar{b}_{ji} g_i$, and thus

$$\begin{aligned} \frac{d}{dt} \sum_{j \in I} g_j^2 h &= 2h \sum_{j \in I} s_j g_j^2 - 2h^2 \sum_{j \in I} g_j^2 \sum_{i=1}^N \bar{b}_{ji} g_i \\ &\leq -2s \sum_{j \in I} g_j^2 h + 2\|b\|_\infty \|g\|_1 \sum_{j \in I} g_j^2 h. \end{aligned}$$

For $j \in I^c$, $s_j = 0$, and

$$\frac{d}{dt}g_j = -\tilde{f}_j h \sum_{i \in I^c} \bar{b}_{ji} g_i - \tilde{f}_j h \sum_{i \in I} \bar{b}_{ji} g_i - g_j h \sum_{i=1}^N \bar{b}_{ji} g_i,$$

so that

$$\begin{aligned} \frac{d}{dt} \left(\sum_{j \in I^c} \frac{g_j^2}{\tilde{f}_j} h \right) &= -2h^2 \sum_{j, i \in I^c} \bar{b}_{ji} g_i g_j - 2h^2 \sum_{j \in I^c} \left(g_j \sum_{i \in I} \bar{b}_{ji} g_i \right) \\ &\quad - 2h^2 \sum_{j \in I^c} \left(\frac{g_j^2}{\tilde{f}_j} \sum_{i=1}^N \bar{b}_{ji} g_i \right) \\ &\leq -2\mu \sum_{j \in I^c} \frac{g_j^2}{\tilde{f}_j} h + 2\|\tilde{f}\|_1^{\frac{1}{2}} \|b\|_{L^\infty} \left(\sum_{j \in I} |g_j| h \right) \left(\sum_{j \in I^c} \frac{g_j^2}{\tilde{f}_j} h \right)^{\frac{1}{2}} \\ &\quad + 2\|b\|_{L^\infty} \|g\|_1 \left(\sum_{j \in I^c} \frac{g_j^2}{\tilde{f}_j} h \right). \end{aligned}$$

Here we have used the Cauchy-Schwarz inequality in bounding the second term.

2. Coupling the two quantities.

Let

$$A_1 = \left(\sum_{j \in I} g_j^2 h \right)^{1/2}, \quad A_2 = \left(\sum_{j \in I^c} \frac{g_j^2}{\tilde{f}_j} h \right)^{1/2},$$

then

$$\begin{cases} \frac{dA_1}{dt} \leq -sA_1 + \|b\|_\infty \|g\|_1 A_1, \\ \frac{dA_2}{dt} \leq -\mu A_2 + \|b\|_{L^\infty} \|g\|_1 A_2 + \|b\|_{L^\infty} \|\tilde{f}\|_1^{\frac{1}{2}} \left(\sum_{j \in I} |g_j| h \right). \end{cases}$$

Further simplification by using $\sum_{j \in I} |g_j| h \leq \sqrt{2} A_1$ and setting $C_2 = \|b\|_{L^\infty} \left(\frac{\|\tilde{f}\|_1}{2} \right)^{1/2}$ leads to

$$(2.5a) \quad \frac{dA_1}{dt} \leq (-s + \|b\|_\infty \|g\|_1) A_1,$$

$$(2.5b) \quad \frac{dA_2}{dt} \leq (-\mu + \|b\|_{L^\infty} \|g\|_1) A_2 + 2C_2 A_1.$$

3. Decay estimates using a Lyapunov functional.

Set

$$L := A_1^2 + \alpha^2 A_2^2,$$

so that

$$(2.6) \quad 2\alpha A_1 A_2 \leq L.$$

Here α is to be determined, so that the exponential decay of L is ensured, yet with largest possible initial data.

A direct calculation gives

$$\begin{aligned} \dot{L} &= 2A_1 \dot{A}_1 + 2\alpha^2 A_2 \dot{A}_2 \\ &\leq -2sA_1^2 - 2\mu\alpha^2 A_2^2 + 4C_2\alpha^2 A_1 A_2 + 2\|b\|_{L^\infty} \|g\|_1 (A_1^2 + \alpha^2 A_2^2). \end{aligned}$$

Proceeding with (2.6) and

$$(2.7) \quad \|g\|_1 \leq \sqrt{2}A_1 + \|\tilde{f}\|_1^{1/2} A_2 \leq \left(\sqrt{2} + \alpha^{-1} \|\tilde{f}\|_1^{1/2} \right) \sqrt{L},$$

we see that

$$\dot{L} \leq - \left[2r - 2C_2\alpha - 2\|b\|_{L^\infty} \left(\sqrt{2} + \alpha^{-1} \|\tilde{f}\|_1^{1/2} \right) \sqrt{L} \right] L.$$

This implies that for any $\alpha \in (0, r/C_2)$, if

$$(2.8) \quad \sqrt{L(0)} < k(\alpha) := \frac{r - C_2\alpha}{\left(\sqrt{2} + \alpha^{-1} \|\tilde{f}\|_1^{1/2} \right) \|b\|_{L^\infty}},$$

then L is decreasing in time, and its decay rate is governed by the linear part as

$$(2.9) \quad L \leq C_3 e^{-2(r-C_2\alpha)t},$$

where C_3 is given by

$$C_3 = \sup_{t \geq 0} \frac{L(0)}{\left[1 + \sqrt{L(0)}/k(\alpha) (e^{-(r-C_2\alpha)t} - 1) \right]^2} = \frac{L(0)}{\left[1 - \sqrt{L(0)}/k(\alpha) \right]^2}.$$

It suffices to select α such that $k(\alpha)$ achieves its maximum. One can verify that

$$\frac{(k_1 - x)x}{x + k_2} \leq (\sqrt{k_1 + k_2} - \sqrt{k_2})^2,$$

and this maximum is achieved at $x = \sqrt{k_2^2 + k_1 k_2} - k_2$. This when applied to $k(\alpha)$ with $k_1 = r/C_2$ and $k_2 = (\|\tilde{f}\|_1/2)^{1/2}$ leads to

$$(2.10) \quad \max k(\alpha) = k(\alpha^*) = \frac{C_2}{\sqrt{2}\|b\|_{L^\infty}} \left(\sqrt{r/C_2 + (\|\tilde{f}\|_1/2)^{1/2}} - \sqrt{(\|\tilde{f}\|_1/2)^{1/2}} \right)^2,$$

where

$$\alpha^* = \sqrt{\|\tilde{f}\|_1/2 + r/C_2(\|\tilde{f}\|_1/2)^{1/2}} - (\|\tilde{f}\|_1/2)^{1/2}.$$

Recall that $C_2 = \|b\|_{L^\infty} \left(\frac{\|\tilde{f}\|_1}{2} \right)^{1/2}$, and

$$\sqrt{L(0)} \leq \max\{1, \alpha\} \|f(0) - \tilde{f}\|_{\tilde{f}} \leq \delta \max\{1, \alpha\}.$$

Hence (2.8) is ensured to hold if we choose δ^* such that

$$(2.11) \quad \delta^* = \frac{k(\alpha^*)}{\max\{1, \alpha^*\}} = \frac{(\alpha^*)^2}{\sqrt{2} \max\{1, \alpha^*\}}, \quad \alpha^* = \sqrt{\frac{r}{\|b\|_{L^\infty}} + \frac{\|\tilde{f}\|_1}{2}} - \sqrt{\frac{\|\tilde{f}\|_1}{2}}.$$

This value of α is indeed less than $r/C_2 = \sqrt{2}r / \left(\|b\|_{L^\infty} \|\tilde{f}\|_1^{1/2} \right)$, as required.

4. Optimal decay rates.

We use C to denote different constants from line to line. Equation (2.5a), combining with (2.7) and (2.9), leads to

$$(2.12) \quad \frac{dA_1}{dt} \leq -sA_1 + Ce^{-kt}A_1, \quad k = r - C_2\alpha^* > 0.$$

Integration of (2.12) gives

$$A_1(t) \leq A_1(0)e^{C\frac{1}{k}(1-e^{-kt})}e^{-st} \leq Ce^{-st}.$$

Substitution of this into (2.5b) yields

$$\frac{dA_2}{dt} \leq -\mu A_2 + Ce^{-kt}A_2 + Ce^{-st}.$$

This upon rewriting gives

$$\frac{d}{dt} \left[A_2 e^{\mu t + C(e^{-kt}-1)/k} \right] \leq C e^{C(e^{-kt}-1)/k} e^{(\mu-s)t} \leq C e^{(\mu-s)t}.$$

Hence,

$$A_2(t) \leq \begin{cases} Ce^{-rt}, & \mu - s \neq 0, \\ C(1+t)e^{-st}, & \mu - s = 0. \end{cases}$$

These when combined with $\|f - \tilde{f}\|_{\tilde{f}}^2 = A_1^2 + A_2^2$ lead to the estimate (2.4). \square

3. CONVERGENCE RATE FOR THE FULLY DISCRETE SCHEME

For the fully discrete scheme

$$(3.1) \quad \frac{f_j^{n+1} - f_j^n}{\Delta t} = f_j^{n+1} \left(\bar{a}_j - h \sum_{i=1}^N \bar{b}_{ji} f_i^n \right), \quad j = 1, \dots, N,$$

we establish the exponential convergence towards the strict ESD and algebraic convergence towards the generic ESD for Δt suitable small.

3.1. Exponential convergence.

Theorem 3.1. *Assume (1.7) holds. Let f_j^n be the numerical solution to (3.1), associated with the strict ESD $\tilde{f} = \{\tilde{f}_j\}$. If Δt satisfies*

$$(3.2) \quad \Delta t \leq \frac{\min\{s, \mu/2\}}{\|b\|_{L^\infty}^2 \|\tilde{f}\|_1^2},$$

then there exists $\delta^* > 0$ such that for any $\delta \in (0, \delta^*)$ if

$$\|f^0 - \tilde{f}\|_{\tilde{f}} \leq \delta,$$

then

$$(3.3) \quad \|f^n - \tilde{f}\|_{\tilde{f}} \leq C(1 + n\Delta t)^\xi \max\{\sqrt{K_s}, K_*\}^n, \quad \xi = 1_{\{\sqrt{K_s} = K_*\}},$$

where

$$K_s = \frac{1}{\sqrt{1 + 2s\Delta t}}, \quad K_* = \frac{1}{\sqrt{1 + 2\mu\Delta t}} \left(1 + \frac{2\Delta t^2 \|b\|_{L^\infty}^2 \|\tilde{f}\|_1^2}{\sqrt{1 + 2\mu\Delta t}} \right) < 1,$$

and C depends on $\mu, s, \|b\|_\infty, \|\tilde{f}\|_1$ and the norm of the initial data.

Proof. The proof follows steps similar to the semi-discrete case.

1. Symmetrization with weight depending on \tilde{f} .

For the strict ESD considered, we substitute $f_j^n = \tilde{f}_j + g_j^n$ into (3.1) so that

$$\frac{g_j^{n+1} - g_j^n}{\Delta t} = s_j g_j^{n+1} - \tilde{f}_j h \sum_{i=1}^N \bar{b}_{ji} g_i^n - g_j^{n+1} h \sum_{i=1}^N \bar{b}_{ji} g_i^n, \quad j = 1, \dots, N.$$

Then

$$\frac{g_j^{n+1} - g_j^n}{\Delta t} = s_j g_j^{n+1} - g_j^{n+1} h \sum_{i=1}^N \bar{b}_{ji} g_i^n, \quad j \in I.$$

Multiplying $g_j^{n+1} h$ on both sides and summing over $j \in I$, we have

$$(3.4) \quad h \sum_{j \in I} \frac{(g_j^{n+1})^2 - g_j^{n+1} g_j^n}{\Delta t} = h \sum_{j \in I} s_j (g_j^{n+1})^2 - h^2 \sum_{j \in I} \left((g_j^{n+1})^2 \sum_{i=1}^N \bar{b}_{ji} g_i^n \right).$$

The left hand side of (3.4) may be written as

$$(3.5) \quad h \sum_{j \in I} \left[\frac{(g_j^{n+1})^2 - (g_j^n)^2}{2\Delta t} + \frac{(g_j^{n+1} - g_j^n)^2}{2\Delta t} \right],$$

and the right hand side of (3.4) is bounded from above by

$$-s \sum_{j \in I} (g_j^{n+1})^2 h + \|b\|_{L^\infty} \|g^n\|_1 \sum_{j \in I} (g_j^{n+1})^2 h.$$

Hence

$$(3.6) \quad h \sum_{j \in I} \frac{(g_j^{n+1})^2 - (g_j^n)^2}{2\Delta t} \leq -s \sum_{j \in I} (g_j^{n+1})^2 h + \|b\|_{L^\infty} \|g^n\|_1 \sum_{j \in I} (g_j^{n+1})^2 h.$$

For $j \in I^c$, $s_j = 0$, and

$$\frac{g_j^{n+1} - g_j^n}{\Delta t} = -\tilde{f}_j h \sum_{i=1}^N \bar{b}_{ji} g_i^n - g_j^{n+1} h \sum_{i=1}^N \bar{b}_{ji} g_i^n.$$

Against $\frac{g_j^{n+1} h}{\tilde{f}_j}$ on both sides, summation over $j \in I^c$ gives

$$(3.7) \quad h \sum_{j \in I^c} \frac{(g_j^{n+1})^2 - g_j^{n+1} g_j^n}{\tilde{f}_j \Delta t} = -h^2 \sum_{j \in I^c} \left(g_j^{n+1} \sum_{i=1}^N \bar{b}_{ji} g_i^n \right) - h^2 \sum_{j \in I^c} \left[\frac{(g_j^{n+1})^2}{\tilde{f}_j} \sum_{i=1}^N \bar{b}_{ji} g_i^n \right].$$

The term on the left hand side is treated same way as in (3.5), and the last term on the right is bounded by

$$\|b\|_{L^\infty} \|g^n\|_1 \sum_{j \in I^c} \frac{(g_j^{n+1})^2}{\tilde{f}_j} h.$$

We focus on the estimate of the first term on the right hand side of (3.7), which can be estimated by

$$\begin{aligned}
&\leq -h^2 \sum_{j \in I^c} \left(g_j^{n+1} \sum_{i \in I^c} \bar{b}_{ji} g_i^{n+1} \right) + h^2 \sum_{j \in I^c} \left[g_j^{n+1} \sum_{i \in I^c} \bar{b}_{ji} (g_i^{n+1} - g_i^n) \right] \\
&\quad - h^2 \sum_{j \in I^c} \left(g_j^{n+1} \sum_{i \in I} \bar{b}_{ji} g_i^n \right) \\
&\leq -\mu \sum_{j \in I^c} \frac{(g_j^{n+1})^2}{\tilde{f}_j} h - \Delta t h^3 \sum_{j \in I^c} \left[g_j^{n+1} \sum_{i \in I^c} \bar{b}_{ji} (\tilde{f}_i + g_i^{n+1}) \sum_{k=1}^N \bar{b}_{ik} g_k^n \right] \\
&\quad + \|b\|_{L^\infty} \left(\sum_{i \in I} |g_i^n| h \right) \|\tilde{f}\|_1^{1/2} \left[\sum_{j \in I^c} \frac{(g_j^{n+1})^2}{\tilde{f}} h \right]^{1/2}.
\end{aligned}$$

Applying the Cauchy-Schwarz inequality to estimate the second term above gives

$$\begin{aligned}
&-\Delta t h^3 \sum_{j \in I^c} \left[g_j^{n+1} \sum_{i \in I^c} \bar{b}_{ji} \Delta t (\tilde{f}_i + g_i^{n+1}) \sum_{k=1}^N \bar{b}_{ik} g_k^n \right] \\
&\leq \Delta t h^2 \left[\sum_{j \in I^c} \frac{(g_j^{n+1})^2}{\tilde{f}} h \right]^{1/2} \left\{ \sum_{j \in I^c} \tilde{f}_j h \left[\sum_{i \in I^c} \bar{b}_{ji} (\tilde{f}_i + g_i^{n+1}) \sum_{k=1}^N \bar{b}_{ik} g_k^n \right]^2 \right\}^{1/2} \\
&\leq \Delta t h^2 \|b\|_{L^\infty} \|\tilde{f}\|_1^{1/2} \sum_{i \in I^c} \left[(\tilde{f}_i + |g_i^{n+1}|) \sum_{k=1}^N \bar{b}_{ik} |g_k^n| \right] \left[\sum_{j \in I^c} \frac{(g_j^{n+1})^2}{\tilde{f}} h \right]^{1/2} \\
&\leq \Delta t \|b\|_{L^\infty}^2 \|\tilde{f}\|_1^{1/2} \|g^n\|_1 \left\{ \|\tilde{f}\|_1 + \|\tilde{f}\|_1^{1/2} \left[\sum_{j \in I^c} \frac{(g_j^{n+1})^2}{\tilde{f}} h \right]^{1/2} \right\} \left[\sum_{j \in I^c} \frac{(g_j^{n+1})^2}{\tilde{f}} h \right]^{1/2}.
\end{aligned}$$

2. Coupling the two quantities.

We set

$$A_1^n = \sum_{j \in I} (g_j^n)^2 h, \quad A_2^n = \sum_{j \in I^c} \frac{(g_j^n)^2}{\tilde{f}_j} h,$$

so that the above estimates may be written as

$$(3.8a) \quad \frac{A_1^{n+1} - A_1^n}{2\Delta t} \leq -s A_1^{n+1} + \|b\|_\infty \|g^n\|_1 A_1^{n+1},$$

$$\begin{aligned}
(3.8b) \quad \frac{A_2^{n+1} - A_2^n}{2\Delta t} &\leq -\mu A_2^{n+1} + \|b\|_{L^\infty} \|g^n\|_1 A_2^{n+1} + \|b\|_{L^\infty} \|\tilde{f}\|_1^{1/2} \left(\sum_{i \in I} |g_i^n| h \right) \sqrt{A_2^{n+1}} \\
&\quad + \Delta t \|b\|_{L^\infty}^2 \|\tilde{f}\|_1 \|g^n\|_1 A_2^{n+1} + \Delta t \|b\|_{L^\infty}^2 \|\tilde{f}\|_1^{3/2} \|g^n\|_1 \sqrt{A_2^{n+1}}.
\end{aligned}$$

3. Decay estimates using a Lyapunov functional.

Set

$$(3.9) \quad L^n := A_1^n + \alpha^2 A_2^n,$$

Next we determine the range of the initial data so that L^n decays in n , with proper choices of α and Δt .

Note that

$$\sum_{i \in I} |g_i^n| h \leq \sqrt{2A_1^n},$$

with which, (3.8) and (3.9), it follows that

$$\begin{aligned} \frac{L^{n+1}-L^n}{2\Delta t} \leq & -rL^{n+1} + \|b\|_{L^\infty} \|g^n\|_1 L^{n+1} + \sqrt{2}\alpha^2 \|b\|_{L^\infty} \|\tilde{f}\|_1^{1/2} \sqrt{A_1^n A_2^{n+1}} \\ & + \Delta t \|b\|_{L^\infty}^2 \|\tilde{f}\|_1 \|g^n\|_1 L^{n+1} + \Delta t \alpha \|b\|_{L^\infty}^2 \|\tilde{f}\|_1^{3/2} \|g^n\|_1 \sqrt{L^{n+1}}. \end{aligned}$$

Proceeding with

$$(3.10) \quad \|g^n\|_1 \leq \sqrt{2A_1^n} + \sqrt{\|\tilde{f}\|_1 A_2^n} \leq \left(\sqrt{2} + \alpha^{-1} \|\tilde{f}\|_1^{1/2} \right) \sqrt{L^n},$$

we see that

$$(3.11) \quad \frac{L^{n+1} - L^n}{2\Delta t} \leq -rL^{n+1} + c_1 \sqrt{L^n} L^{n+1} + c_2 \sqrt{L^n L^{n+1}},$$

where

$$\begin{aligned} c_1 &= \|b\|_{L^\infty} (1 + \Delta t \|b\|_{L^\infty} \|\tilde{f}\|_1) \left(\sqrt{2} + \alpha^{-1} \|\tilde{f}\|_1^{1/2} \right), \\ c_2 &= \sqrt{2}\alpha \|b\|_{L^\infty} \|\tilde{f}\|_1^{1/2} + \Delta t \alpha \|b\|_{L^\infty}^2 \|\tilde{f}\|_1^{3/2} \left(\sqrt{2} + \alpha^{-1} \|\tilde{f}\|_1^{1/2} \right). \end{aligned}$$

Using $\sqrt{L^n L^{n+1}} \leq L^n/2 + L^{n+1}/2$ in (3.11) we obtain

$$L^{n+1}(1 + 2r\Delta t - 2c_1\Delta t\sqrt{L^n} - c_2\Delta t) \leq (1 + c_2\Delta t)L^n.$$

Note that if the time step is taken as

$$(3.12) \quad 0 < \Delta t < \frac{r - \sqrt{2}\alpha \|b\|_{L^\infty} \|\tilde{f}\|_1^{1/2}}{\|b\|_{L^\infty}^2 \|\tilde{f}\|_1^{3/2} \left(\sqrt{2}\alpha + \|\tilde{f}\|_1^{1/2} \right)},$$

then $c_2 < r$. Therefore, $L^{n+1} < L^n$ provided

$$c_2 < r - c_1\sqrt{L^0} \leq r - c_1\sqrt{L^n}.$$

This implies that if

$$(3.13) \quad \sqrt{L^0} < k(\alpha) := \frac{r - c_2}{c_1},$$

then L^n is strictly decreasing in n , and

$$L^{n+1} \leq KL^n, \quad K := \frac{(1 + c_2\Delta t)}{1 + 2r\Delta t - 2c_1\Delta t\sqrt{L^0} - c_2\Delta t} < 1.$$

Therefore exponential decay holds

$$(3.14) \quad L^n \leq K^n L^0.$$

We now check how to choose α so that $k(\alpha)$ is maximized for each fixed Δt satisfying

$$\Delta t < \frac{r}{\|b\|_{L^\infty}^2 \|\tilde{f}\|_1^2}.$$

Note that (3.12) is equivalent to the following requirement

$$(3.15) \quad \alpha < \beta(\Delta t) := \frac{r - \Delta t \|b\|_{L^\infty}^2 \|\tilde{f}\|_1^2}{\sqrt{2\|\tilde{f}\|_1 \|b\|_{L^\infty} (1 + \Delta t \|b\|_{L^\infty} \|\tilde{f}\|_1)}}.$$

Rewriting $k(\alpha)$ as

$$(3.16) \quad k(\alpha) = \frac{\alpha(\beta(\Delta t) - \alpha)}{\alpha + \sqrt{\|\tilde{f}\|_1/2}} \cdot \sqrt{\|\tilde{f}\|_1}.$$

The maximum of this function is achieved at α^* , with

$$\alpha^* = \sqrt{\|\tilde{f}\|_1/2 + (\|\tilde{f}\|_1/2)^{1/2}\beta(\Delta t) - (\|\tilde{f}\|_1/2)^{1/2}} = \frac{\beta(\Delta t)}{1 + \sqrt{1 + (2/\|\tilde{f}\|_1)^{1/2}\beta(\Delta t)}}.$$

Such an α^* clearly satisfies (3.15). Moreover,

$$k(\alpha^*) = \sqrt{\|\tilde{f}\|_1} \left(\sqrt{\beta(\Delta t) + (\|\tilde{f}\|_1/2)^{1/2}} - \sqrt{(\|\tilde{f}\|_1/2)^{1/2}} \right)^2 = \sqrt{2}(\alpha^*)^2.$$

Furthermore,

$$\sqrt{L^0} \leq \max\{1, \alpha\} \|f^0 - \tilde{f}\|_{\tilde{f}} \leq \delta \max\{1, \alpha\}.$$

Hence (3.13) is ensured to hold if we choose δ^* such that

$$(3.17) \quad \delta^* = \frac{k(\alpha^*)}{\max\{1, \alpha^*\}} = \frac{\sqrt{2}(\alpha^*)^2}{\max\{1, \alpha^*\}}.$$

4. Optimal decay rates.

In the estimate to follow, we use C to denote different constants from line to line if applicable. From (3.8a), combining with (3.10) and (3.14), it follows

$$\frac{A_1^{n+1} - A_1^n}{2\Delta t} \leq - (s - \gamma K^{n/2}) A_1^{n+1},$$

for $\gamma = \|b\|_{L^\infty} \left(\sqrt{2} + \alpha^{-1} \|\tilde{f}\|_1^{1/2} \right) \sqrt{L^0} < s$ which can be obtained from (3.16). Then

$$(3.18) \quad \begin{aligned} A_1^n &\leq \frac{A_1^{n-1}}{1+2s\Delta t-2\gamma\Delta tK^{(n-1)/2}} \\ &\leq \frac{A_1^0}{(1+2s\Delta t)^n} \prod_{i=0}^{n-1} \frac{1+2s\Delta t}{1+2s\Delta t-2\gamma\Delta tK^{i/2}} \\ &\leq C_\gamma A_1^0 K_s^n, \end{aligned}$$

where $K_s = \frac{1}{1+2s\Delta t}$. In fact, the product may be estimated as follows

$$\begin{aligned} &\leq \prod_{i=0}^{n-1} (1 + 2\gamma\Delta t K^{i/2}) \\ &\leq \exp \left(\sum_{i=0}^{n-1} \log (1 + 2\gamma\Delta t K^{i/2}) \right) \\ &\leq \exp \left(2\gamma\Delta t \sum_{i=0}^{n-1} (\sqrt{K})^i \right) \leq \exp \left(\frac{2\gamma\Delta t}{1 - \sqrt{K}} \right), \end{aligned}$$

leading to the claimed bound.

We now estimate the decay rate of A_2^n . Substitution of (3.10), (3.14) and (3.18) into (3.8b) yields

$$\begin{aligned} \frac{A_2^{n+1} - A_2^n}{2\Delta t} &\leq -\mu A_2^{n+1} + c_1 \sqrt{L^0} K^{n/2} A_2^{n+1} + \sqrt{2} \|b\|_{L^\infty} \|\tilde{f}\|_1^{1/2} \sqrt{A_1^n} \sqrt{A_2^{n+1}} \\ &\quad + \Delta t \|b\|_{L^\infty}^2 \|\tilde{f}\|_1^{3/2} (\sqrt{2A_1^n} + \sqrt{\|\tilde{f}\|_1 A_2^n}) \sqrt{A_2^{n+1}} \\ &\leq -\mu A_2^{n+1} + c_1 \sqrt{L^0} K^{n/2} A_2^{n+1} + C_1 K_s^{n/2} \sqrt{A_2^{n+1}} \\ &\quad + \Delta t \|b\|_{L^\infty}^2 \|\tilde{f}\|_1^2 \sqrt{A_2^n A_2^{n+1}}, \end{aligned}$$

where $C_1 = \sqrt{2} \|b\|_{L^\infty} \|\tilde{f}\|_1^{1/2} (1 + \|b\|_{L^\infty} \|\tilde{f}\|_1 \Delta t) \sqrt{C_\gamma A_1^0}$. Hence

$$A_2^{n+1} - 2e_n d_n^2 \sqrt{A_2^{n+1}} - d_n^2 A_2^n \leq 0,$$

where

$$\begin{aligned} d_n &= \frac{1}{\left[1 + 2\Delta t \left(\mu - c_1 \sqrt{L^0} K^{n/2}\right)\right]^{1/2}}, \\ e_n &= C_1 \Delta t K_s^{n/2} + \Delta t^2 \|b\|_{L^\infty}^2 \|\tilde{f}\|_1^2 \sqrt{A_2^n}. \end{aligned}$$

This gives

$$\begin{aligned} \sqrt{A_2^{n+1}} &\leq e_n d_n^2 + \sqrt{e_n^2 d_n^4 + d_n^2 A_2^n} \\ &\leq 2e_n d_n^2 + d_n \sqrt{A_2^n} \\ &\leq 2C_1 \Delta t K_s^{n/2} d_n^2 + \tilde{d}_n \sqrt{A_2^n}, \quad \tilde{d}_n := d_n (1 + 2\Delta t^2 \|b\|_{L^\infty}^2 \|\tilde{f}\|_1^2 d_n). \end{aligned}$$

By induction,

$$\sqrt{A_2^n} \leq \left(\prod_{i=0}^{n-1} \tilde{d}_i \right) \sqrt{A_2^0} + 2C_1 \Delta t \sum_{i=0}^{n-1} \left(K_s^{i/2} d_i^2 \prod_{j=i+1}^{n-1} \tilde{d}_j \right).$$

For fixed Δt , $d_\infty := \sqrt{K_\mu} \leq d_j \leq 1$, a similar estimate as in (3.18) gives

$$\begin{aligned} \prod_{j=i}^{n-1} \tilde{d}_j &= (\tilde{d}_\infty)^{n-i} \prod_{j=i}^{n-1} \frac{d_j}{d_\infty} \cdot \prod_{j=i}^{n-1} \frac{1 + 2\Delta t^2 \|b\|_{L^\infty}^2 \|\tilde{f}\|_1^2 d_j}{1 + 2\Delta t^2 \|b\|_{L^\infty}^2 \|\tilde{f}\|_1^2 d_\infty} \\ &\leq (\tilde{d}_\infty)^{n-i} \prod_{j=i}^{n-1} \frac{d_j}{d_\infty} \cdot \prod_{j=i}^{n-1} \frac{d_j}{d_\infty} \\ &\leq \exp\left(\frac{2c_1 \sqrt{L_0} \Delta t}{1 - \sqrt{K}}\right) (K_*)^{(n-i)}, \quad \text{for } i = 0, 1, 2, \dots, n-1, \end{aligned}$$

where

$$K_* := \tilde{d}_\infty = \sqrt{K_\mu} \left(1 + 2\Delta t^2 \|b\|_{L^\infty}^2 \|\tilde{f}\|_1^2 \sqrt{K_\mu}\right),$$

which is strictly less than one provided (3.2) is satisfied. In fact, from (3.2) it follows that

$$\begin{aligned} \Delta t \frac{\|b\|_{L^\infty}^2 \|\tilde{f}\|_1^2}{\mu} &\leq \frac{1}{2} < 1 - \frac{1}{1 + \sqrt{1 + 2\mu\Delta t}} \\ &= \frac{\sqrt{1 + 2\mu\Delta t}(\sqrt{1 + 2\mu\Delta t} - 1)}{2\mu\Delta t}, \end{aligned}$$

which yields

$$K_* = \sqrt{K_\mu} + 2\Delta t^2 \|b\|_{L^\infty}^2 \|\tilde{f}\|_1^2 K_\mu < 1.$$

Furthermore,

$$\begin{aligned} 2C_1\Delta t \sum_{i=0}^{n-1} \left(K_s^{i/2} d_i^2 \prod_{j=i+1}^{n-1} \tilde{d}_j \right) &\leq 2C_1\Delta t \sum_{i=0}^{n-1} \left(K_s^{i/2} d_i \prod_{j=i}^{n-1} \tilde{d}_j \right) \\ &\leq C\Delta t K_*^n \sum_{i=0}^{n-1} \left(\frac{\sqrt{K_s}}{K_*} \right)^i. \end{aligned}$$

This is bounded by $CK_s^{n/2}n\Delta t$ if $K_* = \sqrt{K_s}$, and if $K_* \neq \sqrt{K_s}$,

$$\Delta t K_* \frac{K_*^n - \sqrt{K_s}^n}{K_* - \sqrt{K_s}} \leq \frac{K_*\Delta t}{|K_* - \sqrt{K_s}|} \max\{\sqrt{K_s}, K_*\}^n \leq C \max\{\sqrt{K_s}, K_*\}^n.$$

The consistency of this bound with the semi-discrete case can be seen from the fact that

$$\lim_{\Delta t \rightarrow 0} \frac{K_*\Delta t}{|K_* - \sqrt{K_s}|} = \frac{1}{|s - \mu|}.$$

In summary, we have

$$\sqrt{A_2^n} \leq C \begin{cases} \max\{\sqrt{K_s}, K_*\}^n, & K_* \neq \sqrt{K_s}, \\ K_s^{n/2} n\Delta t, & K_* = \sqrt{K_s}. \end{cases}$$

These when combined with $\|f - \tilde{f}\|_{\tilde{f}} \leq \sqrt{A_1} + \sqrt{A_2}$ lead to the estimate (3.3). \square

3.2. Algebraic convergence. It was shown in [12] that the numerical solution of (3.1) converges to the ESD in weighted norm $\|\cdot\|_b$. In this section we investigate the convergence rate of the numerical solution toward the ESD in this norm.

Define the relative entropy

$$(3.19) \quad F^n = \sum_{j=1}^N \left(\tilde{f}_j \log \left(\frac{\tilde{f}_j}{f_j^n} \right) + f_j^n - \tilde{f}_j \right) h,$$

and a nonlinear function

$$H(f) = \frac{f^T B f}{2} h^2 - a^T f h,$$

with $f = (f_1, f_2, \dots, f_N)^T$ and $a = (\bar{a}_1, \bar{a}_2, \dots, \bar{a}_N)^T$.

For later use, we present a uniform l^1 -bound of the numerical solution when $b \geq b_f > 0$.

Lemma 3.1. *Assume (1.7) holds. Let f_j^n be the numerical solution generated from scheme (3.1) with nonnegative initial data $f_j^0 \geq 0$ for all $j = 1, \dots, N$, and $\|f^0\|_1 < \infty$. Then for any $n > 0$,*

$$(3.20) \quad \|f^n\|_1 \leq \max \left\{ \|f^0\|_1, \frac{\|a\|_{L^\infty}}{b_f} \right\},$$

provided

$$(3.21) \quad \Delta t \leq \frac{1}{\|a\|_{L^\infty}}.$$

Proof. From (3.1) it follows that if $f_j^n \geq 0$ and (3.21) holds, then $f_j^{n+1} \geq 0$, hence the numerical solution remains non-negative at all time steps.

Let $M^n = h \sum_{j=1}^N f_j^n = \|f^n\|_1$ and $\gamma = \frac{\|a\|_{L^\infty}}{b_f}$. From scheme (3.1) it follows

$$\begin{aligned} M^{n+1} - M^n &= \Delta t \left(h \sum_{j=1}^N f_j^{n+1} \bar{a}_j - h \sum_{j=1}^N f_j^{n+1} h \sum_{i=1}^N \bar{b}_{ji} f_i^n \right) \\ &\leq \Delta t (\|a\|_{L^\infty} M^{n+1} - b_f M^{n+1} M^n) \\ &= -\Delta t b_f M^{n+1} (M^n - \gamma). \end{aligned}$$

There are two cases to distinguish:

- i) if $M^n \geq \gamma$, then $M^{n+1} \leq M^n$;
- ii) if $M^n < \gamma$, we rewrite

$$M^{n+1} - \gamma = (M^n - \gamma)(1 - \Delta t M^{n+1} b_f).$$

According to (3.21), we have

$$M^{n+1} - \gamma \leq (M^n - \gamma) \left(1 - \frac{M^{n+1}}{\gamma}\right),$$

which leads to $M^{n+1} \leq \gamma$. Hence,

$$M^{n+1} \leq \max\{M^n, \gamma\} \leq \dots \leq \max\{M^0, \gamma\},$$

which is as desired. □

Lemma 3.2. [12, Corrolary 3.1] *Assume (1.7) holds. Let f_j^n be the numerical solution generated from scheme (3.1) with positive initial data $f_j^0 > 0$ for all $j = 1, \dots, N$. Then*

$$F^{n+1} - F^n \leq -\frac{1}{2} \Delta t \|f^n - \tilde{f}\|_b^2,$$

provided time step Δt is suitably small.

This implies the following assertions:

$$(3.22) \quad \lim_{n \rightarrow \infty} F^n = 0, \quad \lim_{n \rightarrow \infty} \|f^n - \tilde{f}\|_b = 0.$$

The main aim here is to obtain the convergence rate toward the ESD.

Theorem 3.2. *Assume (1.7) holds, and $F^0 < +\infty$. Let f_j^n be the numerical solution generated from scheme (3.1) with positive initial data $f_j^0 > 0$ for all $j = 1, \dots, N$, $\tilde{f} = \{\tilde{f}_j\}$ is the discrete ESD. Then*

$$(3.23) \quad \|f^n - \tilde{f}\|_b^2 \leq \frac{2F^0}{\Delta t n},$$

provided $\Delta t \leq \tau$, where

$$(3.24) \quad \tau = \min \left\{ \frac{\lambda_{\min}}{\lambda_{\max}(2C_1 + 2C_2\lambda_{\max} + C_2\lambda_{\min})}, \frac{2}{C_2\|b\|_{L^\infty}} \right\},$$

where $C_1 = \|a\|_{L^\infty} + \|b\|_{L^\infty}\|\tilde{f}\|_1$ and $C_2 = \max\{\|f^0\|_1, \frac{\|a\|_{L^\infty}}{b_f}\}$.

Proof. We proceed in two steps:

i) we first establish for the relative entropy F^n the dissipation inequality of the form

$$(3.25) \quad F^{n+1} - F^n \leq -\Delta t[H(f^{n+1}) - H(\tilde{f})];$$

ii) we then show that $H(f^n)$ is decreasing in n , i.e.,

$$(3.26) \quad H(f^{n+1}) - H(f^n) \leq 0.$$

We postpone the proof of these two inequalities, while we now use them to prove estimate (3.23). The summation of (3.25) in n gives

$$(3.27) \quad \Delta t \sum_{i=0}^{+\infty} [H(f^{i+1}) - H(\tilde{f})] \leq F^0 - F^\infty = F^0.$$

On the other hand, for any large number n ,

$$(3.28) \quad \begin{aligned} \Delta t \sum_{i=0}^{+\infty} [H(f^{i+1}) - H(\tilde{f})] &\geq \Delta t \sum_{i=0}^{n-1} [H(f^{i+1}) - H(\tilde{f})] \\ &\geq n\Delta t[H(f^n) - H(\tilde{f})], \end{aligned}$$

where we have used (3.26). Combining (3.27) and (3.28), we have

$$0 \leq H(f^n) - H(\tilde{f}) \leq \frac{F^0}{n\Delta t},$$

which when combined with

$$H(f^n) - H(\tilde{f}) = \frac{1}{2}\|f^n - \tilde{f}\|_b^2 - h \sum_{j=1}^N s_j[\tilde{f}]f_j^n \geq \frac{1}{2}\|f^n - \tilde{f}\|_b^2.$$

gives the desired estimate (3.23).

Finally we specify the restrictions on the time step for both (3.25) and (3.26) to hold true. Denote $\|\cdot\|$ the usual Euclidean norm of a vector.

Scheme (3.1) can be written as

$$(3.29) \quad f_j^{n+1} = \frac{f_j^n}{1 - \Delta t \bar{a}_j + h\Delta t (Bf^n)_j},$$

if Δt is suitably small, for example, for

$$(3.30) \quad \Delta t < \|a\|_{L^\infty}^{-1},$$

we have $f_j^{n+1} > 0$ for $f_j^n > 0$. This positivity property will be used below.

We now prove (3.25). Using $\log x \leq x - 1$ for any $x > 0$ and the scheme (3.1), we obtain

$$\begin{aligned} F^{n+1} - F^n &= h \sum_{j=1}^N \left(\tilde{f}_j \log \frac{f_j^n}{f_j^{n+1}} + f_j^{n+1} - f_j^n \right) \\ &\leq h \sum_{j=1}^N \left(\tilde{f}_j \frac{f_j^n - f_j^{n+1}}{f_j^{n+1}} + f_j^{n+1} - f_j^n \right) \\ &= \Delta t h \sum_{j=1}^N (\bar{a}_j - h \sum_{i=1}^N \bar{b}_{ji} f_i^n) (f_j^{n+1} - \tilde{f}_j). \end{aligned}$$

Proceeding with $g^n := f^n - \tilde{f}$, we have

$$\begin{aligned} (3.31) \quad F^{n+1} - F^n &\leq -\Delta t h^2 g^n \cdot B g^{n+1} + \Delta t h \sum_{j=1}^N (\bar{a}_j - h \sum_{i=1}^N \bar{b}_{ji} \tilde{f}_i) g_j^{n+1} \\ &= -\Delta t h^2 g^{n+1} \cdot B g^{n+1} + \Delta t h^2 (g^{n+1} - g^n) \cdot B g^{n+1} \\ &\quad + \Delta t h \sum_{j=1}^N s_j [\tilde{f}] f_j^{n+1} \\ &\leq -\Delta t h^2 g^{n+1} \cdot B g^{n+1} + \Delta t h^2 \|B\|_2 \|g^{n+1} - g^n\| \|g^{n+1}\| \\ &\quad + \Delta t h \sum_{j=1}^N s_j [\tilde{f}] f_j^{n+1}. \end{aligned}$$

Next, we show there exists C^* , which may depend on Δt , such that

$$(3.32) \quad \|g^{n+1} - g^n\| \leq C^* \Delta t \|g^{n+1}\|.$$

Using the fact that $\tilde{f}_j \left(\bar{a}_j - h \sum_{i=1}^N \bar{b}_{ji} \tilde{f}_i \right) = 0$, we have

$$\begin{aligned} (g^{n+1} - g^n)_j &= \Delta t f_j^{n+1} \left[\bar{a}_j - h \sum_{i=1}^N \bar{b}_{ji} (g_i^n + \tilde{f}_i) \right] \\ &= \Delta t \left[(f_j^{n+1} - \tilde{f}_j) \left(\bar{a}_j - h \sum_{i=1}^N \bar{b}_{ji} \tilde{f}_i \right) - h f_j^{n+1} \sum_{i=1}^N \bar{b}_{ji} g_i^n \right]. \end{aligned}$$

Hence,

$$\begin{aligned} \|g^{n+1} - g^n\| &\leq \Delta t \|s[\tilde{f}]\|_{L^\infty} \|g^{n+1}\| + \Delta t h \|f^{n+1}\|_\infty \|B\|_2 \|g^n\| \\ &\leq \Delta t (\|s[\tilde{f}]\|_{L^\infty} + \|f^{n+1}\|_1 \|B\|_2) \|g^{n+1}\| + \Delta t \|f^{n+1}\|_1 \|B\|_2 \|g^{n+1} - g^n\| \\ &\leq c_1 \Delta t \|g^{n+1}\| + c_2 \Delta t \|g^{n+1} - g^n\|, \end{aligned}$$

where in virtue of Lemma 3.1,

$$(3.33) \quad c_2 = \max\left\{ \|f^0\|_1, \frac{\|a\|_{L^\infty}}{b_f} \right\} \|B\|_2, \quad c_1 = \|a\|_{L^\infty} + \|b\|_{L^\infty} \|\tilde{f}\|_1 + c_2.$$

This has proved (3.32) with

$$C^* = \frac{c_1}{1 - c_2 \Delta t}$$

for

$$(3.34) \quad \Delta t < \frac{1}{c_2}.$$

Substituting (3.32) into (3.31) and using $\lambda_{\min}\|g^{n+1}\|^2 \leq g^{n+1} \cdot Bg^{n+1}$, we have

$$\begin{aligned} F^{n+1} - F^n &\leq -\Delta th^2 [g^{n+1} \cdot Bg^{n+1} - \Delta t C^* \|B\|_2 \|g^{n+1}\|^2] + \Delta th \sum_{j=1}^N s_j [\tilde{f}] f_j^{n+1} \\ &\leq -\frac{1}{2} \Delta th^2 g^{n+1} \cdot Bg^{n+1} + \Delta th \sum_{j=1}^N s_j [\tilde{f}] f_j^{n+1}, \end{aligned}$$

as long as $\Delta t \leq \frac{\lambda_{\min}}{2C^*\|B\|_2} = \frac{\lambda_{\min}}{2C^*\lambda_{\max}}$, that is

$$(3.35) \quad \Delta t \leq \frac{\lambda_{\min}}{2c_1\lambda_{\max} + c_2\lambda_{\min}}.$$

We proceed

$$\begin{aligned} F^{n+1} - F^n &\leq -\Delta th \left(\frac{f^{n+1} \cdot Bf^{n+1}}{2} h - f^{n+1} \cdot B\tilde{f}h + \frac{\tilde{f} \cdot B\tilde{f}}{2} h - a \cdot f^{n+1} + f^{n+1} \cdot B\tilde{f}h \right) \\ (3.36) \quad &= -\Delta t [H(f^{n+1}) - H(\tilde{f})]. \end{aligned}$$

We next prove (3.26). Using the fact that B is symmetric and scheme (3.1), we calculate

$$\begin{aligned} H(f^{n+1}) - H(f^n) &= \left[\frac{1}{2} (f^{n+1} - f^n)^T B (f^{n+1} + f^n) h^2 - a^T (f^{n+1} - f^n) h \right] \\ &= - \sum_{j=1}^N \left[(f_j^{n+1} - f_j^n) \left(\bar{a}_j - h \sum_{i=1}^N \bar{b}_{ji} \frac{f_i^{n+1} + f_i^n}{2} \right) \right] h \\ &= -\Delta th \sum_{j=1}^N f_j^{n+1} (\bar{a}_j - h \sum_{i=1}^N \bar{b}_{ji} f_i^n)^2 \\ &\quad + \Delta th^2 \sum_{j=1}^N \left[f_j^{n+1} (\bar{a}_j - h \sum_{i=1}^N \bar{b}_{ji} f_i^n) \sum_{k=1}^N \bar{b}_{jk} \frac{f_k^{n+1} - f_k^n}{2} \right] \\ &=: -T_1 + T_2. \end{aligned}$$

Since $T_1 \geq 0$, we only need to show $T_2 \leq C\Delta t T_1$ for suitably small Δt . Using scheme (3.1), we obtain

$$\begin{aligned} T_2 &= \frac{(\Delta t)^2 h^2}{2} \sum_{j,k=1}^N \bar{b}_{jk} f_k^{n+1} f_j^{n+1} (\bar{a}_j - h \sum_{i=1}^N \bar{b}_{ji} f_i^n) (\bar{a}_k - h \sum_{i=1}^N \bar{b}_{ki} f_i^n) \\ (3.37) \quad &\leq \frac{(\Delta t)^2 h^2}{2} \sum_{j,k=1}^N \bar{b}_{jk} f_k^{n+1} f_j^{n+1} (\bar{a}_j - h \sum_{i=1}^N \bar{b}_{ji} f_i^n)^2 \\ &\leq \frac{(\Delta t)^2 h^2}{2} (\max_j \sum_{k=1}^N \bar{b}_{jk} f_k^{n+1}) \sum_{j=1}^N f_j^{n+1} (\bar{a}_j - h \sum_{i=1}^N \bar{b}_{ji} f_i^n)^2, \end{aligned}$$

where we have used the symmetry of B in the second inequality. Note that

$$(3.38) \quad h \max_j \sum_{k=1}^N \bar{b}_{jk} f_k^{n+1} \leq \|b\|_{L^\infty} \|f^{n+1}\|_1.$$

Then

$$T_2 \leq \frac{\Delta t}{2} \|b\|_{L^\infty} \|f^{n+1}\|_1 T_1.$$

Thus,

$$H(f^{n+1}) - H(f^n) \leq -(1 - \frac{\Delta t}{2} \|b\|_{L^\infty} \|f^{n+1}\|_1) T_1 \leq 0,$$

if

$$(3.39) \quad \Delta t \leq \frac{2\|B\|_2}{c_2\|b\|_{L^\infty}}.$$

(3.26) is established. Combining (3.30), (3.34), (3.35) and (3.39), we have the restriction (3.24). \square

Remark 3.1. *When $0 < b_f \leq \bar{b}_{ji}$ in (1.7) is weakened to $\bar{b}_{ji} \geq 0$, we are still able to show the same convergence rate, but with a different time step restriction.*

A more precise statement is as follows

Theorem 3.3. *Let $\bar{b}_{ji} \geq 0$ and other assumptions remain the same as those in Theorem 3.2, then*

$$(3.40) \quad \|f^n - \tilde{f}\|_b^2 \leq \frac{2F^0}{n\Delta t},$$

provided that

$$(3.41) \quad \Delta t \leq \min \left\{ \frac{\lambda_{\min}}{\lambda_{\max}(2C_1 + 2\lambda_{\max}S(F^0) + \lambda_{\min}S(F^0))}, \frac{h}{\|b\|_{L^\infty}S(F^0)} \right\}.$$

Proof. We use the fact (refer to the proof in [12, Theorem 3.1]) that there exists a non-decreasing, positive function S such that

$$(3.42) \quad h\|f^n\|_\infty \leq S(F^n),$$

and F^n is decreasing in n . Thus, $c_2 = S(F^0)\|B\|_2$ in (3.33). Combining this with estimate

$$h \max_j \sum_{k=1}^N \bar{b}_{jk} f_k^{n+1} \leq 2\|b\|_{L^\infty} \|f^n\|_\infty \leq 2\|b\|_{L^\infty} h^{-1} S(F^0)$$

instead of (3.38). The remaining proof follows the same strategy as in the proof of Theorem 3.2, then (3.40) is established if condition (3.41) holds. \square

4. CONCLUDING REMARKS

In this work, we have investigated the discrete dynamics of an integro-differential model that describes the evolution of a population structured with respect to a continuous trait. The discrete model considered is the entropy satisfying finite volume scheme proposed in [12]. Several time-asymptotic convergence rates towards the discrete evolutionary stable distribution (ESD) are established. More precisely, we have obtained the following results:

- For the discrete ESD satisfying a strict sign condition, we have established the exponential convergence rate of numerical solutions towards such a strict ESD for both the semi-discrete scheme and the fully discrete scheme. However, the convergence rate is typically mesh dependent, as a similar result is not expected for the continuous model.
- For general discrete ESD, we proved that numerical solutions of the fully discrete scheme converge towards the discrete ESD at an rate $1/n$, which is faster than the rate $O(\log t/t)$ obtained in [11] for the continuous model.

These results, proved for the one dimensional case, are expected to hold for arbitrary dimensions.

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REFERENCES

- [1] Barabás G, Meszzena G (2009) When the exception becomes the rule: the disappearance of limiting similarity in the Lotka-Volterra model. *J Math Biol* 258(1):89-94.
- [2] Barles G, Mirrahimi S, Perthame B (2009) Concentration in Lotka-Volterra parabolic or integral equations: a general convergence result. *Methods and Applications of Analysis* 16(3):321-340.
- [3] Brger R (2000) *The mathematical theory of selection, recombination and mutation*. Wiley, New York
- [4] Champagnat N, Ferrière R, Méléard S (2006) Unifying evolutionary dynamics: from individual stochastic processes to macroscopic models. *Theor Popul Biol* 69:297-321.
- [5] Champagnat N, Ferrière R, Méléard S (2008) From individual stochastic processes to macroscopic models in adaptive evolution. *Stoch Models* 24(1):2-44.
- [6] Cohen Y, Galiano G (2013) Evolutionary distributions and competition by way of reaction-diffusion and by way of convolution. *Bulletin of Mathematical Biology*, 75(12):2305-2323.
- [7] Desvillettes L, Jabin PE, Mischler S, Raoul G (2008) On selection dynamics for continuous structured populations. *Commun Math Sci* 6(3):729-747.
- [8] Dieckmann U, Law R (1996) The dynamical theory of coevolution: a derivation from a stochastic ecological processes. *J Math Biol* 34:579-612.
- [9] Doebeli M, Blok HJ, Leimar O, Dieckmann U (2007) Multimodal pattern formation in phenotype distributions of sexual populations. *Proc R Soc B* 274:347-357.
- [10] Genieys S, Volpert V, Auger P (2006) Pattern and waves for a model in population dynamics with nonlocal consumption of resources. *Math Model Nat Phenom* 1:65-82.
- [11] Jabin PE, Raoul G (2011) On selection dynamics for competitive interactions. *J Math Biol* 63(3):493-517.
- [12] Liu HL, Cai WL, Su N (2014) Entropy satisfying schemes for computing selection dynamics in competitive interactions. Submitted.
- [13] Lorz A, Mirrahimi S, Perthame B (2011) Dirac mass dynamics in multidimensional nonlocal parabolic equations. *Communications in Partial Differential Equations*, 36(6):1071-1098.
- [14] Mirrahimi S, Perthame B, Bouin E, Millien P (2011) Population formulation of adaptative meso-evolution: theory and numerics. *Mathematics and Biosciences in Interaction* 159-174.
- [15] Mirrahimi S, Perthame B, Wakano JY (2012) Direct competition results from strong competitor for limited resource (preprint).
- [16] Perthame B, Barles G (2008) Dirac concentrations in Lotka-Volterra parabolic PDEs. *Indiana Univ Math J* 57(7):3275-3301.
- [17] Roughgarden J (1979) *Theory of population genetics and evolutionary ecology*. Macmillan, New York
- [18] Raoul G (2011) Long time evolution of populations under selection and vanishing mutations. *Acta Appl Math* 114(1-2):1-14.
- [19] Raoul G (2012) Local stability of evolutionary attractors for continuous structured populations. *Monatsh Math* 165(1):117-144.

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