

On the Convergence of Difference Approximations to Scalar Conservation Laws*

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Abstract. We present a unified treatment of explicit in time, two-level, second-order resolution (SOR), total-variation diminishing (TVD), approximations to scalar conservation laws. The schemes are assumed only to have conservation form and incremental form. We introduce a modified flux and a viscosity coefficient and obtain results in terms of the latter. The existence of a cell entropy inequality is discussed and such an equality for all entropies is shown to imply that the scheme is an E scheme on monotone (actually more general) data, hence at most only first-order accurate in general. Convergence for TVD-SOR schemes approximating convex or concave conservation laws is shown by enforcing a single discrete entropy inequality.

Introduction. Recently there has been an enormous amount of activity related to the construction and analysis of finite-difference approximations which approximate nonlinear hyperbolic systems of conservation laws and which are supposed to have the following properties:

- (1) Limit solutions which satisfy a geometric and/or analytic entropy condition.
- (2) A bound on the variation of the approximate solutions at least in the scalar and linear systems case. This bound is such as to imply the absence of spurious oscillations in the approximate solutions.
- (3) At least second-order accuracy in regions of smoothness, except for certain isolated points as described below.

Some examples of the successful computational consequences of this activity can be found in the proceedings of the sixth AIAA Computational Fluid Dynamics Conference, and elsewhere; see, e.g., the bibliography in [21].

Some of the earliest work in the design of schemes having properties (2) and (3) above was done by Van Leer [27], [28]. There he introduced the concepts of flux limiters and higher-order Riemann solvers. Recently, Harten [10], [11] obtained conditions which he showed to be compatible with second-order accuracy, and which guarantee that a scalar one-dimensional scheme is TVD, that is, total-variation diminishing. He constructed a scheme having that property and formally extended it to systems using a field-by-field limiter and Roe's decomposition [22].

We would also like to mention the work of Boris and Book [1] concerning FCT schemes. They also used flux limiters to suppress oscillations in their schemes.

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Harten's construction in [10] was done first for a fully discrete, explicit in time approximation. P. K. Sweby [24] has investigated the properties of various limiters in this context.

We shall use the term "high resolution scheme" to mean a formal extension to systems via a field-by-field decomposition of a scalar, higher than first-order accurate, variation diminishing scheme. These schemes do not, in general, satisfy the entropy condition; e.g., expansion shocks exist as stable solutions of high resolution schemes based on Roe's (unmodified) scheme. In [20], Osher's decomposition and certain flux limiters were used to prove that limit solutions to a class of semidiscrete, time continuous high resolution schemes do satisfy the entropy condition for hyperbolic systems of conservation laws. Convergence of other classes of semidiscrete, time continuous, high resolution approximations to scalar convex conservation laws was proven in [19] and [20].

A systematic recipe for constructing semidiscrete high resolution schemes whose formal accuracy is higher than two (away from the isolated points) using a minimal band width, was presented in [21].

In the present paper we are considering two-level, (and for simplicity only) explicit finite-difference approximations to a scalar conservative law having two properties:

- (1) Lax-Wendroff conservative form [15], and
- (2) Roe's incremental form [22].

In Section 1 we introduce a modified flux and a viscosity coefficient, both grid-dependent quantities, with which we set up and advocate a third form of these schemes, the viscosity form. A sufficient TVD criterion is then presented in terms of this form.

In Section 2 we consider schemes which have a 3-point stencil. In this case the modified flux coincides with the original (differential equation's) flux; accuracy is consequently limited to first order.

Thus, in Section 3, we consider wider stencils. A comprehensive framework of second-order resolution (SOR), total-variation diminishing (TVD) schemes is presented. The viscous form of the scheme again plays a key role here.

Two concrete examples for the construction of SOR-TVD schemes according to the above guidelines are detailed in Section 4; both approaches use the viscous form of the underlying 3-point TVD schemes. Special attention is called to the second ("piecewise constant viscosity modification") recipe in Corollary 4.9, which results in an easily implemented and highly attractive two-step formulation.

These four sections comprise Part I. In our attempt at a unified treatment of this subject, we derived (and in some cases rederived) some notable conclusions:

- (i) Three-point TVD schemes are, at most, first-order accurate.
- (ii) Accuracy at nonsonic critical points is limited to first order.
- (iii) Our piecewise linear flux correction (Theorem 4.3) extends Harten's recipe [10], under a CFL of limitation $2/3$, (Remark 4.10), in agreement with that found by Sweby in [24].
- (iv) Our recipe for converting three-point TVD schemes to five-point SOR-TVD schemes seems quite general and attractive from both a computational and a theoretical point of view. The latter point of view will be used in Part II.

The discussion concerning a cell entropy inequality begins in Section 5, using the canonical Godunov scheme. Section 6 contains the heart of our discussion. The *grid-dependent* modified flux of Section 1 is extended to a modified flux function defined on the *intervals connecting the grid values*. The definition involves a piecewise linear function (i.e., the “double wing” seen in Figures 6.4) and the remark made after the statement of Theorem 6.11, which takes into account the presence or absence of critical points.

The schemes under consideration are then expressed as convex combinations of Riemann problem solvers, i.e., of Godunov-type schemes, using the above modified flux functions. We obtain an entropy in cell inequality involving a familiar (integral) residual term.

Section 7 discusses the significance of E fluxes (introduced in [18], cf. also [26]). We show that for *any* approximate entropy in cell inequality to be valid for *all* intervals, the underlying scheme must have an E flux, hence be at most first-order accurate (Remark 7.2). Conversely, the existence of an E flux implies the nonpositivity of the residual term mentioned above, hence implies a general cell entropy inequality. To obtain convergence for SOR schemes, we give up the requirement that *all* discrete entropy inequalities be valid. This limits our convergence proof to convex (or concave) conservation laws, as in [19], [20].

In Section 8 we treat SOR-TVD schemes satisfying a single quadratic cell entropy inequality. Here we use the viscosity modification recipe of Section 4. A cell entropy inequality follows by estimating the residual integral presented in Section 6. The rather delicate tuning of the modified flux so as to comply with the three criteria SOR, TVD, and entropy inequality, leads us to a wide class of explicit in time SOR-TVD schemes (again only for the convex or concave scalar case).

The main convergence results are contained in Theorems 8.4 and 8.5.

PART I. TOTAL-VARIATION DIMINISHING SCHEMES

1. The Modified Flux and the Numerical Viscosity Coefficient. We study two-step difference schemes of the form

$$(1.1) \quad v_\nu(t+k) = H(v_{\nu-p}(t), \dots, v_{\nu+p}(t); f, \lambda),$$

serving as consistent approximations to the scalar conservation law

$$(1.2) \quad \frac{\partial u}{\partial t} + \frac{\partial f}{\partial x} = 0.$$

Here, $v_\nu(t) \equiv v(x_\nu, t)$ denotes the approximation value at the grid point ($x_\nu \equiv \nu \Delta x, t$), $k \equiv \Delta t$ and Δx are, respectively, the temporal and spatial mesh size with fixed mesh ratio $\lambda = k/\Delta x$, and p a natural number.

We postulate

Assumption 1.1. (i) The scheme (1.1) admits a *conservative form* in the sense of Lax-Wendroff [15],

$$(1.3) \quad H(v_{\nu-p}, \dots, v_{\nu+p}; f, \lambda) = v_\nu - \lambda(h_{\nu+1/2} - h_{\nu-1/2}),$$

where $h_{\nu+1/2}$ stands for the Lipschitz continuous *numerical flux*

$$h_{\nu+1/2} \equiv h(v_{\nu-p+1}, \dots, v_{\nu+p}; f, \lambda),$$

consistent with the differential one,

$$h(w, w, \dots, w; f, \lambda) = f(w).$$

(ii) The scheme (1.1) can also be written in an *incremental form*

$$(1.4) \quad H(v_{\nu-p}, \dots, v_{\nu+p}; f, \lambda) = v_{\nu} + C_{\nu+1/2}^+ \Delta v_{\nu+1/2} - C_{\nu-1/2}^- \Delta v_{\nu-1/2},$$

where we use the standard notation

$$\Delta v_{\nu+1/2} = v_{\nu+1} - v_{\nu}.$$

Equating the right-hand sides of (1.3) and (1.4), rearranging and dividing by λ , we find

$$(1.5a) \quad h_{\nu+1/2} + \frac{1}{\lambda} C_{\nu+1/2}^+ \Delta v_{\nu+1/2} = h_{\nu-1/2} + \frac{1}{\lambda} C_{\nu-1/2}^- \Delta v_{\nu-1/2}.$$

We term the equated grid-dependent quantities in (1.5a) as the *modified flux* associated with scheme (1.1):

$$(1.5b) \quad \begin{aligned} g_{\nu} &= h_{\nu \pm 1/2} + \frac{1}{\lambda} C_{\nu \pm 1/2}^{\pm} \Delta v_{\nu \pm 1/2} \\ &= \frac{1}{2} \left[h_{\nu-1/2} + h_{\nu+1/2} + \frac{1}{\lambda} C_{\nu-1/2}^- \Delta v_{\nu-1/2} + \frac{1}{\lambda} C_{\nu+1/2}^+ \Delta v_{\nu+1/2} \right]. \end{aligned}$$

As in [25, Section 2], we now use the consistency relation

$$(1.6a) \quad C_{\nu+1/2}^- - C_{\nu+1/2}^+ = \lambda \frac{\Delta g_{\nu+1/2}}{\Delta v_{\nu+1/2}}$$

to conclude that with a *given modified flux*, g_{ν} , there is only *one* degree of freedom in setting up the recipe of the difference scheme (1.1): This will be manifested in terms of the quantity,

$$(1.6b) \quad Q_{\nu+1/2} = C_{\nu+1/2}^- + C_{\nu+1/2}^+.$$

Indeed, by averaging (1.3) and (1.4) we find

$$\begin{aligned} v_{\nu}(t+k) &= v_{\nu}(t) - \frac{\lambda}{2} (h_{\nu+1/2} - h_{\nu-1/2}) \\ &\quad + \frac{1}{2} \left(C_{\nu+1/2}^+ \Delta v_{\nu+1/2} - C_{\nu+1/2}^- \Delta v_{\nu-1/2} \right); \end{aligned}$$

adding and subtracting $\frac{1}{2}(C_{\nu+1/2}^- \Delta v_{\nu+1/2} - C_{\nu-1/2}^+ \Delta v_{\nu-1/2})$ on the right-hand side, we arrive at

$$\begin{aligned} v_{\nu}(t+k) &= v_{\nu}(t) - \frac{\lambda}{2} \left[\left(h_{\nu+1/2} + \frac{1}{\lambda} C_{\nu+1/2}^- \Delta v_{\nu+1/2} \right) \right. \\ &\quad \left. - \left(h_{\nu-1/2} + \frac{1}{\lambda} C_{\nu-1/2}^+ \Delta v_{\nu-1/2} \right) \right] \\ &\quad + \frac{1}{2} \left[\left(C_{\nu+1/2}^- + C_{\nu+1/2}^+ \right) \Delta v_{\nu+1/2} - \left(C_{\nu-1/2}^- + C_{\nu-1/2}^+ \right) \Delta v_{\nu-1/2} \right]. \end{aligned}$$

In view of (1.5b) and (1.6b), the scheme (1.1) is thus finally recast into its *viscous form*

$$(1.7a) \quad v_{\nu}(t+k) = v_{\nu}(t) - \frac{\lambda}{2} (g_{\nu+1} - g_{\nu-1}) + \frac{1}{2} [\Delta(Q_{\nu-1/2} \Delta v_{\nu-1/2}(t))],$$

expressed in terms of the *numerical viscosity coefficient* [10], [26]

$$(1.7b) \quad Q_{\nu+1/2} = C_{\nu+1/2}^- + C_{\nu+1/2}^+ = \lambda \frac{g_\nu + g_{\nu+1} - 2h_{\nu+1/2}}{\Delta v_{\nu+1/2}}.$$

Let $\text{TV}[v(t)] \equiv \sum_\nu |v_{\nu+1}(t) - v_\nu(t)|$ denote the total variation of the computed solution at time level t ; the following lemma provides us with a sufficient criterion for the scheme (1.1) to be total-variation diminishing (TVD), in the spirit of [10].

LEMMA 1.2. (*Total-variation diminishing.*) *The scheme (1.1) has a diminishing total variation,*

$$\text{TV}[v(t+k)] \leq \text{TV}[v(t)],$$

provided its numerical viscosity coefficient $Q_{\nu+1/2}$ satisfies

$$(1.8) \quad \lambda \left| \frac{\Delta g_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \leq Q_{\nu+1/2} \leq 1.$$

Proof. By averaging (1.6a) and (1.6b) we find

$$(1.9) \quad C_{\nu+1/2}^\pm = \frac{1}{2} \left(Q_{\nu+1/2} \mp \lambda \frac{\Delta g_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right).$$

In view of these last relations, the inequalities (1.8) boil down to

$$(1.10) \quad C_{\nu+1/2}^+ \geq 0, \quad C_{\nu+1/2}^- \geq 0, \quad 1 - C_{\nu+1/2}^- - C_{\nu+1/2}^+ \geq 0,$$

and TVD follows along the lines of [10]. $\square$

2. Three-Point TVD Schemes. In the case of 3-point schemes, $p = 1$, it was shown in [25, Lemma 2.1] that there exists *one and only one* incremental form (1.4) whose incremental coefficients are

$$(2.1) \quad C_{\nu+1/2}^+ = \lambda \frac{f_\nu - h_{\nu+1/2}}{\Delta v_{\nu+1/2}}, \quad C_{\nu+1/2}^- = \lambda \frac{f_{\nu+1} - h_{\nu+1/2}}{\Delta v_{\nu+1/2}}.$$

Inserted into (1.5b) we find that the modified flux in this case coincides with the original one, $g_\nu = f_\nu$, and the scheme is therefore completely determined by its numerical viscosity coefficient, $Q_{\nu+1/2}$,

$$(2.2) \quad v_\nu(t+k) = v_\nu(t) - \frac{\lambda}{2} (f_{\nu+1} - f_{\nu-1}) + \frac{1}{2} [\Delta(Q_{\nu-1/2} \Delta v_{\nu-1/2})].$$

Abbreviating $a(w) = df(w)/dw$, the following notation will be used throughout:

$$(2.3) \quad a_{\nu+1/2} = \begin{cases} \frac{\Delta f_{\nu+1/2}}{\Delta v_{\nu+1/2}} & \text{if } \Delta v_{\nu+1/2} \neq 0, \\ a\left(\frac{v_\nu + v_{\nu+1}}{2}\right) & \text{if } \Delta v_{\nu+1/2} = 0. \end{cases}$$

The TVD constraint (1.8) now reads

$$(2.4) \quad \lambda |a_{\nu+1/2}| \leq Q_{\nu+1/2} \leq 1,$$

and the following list quotes the schemes in this 3-point TVD category most frequently referred to [6], [7], [16], [22]:

$$(2.5a) \quad Q_{\nu+1/2}^{\text{LF}} = 1 \quad (\text{Lax-Friedrichs scheme}),$$

$$(2.5b) \quad Q_{\nu+1/2}^{\text{EO}} = \frac{\lambda}{\Delta v_{\nu+1/2}} \int_{v_\nu}^{v_{\nu+1}} |a(v)| dv \quad (\text{Engquist-Osher scheme}),$$

$$(2.5c) \quad Q_{\nu+1/2}^{\text{G}} = \lambda \max_{(v-v_\nu) \cdot (v-v_{\nu+1}) \leq 0} \frac{f_\nu + f_{\nu+1} - 2f(v)}{\Delta v_{\nu+1/2}} \quad (\text{Godunov scheme}),$$

$$(2.5d) \quad Q_{\nu+1/2}^{\text{R}} = \lambda \cdot |a_{\nu+1/2}| \quad (\text{Roe-Murman scheme}).$$

In fact, in the special case under consideration of 3-point schemes, condition (2.4) is *necessary* as well as sufficient for TVD; see [25, Corollary 2.3]. Hence, the following 3-point schemes are *not* TVD ones:

$$(2.6a) \quad Q_{\nu+1/2}^{\text{LW}} = (\lambda a_{\nu+1/2})^2 \quad (\text{Lax-Wendroff scheme}),$$

$$(2.6b) \quad Q_{\nu+1/2}^{\text{E}} = 0 \quad (\text{forward Euler scheme}).$$

A special significance is attached to the Lax-Wendroff scheme

$$(2.7) \quad v_\nu(t+k) = v_\nu(t) - \frac{\lambda}{2}(f_{\nu+1} - f_{\nu-1}) + \frac{1}{2} [\Delta ((\lambda a_{\nu-1/2})^2 \Delta v_{\nu-1/2})];$$

when seeking second-order accurate schemes, the further limitation placed on $Q_{\nu+1/2}$ in the 3-point case singles out the Lax-Wendroff choice, (2.6a). In view of the above, we therefore conclude that *3-point TVD schemes are at most first-order accurate* [10], [18], [25].

At this point we are widening our discussion to include schemes whose stencil occupies more than three points: The further freedom in setting up the modified flux g_ν in this case will enable us to achieve higher (than one) degree of accuracy.

3. Second-Order Resolution Schemes. We start by rewriting the modified flux, g_ν , in terms of the *correction* $\tilde{g}_\nu$ to the original one,

$$(3.1) \quad g_\nu = f_\nu + \frac{1}{\lambda} \tilde{g}_\nu.$$

Inserted into (1.7a), our scheme now reads

$$(3.2) \quad \begin{aligned} v_\nu(t+k) &= v_\nu(t) - \frac{\lambda}{2} \left[\left(f_{\nu+1} + \frac{1}{\lambda} \tilde{g}_{\nu+1} \right) - \left(f_{\nu-1} + \frac{1}{\lambda} \tilde{g}_{\nu-1} \right) \right] \\ &\quad + \frac{1}{2} [\Delta(Q_{\nu-1/2} \Delta v_{\nu-1/2})] \\ &= v_\nu(t) - \frac{\lambda}{2} (f_{\nu+1} - f_{\nu-1}) + \frac{1}{2} [\Delta(Q_{\nu-1/2} \Delta v_{\nu-1/2})], \end{aligned}$$

where $Q_{\nu+1/2}$ stands for the *modified viscosity coefficient* given by

$$(3.3) \quad Q_{\nu+1/2} = Q_{\nu+1/2} - \frac{\tilde{g}_\nu + \tilde{g}_{\nu+1}}{\Delta v_{\nu+1/2}} = \lambda \left[\frac{f_\nu + f_{\nu+1} - 2h_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right].$$

In other words, *we rewrite our scheme with respect to the original flux, f_ν , modifying the viscosity coefficient instead.*

Remark 3.1. The viscosity coefficient discussed in [10] corresponds to the *modified* one in (3.3); indeed, in the 3-point case treated in [25], where $\tilde{g}_\nu = \lambda(g_\nu - f_\nu) = 0$, it coincides with (regular) viscosity in (1.7b).

Comparison of (3.2) and the Lax-Wendroff scheme (2.7) leads us to the following second-order accuracy requirement:

$$(3.4a) \quad \mathbf{Q}_{\nu+1/2} \equiv Q_{\nu+1/2} - \frac{\tilde{g}_\nu + \tilde{g}_{\nu+1}}{\Delta v_{\nu+1/2}} = \lambda^2(a^2)_{\nu+1/2};$$

the deliberately vague notation of $\lambda^2(a^2)_{\nu+1/2}$ on the right stands for $Q_{\nu+1/2}^{\text{LW}} = \lambda^2(a_{\nu+1/2})^2$ modulo first-order errors, i.e.,

$$(3.4b) \quad \lambda^2(a^2)_{\nu+1/2} \equiv \lambda^2(a_{\nu+1/2})^2 + O(|\Delta v|)_{\nu+1/2}.$$

Together with the TVD constraint (1.8), we finally arrive at a general description of second-order TVD schemes, which is summarized in the following

LEMMA 3.2. (*Second-order TVD schemes.*) *Consider the difference scheme (1.1) written in its viscous form*

$$(3.5a) \quad v_\nu(t+k) = v_\nu(t) - \frac{\lambda}{2}(g_{\nu+1} - g_{\nu-1}) + \frac{1}{2} [\Delta(Q_{\nu-1/2} \Delta v_{\nu-1/2})],$$

with modified flux, g_ν , and viscosity coefficient, $Q_{\nu+1/2}$, given respectively by

$$(3.5b) \quad g_\nu = f_\nu + \frac{1}{\lambda} \tilde{g}_\nu,$$

$$(3.5c) \quad Q_{\nu+1/2} = \lambda^2(a^2)_{\nu+1/2} + \frac{\tilde{g}_\nu + \tilde{g}_{\nu+1}}{\Delta v_{\nu+1/2}}.$$

The scheme (3.5) is second-order TVD, provided the correction terms $\tilde{g}_\nu$ and $\lambda^2(a^2)_{\nu+1/2}$ on the right-hand side of (3.5b–c) satisfy the following two requirements of

(i) *second-order accuracy:*

$$(3.6a) \quad \lambda^2(a^2)_{\nu+1/2} = (\lambda a_{\nu+1/2})^2 + O(|\Delta v|)_{\nu+1/2};$$

(ii) *total-variation diminishing:*

$$(3.6b) \quad \left| \lambda a_{\nu+1/2} + \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \leq \lambda^2(a^2)_{\nu+1/2} + \frac{\tilde{g}_\nu + \tilde{g}_{\nu+1}}{\Delta v_{\nu+1/2}} \leq 1.$$

As we shall see below, one cannot satisfy both requirements (3.6a) and (3.6b) at the nonsonic critical grid values v_ν where $\Delta v_{\nu-1/2} \cdot \Delta v_{\nu+1/2} < 0 \neq a(v_\nu)$, and, therefore, second-order accuracy must be given up at these values. Difference schemes with (formal) second-order accuracy at all but those critical grid values are classified as having second-order resolution, after Harten [10]. Thanks to the TVD property, no new such first-order accurate critical grid values are added during the computation (since the scheme is monotonicity-preserving [10], [25]), and the *overall* second-order accuracy does not seem to be degraded in this case, at least in the L^1 -norm.

There are various approaches to the construction of second-order resolution (*abbreviated hereafter SOR*) schemes, the main three of which are the following:

(i) The *MUSCL* approach [27], [19] directly modifies the first-order *numerical flux*, $h_{\nu+1/2}$, using a Godunov-like solver for a limited slope piecewise linear grid data.

(ii) The *modified equation* approach [10], [11] directly modifies the *flux* $f(v_\nu)$, based on considerations of the modified equation associated with a first-order scheme, so that the limited modified flux constructed guarantees second-order TVD resolution.

(iii) The *flux limiter* approach [24] directly modifies the first-order *incremental* coefficients, $C_{\nu+1/2}^\pm$, using a class of flux limiters which preserve TVD and maintain second-order resolution; in particular, such flux limiters can be chosen to interpret MUSCL and modified flux-type schemes due to Van Leer, Roe, Harten, Chakravarthy, Osher, and others; see [24].

All the above approaches can be entertained, of course, within the general framework provided in Lemma 3.2. A common “limiting” feature shared by all of these approaches can be directly derived from the TVD constraints in that lemma: The grid values $\tilde{g}_\nu$ in (3.5) are expected to form a first-order “smooth” grid function correction, at least at the generic noncritical zones; since $\lambda^2(a^2)_{\nu+1/2}$ is determined up to first-order perturbations, see (3.6a), we can absorb such a perturbation, $\pm\Delta\tilde{g}_{\nu+1/2}/\Delta v_{\nu+1/2} = O(|\Delta v|_{\nu+1/2})$, into the middle term of (3.6b), and the right-hand inequality now reads

$$(3.7a) \quad \lambda^2(a^2)_{\nu+1/2} + 2\frac{\tilde{g}_{\nu+1}}{\Delta v_{\nu+1/2}} \leq 1, \quad \lambda^2(a^2)_{\nu+1/2} + 2\frac{\tilde{g}_\nu}{\Delta v_{\nu+1/2}} \leq 1.$$

The left-hand inequality in (3.6b) gives us

$$(3.7b) \quad \begin{aligned} \lambda a_{\nu+1/2} - \lambda^2(a^2)_{\nu+1/2} &\leq 2\frac{\tilde{g}_\nu}{\Delta v_{\nu+1/2}}, \\ -\lambda a_{\nu+1/2} - \lambda^2(a^2)_{\nu+1/2} &\leq 2\frac{\tilde{g}_{\nu+1}}{\Delta v_{\nu+1/2}}. \end{aligned}$$

We shift indices in the first and fourth inequalities in (3.7). The resulting inequalities boil down to the following inequality to be satisfied by TVD schemes:

$$(3.7c) \quad \frac{1}{2}(\pm\lambda a_{\nu\pm 1/2} - \lambda^2(a^2)_{\nu\pm 1/2}) \leq \frac{\tilde{g}_\nu}{\Delta v_{\nu\pm 1/2}} \leq \frac{1}{2}(1 - \lambda^2(a^2)_{\nu\pm 1/2}).$$

Hence, $\tilde{g}_\nu/\Delta v_{\nu\pm 1/2}$ *must* lie between the above two bounds, *both determined up to first-order perturbations*. In particular, in view of the second-order accuracy requirement in (3.6a), the sign of the parenthesis of the left of (3.7c) is determined by that of the first term $\pm\lambda a_{\nu\pm 1/2}$; further, if we are at a nonsonic value, $a(v_\nu) \neq 0$, then this first term cannot be absorbed as a first-order perturbation of the second one, and (3.7c) *requires* the middle terms $\tilde{g}_\nu/\Delta v_{\nu\pm 1/2}$ to be positive. *This implies that the second-order accuracy must be given up at the nonsonic critical values v_ν where $\Delta v_{\nu-1/2} \cdot \Delta v_{\nu+1/2} < 0 \neq a(v_\nu)$, [25], [20].*

We now describe two specific recipes which convert arbitrary 3-point TVD schemes into second-order resolution ones, in the spirit of the above guidelines. Indeed, although the reasoning may be different, the various “conversion” recipes, including the two below, end up with difference schemes which bear close similarities to each other, as dictated by the framework provided in Lemma 3.2.

4. Two Recipes for the Construction of SOR-TVD Schemes. We start by identifying a 3-point scheme by its numerical viscosity coefficient $Q_{\nu+1/2} = Q(f; v_\nu, v_{\nu+1}, \lambda)$; it is considered here as a *functional* in the original flux $f(\cdot)$ with

the further dependence on the grid values $v_\nu, v_{\nu+1}$ and the mesh ratio λ . In view of the *necessity* of the TVD constraint on the left of (2.4), we may assume without loss of generality that the following *functional* inequality is obeyed:

$$(4.1) \quad \lambda \left| \frac{z(w_2) - z(w_1)}{w_2 - w_1} \right| \leq Q(z; w_1, w_2, \lambda).$$

Thus, the TVD requirement is reduced in this case to the CFL-like condition given on the right of (2.4),

$$(4.2) \quad Q_{\nu+1/2}(f) \equiv Q(f; v_\nu, v_{\nu+1}, \lambda) \leq 1.$$

Remark 4.1. The above description seems to exclude several difference schemes, where it is actually the inequality on the *left* of (2.4), $\lambda |\Delta f_{\nu+1/2} / \Delta v_{\nu+1/2}| \leq Q_{\nu+1/2}(f)$, which leads to the CFL limitation, as in the case, for example, of the Lax-Friedrichs scheme (2.5a) with $Q_{\nu+1/2}^{\text{LF}} = 1$. *Together with the TVD constraint* (2.4), however, the Lax-Friedrichs scheme, for example, can be equivalently represented by a numerical viscosity coefficient $Q_{\nu+1/2} = \text{Max}(1, \lambda |\Delta f_{\nu+1/2}| / |\Delta v_{\nu+1/2}|)$ which fits into our above interpretation. In general, an expansion of $Q(\cdot; \cdot, \cdot, \lambda)$ in powers of λ is called for, but we shall not elaborate on that here.

To construct an SOR-TVD scheme, the two first-order correction terms $\tilde{g}_\nu$ and $\lambda^2(a^2)_{\nu+1/2}$ are to be determined; see (3.5). Taking advantage of relation (3.5c), we shall also make use of the further (first-order) flexibility in determining the viscosity coefficient, $Q_{\nu+1/2}^{\text{SOR}}$.

Let us denote

$$(4.3a) \quad s_\nu = \frac{1}{2}(s_{\nu-1/2} + s_{\nu+1/2}), \quad s_{\nu+1/2} \equiv \text{sgn}(\Delta v_{\nu+1/2}),$$

and, in view of (3.7c), set the flux correction to be

$$(4.3b) \quad \tilde{g}_\nu = \frac{s_\nu}{2} \text{Min}([Q_{\nu\pm 1/2}(f) - \lambda^2(a^2)_{\nu\pm 1/2}] \cdot |\Delta v_{\nu\pm 1/2}|).$$

Here, the term $\lambda^2(a^2)_{\nu+1/2}$ is chosen so that

$$(4.3c) \quad \lambda^2(a^2)_{\nu+1/2} = (\lambda a_{\nu+1/2})^2 + O(|\Delta v|)_{\nu+1/2}, \quad \lambda^2(a^2)_{\nu+1/2} \leq Q_{\nu+1/2}(f)$$

with otherwise arbitrary first-order perturbation $O(|\Delta v|)_{\nu+1/2}$. The modified flux is then given by

$$(4.3d) \quad g_\nu = f(v_\nu) + \frac{1}{\lambda} \tilde{g}_\nu,$$

and we are considering the difference scheme whose viscous form is expressed in terms of that modified flux,

$$(4.3e) \quad v_\nu(t+k) = v_\nu(t) - \frac{\lambda}{2}(g_{\nu+1} - g_{\nu-1}) + \frac{1}{2}[\Delta(Q_{\nu-1/2}^{\text{SOR}} \Delta v_{\nu-1/2})].$$

We have

THEOREM 4.2. *The difference scheme (4.3) is SOR-TVD provided its viscosity coefficient, $Q_{\nu+1/2}^{\text{SOR}}$, satisfies*

$$(4.4a) \quad \left| \lambda a_{\nu+1/2} + \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \leq Q_{\nu+1/2}^{\text{SOR}} \leq 1,$$

$$(4.4b) \quad |Q_{\nu+1/2}^{\text{SOR}} - Q_{\nu+1/2}(f)| = O(|\Delta v|)_{\nu+1/2}.$$

Remark. We note that the definition of $(a^2)_{\nu+1/2}$ in (4.3c) still allows the flexibility of (some limited amount of) first-order perturbation, as long as $Q_{\nu+1/2}(f) - \lambda^2(a^2)_{\nu+1/2}$ remains nonnegative; the latter guarantees that the sign of $\tilde{g}_\nu$ agrees with $s_{\nu\pm 1/2}$.

Proof. The first condition, (4.4a) is nothing but the TVD requirement (3.6b); consult (1.8). Away from the critical values where $s_\nu = 0$, (4.4b) implies second-order accuracy; consult (3.4),

$$\begin{aligned} Q_{\nu+1/2}^{\text{SOR}} - \frac{\tilde{g}_\nu + \tilde{g}_{\nu+1}}{\Delta v_{\nu+1/2}} &= Q_{\nu+1/2}^{\text{SOR}} - [Q_{\nu+1/2}(f) - \lambda^2(a^2)_{\nu+1/2}] + O(|\Delta v|)_{\nu+1/2} \\ &= [Q_{\nu+1/2}^{\text{SOR}} - Q_{\nu+1/2}(f)] + \lambda^2(a_{\nu+1/2})^2 + O(|\Delta v|)_{\nu+1/2} \\ &= \lambda^2(a_{\nu+1/2})^2 + O(|\Delta v|)_{\nu+1/2}. \quad \square \end{aligned}$$

Two concrete choices for SOR viscosity will now be discussed.

As a building block for the first, we introduce the piecewise linear flux *correction*

$$(4.5a) \quad \tilde{g}_{\nu+1/2}(v) = \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}}(v - v_\nu) + \tilde{g}_\nu$$

and use the original 3-point viscosity *functional*, setting

$$(4.5b) \quad Q_{\nu+1/2}^{\text{SOR}} = Q \left(f(\cdot) + \frac{1}{\lambda} \tilde{g}_{\nu+1/2}(\cdot); v_\nu, v_{\nu+1}, \lambda \right).$$

Lipschitz continuity of the viscosity functional implies¹

$$(4.6) \quad \left| Q_{\nu+1/2}^{\text{SOR}} - Q_{\nu+1/2}(f) \right| \leq L \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right|,$$

and we end up with

THEOREM 4.3. *(Piecewise linear flux modification.) The difference scheme (4.3), (4.5) is SOR-TVD under the CFL-like condition*

$$(4.7) \quad Q_{\nu+1/2}(f) \leq \frac{2}{2 + L}.$$

Proof. The SOR requirement in (4.4b) is fulfilled because of (4.6), whose right-hand side is of order $(O|\Delta v|)_{\nu+1/2}$; the TVD requirement on the left of (4.4a) follows from (4.1), and we are left with the inequality on the right, requiring

$$(4.8) \quad Q_{\nu+1/2}^{\text{SOR}} \leq 1.$$

¹To be precise, the constant L stands for the Lipschitz constant times the (assumed finite) maximal bound of the ratios $|Q(\tilde{g}_{\nu+1/2}(\cdot); v_\nu, v_{\nu+1}, \lambda)|/|\Delta \tilde{g}_{\nu+1/2}/\Delta v_{\nu+1/2}|$.

Indeed, in view of (4.6) we have

$$(4.9a) \quad Q_{\nu+1/2}^{\text{SOR}} \leq Q_{\nu+1/2}(f) + L \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right|;$$

since $\tilde{g}_{\nu+1}$ and $\tilde{g}_{\nu}$ cannot have different signs by the choice of s_{ν} in (4.3a), we also find

$$(4.9b) \quad \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \leq \frac{\text{Max}(|\tilde{g}_{\nu}|, |\tilde{g}_{\nu+1}|)}{|\Delta v_{\nu+1/2}|} \leq \frac{1}{2} [Q_{\nu+1/2}(f) - \lambda^2 (a^2)_{\nu+1/2}].$$

The last two inequalities yield the crude bound

$$(4.10) \quad \begin{aligned} Q_{\nu+1/2}^{\text{SOR}} &\leq Q_{\nu+1/2}(f) + \frac{L}{2} [Q_{\nu+1/2}(f) - \lambda^2 (a^2)_{\nu+1/2}] \\ &\leq \left(1 + \frac{L}{2}\right) Q_{\nu+1/2}(f), \end{aligned}$$

and augmented with the assumed CFL-like condition (4.7), the TVD constraint (4.8) is now satisfied. $\square$

Remark 4.4. With the primary examples of Engquist-Osher, Godunov and Roe-Murman schemes in (2.5b–d), the Lipschitz constant L in (4.6) may be taken to be $L = 1$, and the CFL limit in (4.7) is found to equal (compare, e.g., [24, Section 3])

$$(4.11) \quad Q_{\nu+1/2}(f) \leq 2/3.$$

(Taking into account the negative term $-\lambda^2 (a^2)_{\nu+1/2}$ which was ignored in the crude bound of (4.10), the CFL limit may be further improved to be

$$(4.12) \quad Q_{\nu+1/2}(f) \leq \frac{2 + \lambda^2 (a^2)_{\nu+1/2}}{3}.)$$

Remark 4.5. In the case of the Lax-Friedrichs scheme, where $Q_{\nu+1/2}^{\text{LF}} = 1$ violates (4.11), we may take L to be zero, which leaves us with the original 3-point CFL limitation

$$(4.13) \quad \lambda |a_{\nu+1/2}| \leq 1.$$

Remark 4.6. Consider the special choice $(a^2)_{\nu+1/2} = (a_{\nu+1/2})^2$ in (4.3b), which is admissible in view of (2.4). The piecewise linear flux modification in Theorem 4.3 extends Harten's recipe [10], which was restricted to viscosity *functions* of the form $Q_{\nu+1/2} = Q(\lambda |a_{\nu+1/2}|)$; (hence, the Engquist-Osher and Godunov schemes, for example, were excluded from the discussion [10]).

We now turn to a second more attractive choice of SOR viscosity, inspired by Theorem 4.2 (see also [10, p. 368]). Here, we use a piecewise constant viscosity modification of the form

$$(4.14) \quad Q_{\nu+1/2}^{\text{SOR}} = Q_{\nu+1/2}(f) + \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right|,$$

which leads to

THEOREM 4.7. (*Piecewise constant viscosity modification.*) *The difference scheme (4.3), (4.14) is SOR-TVD under the CFL-like condition*

$$(4.15) \quad Q_{\nu+1/2}(f) \leq 2/3.$$

Proof. The SOR requirement in (4.4b) is directly verified by the very definition of $Q_{\nu+1/2}^{\text{SOR}}$ in (4.14); the TVD requirement on the left of (4.4a) follows from (4.1),

$$(4.16) \quad \left| \lambda a_{\nu+1/2} + \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \leq \lambda \left| \frac{f(v_{\nu+1}) - f(v_{\nu})}{v_{\nu+1} - v_{\nu}} \right| + \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \leq Q_{\nu+1/2}^{\text{SOR}},$$

and we are left with the inequality on the right, requiring

$$(4.17) \quad Q_{\nu+1/2}^{\text{SOR}} \leq 1.$$

Indeed, the same estimate used in Theorem 4.3 before, see (4.9b), together with the definition of $Q_{\nu+1/2}^{\text{SOR}}$, (4.14), gives us the pessimistic bound

$$(4.18) \quad Q_{\nu+1/2}^{\text{SOR}} \leq Q_{\nu+1/2}(f) + \frac{1}{2}[Q_{\nu+1/2}(f) - \lambda^2(a^2)_{\nu+1/2}] \leq \frac{3}{2}Q_{\nu+1/2}(f),$$

and (4.17) follows in view of the assumed CFL-like condition (4.15). $\square$

Remark 4.8. As before, by improving the crude bound of (4.18), the CFL-limit may be improved to be that of (4.12).

We note that the piecewise constant viscosity modification involves *linear* corrections of the numerical flux f_{ν} in (4.3d) and the viscosity coefficient in (4.14); both are appearing *linearly* in the original 3-point scheme we are starting with; see (2.2). *Hence, the resulting modification recipe does not change the underlying 3-point TVD code, but, rather, adds to it an antidiffusive term which boils down to*

$$(4.19) \quad -\frac{1}{2}(\tilde{g}_{\nu+1} - \tilde{g}_{\nu-1}) + \frac{1}{2} \left(\Delta \left(\left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \Delta v_{\nu+1/2} \right) \right).$$

To make our point more precise, we state as our final result of this first part, the following

COROLLARY 4.9. (*Piecewise constant viscosity modification—revisited.*) *Consider the 3-point scheme*

$$(4.20a) \quad v_{\nu}(t+k) = H(v_{\nu-1}, v_{\nu}, v_{\nu+1}; f, \lambda) \equiv v_{\nu} - \lambda(h_{\nu+1/2} - h_{\nu-1/2})$$

and assume the TVD-like constraint, compare (2.4),

$$(4.20b) \quad \lambda|a_{\nu+1/2}| \leq Q_{\nu+1/2}(f) \leq \frac{2}{3}, \quad Q_{\nu+1/2}(f) = \lambda \frac{f_{\nu} + f_{\nu+1} - 2h_{\nu+1/2}}{\Delta v_{\nu+1/2}},$$

holds. Then the modified scheme

$$(4.21a) \quad v_{\nu}(t+k) = H(v_{\nu-1}, v_{\nu}, v_{\nu+1}; f, \lambda) - \left[\left(\frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right)^- \Delta v_{\nu+1/2} + \left(\frac{\Delta \tilde{g}_{\nu-1/2}}{\Delta v_{\nu-1/2}} \right)^+ \Delta v_{\nu-1/2} \right]$$

is SOR-TVD. Here, $\tilde{g}_{\nu}$ is the flux correction, given by

$$(4.21b) \quad \tilde{g}_{\nu} = \frac{s_{\nu}}{2} \text{Min}([Q_{\nu\pm 1/2}(f) - \lambda^2(a^2)_{\nu\pm 1/2}] \cdot |\Delta v_{\nu\pm 1/2}|),$$

where

$$(4.21c) \quad \lambda^2(a^2)_{\nu+1/2} = \lambda^2(a_{\nu+1/2})^2 + O(|\Delta v|)_{\nu+1/2} \leq Q_{\nu+1/2}(f).$$

Alternatively, we can rewrite (4.21a) in the conservative form

$$(4.22a) \quad \begin{aligned} v_\nu(t+k) &= H(v_{\nu-1}, v_\nu, v_{\nu+1}; f, \lambda) - [\tilde{h}_{\nu+1/2} - \tilde{h}_{\nu-1/2}] \\ &= v_\nu(t) - \lambda \left[\left(h_{\nu+1/2} + \frac{1}{\lambda} \tilde{h}_{\nu+1/2} \right) - \left(h_{\nu-1/2} + \frac{1}{\lambda} \tilde{h}_{\nu-1/2} \right) \right]. \end{aligned}$$

Here $\tilde{h}_{\nu+1/2}$ is the numerical flux correction given by

$$(4.22b) \quad \tilde{h}_{\nu+1/2} = \frac{1}{2} [\tilde{g}_\nu + \tilde{g}_{\nu+1} - s_{\nu+1/2} |\tilde{g}_{\nu+1} - \tilde{g}_\nu|].$$

Remark 4.10. The Lax-Friedrichs scheme violates the CFL limitation (4.20b). Nonetheless, an “antidiffusive” modification follows in this case directly from the flux modification of Theorem 4.3, with numerical flux correction $\tilde{h}_{\nu+1/2}^{\text{LF}} = \frac{1}{2}(\tilde{g}_\nu + \tilde{g}_{\nu+1})$.

Remark 4.11. The conservative form of scheme (4.22) is identical with the corrective type of Harten’s artificial compression method [8], [9], the difference lying, of course, in the exact details of the flux correction used. In particular, we have an operator splitting which could be easily implemented as a two-step predictor-corrector method [8, Section 6].

PART II. ON A CELL ENTROPY INEQUALITY

5. Godunov’s Scheme and Its Cell Entropy Inequality. An entropy pair (U, F) associated with the conservative model (1.2) consists of a *convex entropy function*, $U(\cdot)$, augmented with an *entropy flux*, $F(\cdot)$, such that $U'f' = F'$. The requirement of having an *entropy inequality* of the form

$$(5.1) \quad \frac{\partial U}{\partial t}(u(x, t)) + \frac{\partial F}{\partial x}(u(x, t)) \leq 0$$

for *all* entropy pairs singles out the unique physically relevant (weak) solution of (1.2); e.g., [13], [14]. Accordingly, we are seeking conditions which guarantee that the difference scheme (1.1) will satisfy a *cell entropy inequality* of the form [12]

$$(5.2a) \quad U(v_\nu(t+k)) \leq U(v_\nu(t)) - \lambda(F_{\nu+1/2} - F_{\nu-1/2});$$

here, $F_{\nu+1/2} = F(v_{\nu-q+1}, \dots, v_{\nu+q}; f, \lambda)$ is a *numerical entropy flux*, assumed to be consistent with the differential one²

$$(5.2b) \quad F(w, w, \dots, w; f, \lambda) = F(w).$$

If this holds for a wide enough class of entropy functions, the cell entropy inequalities (5.2) will guarantee the *uniqueness* of limits of all *converging subsequences* of the total-variation bounded solutions, and there follows the convergence of the whole computed sequence to the unique physically relevant (weak) solution; e.g., [3], [12], [18], [23], [26].

As an example, we begin with the all-important cell entropy inequality associated with *Godunov’s scheme*: By averaging two Riemann solvers, one obtains Godunov’s scheme, whose numerical flux is given by [17], [18]

$$(5.3) \quad \begin{aligned} h_{\nu+1/2}^G &= s_{\nu+1/2} \min_{(v-v_\nu) \cdot (v-v_{\nu+1}) < 0} [s_{\nu+1/2} f(v)], \\ s_{\nu+1/2} &\equiv \text{sgn}(\Delta v_{\nu+1/2}). \end{aligned}$$

²Both the differential and the numerical entropy fluxes are denoted by F ; the distinction is made by the number of arguments.

Assuming that the minimum on the right-hand side is attained at a (not *necessarily unique*) value $v_{\nu+1/2}^G = v^G(v_\nu, v_{\nu+1}, f)$,

$$(5.4a) \quad s_{\nu+1/2} \underset{(v-v_\nu) \cdot (v-v_{\nu+1}) < 0}{\text{Min}} [s_{\nu+1/2} f(v)] = f(v_{\nu+1/2}^G),$$

then the Godunov scheme

$$(5.4b) \quad v(t+k) = H^G(v_{\nu-1}, v_\nu, v_{\nu+1}; f, \lambda)$$

boils down to

$$(5.4c) \quad H^G(v_{\nu-1}, v_\nu, v_{\nu+1}; f, \lambda) = v_\nu - \lambda[f(v_{\nu+1/2}^G) - f(v_{\nu-1/2}^G)].$$

Integrating over a typical cell the *differential* entropy inequality (5.1), Jensen's inequality yields the following well-known lemma; e.g., [3], [18], [26].

LEMMA 5.1. (*Cell entropy inequality of Godunov scheme.*) *Consider the Godunov scheme (5.4) satisfying the CFL condition*

$$(5.5a) \quad \lambda \underset{v}{\text{Max}} |f'(v)| \leq 1.^3$$

Then the following cell entropy inequality holds,

$$(5.5b) \quad U[H^G(v_{\nu-1}, v_\nu, v_{\nu+1}; f, \lambda)] \leq U(v_\nu) - \lambda[F(v_{\nu+1/2}^G) - F(v_{\nu-1/2}^G)],$$

for all entropy pairs (U, F) . Equivalently, (5.5b) can be rewritten as

$$(5.5c) \quad U[H^G(v_{\nu-1}, v_\nu, v_{\nu+1}; f, \lambda)] \leq U(v_\nu) - \lambda \int_{v_{\nu-1/2}^G}^{v_{\nu+1/2}^G} U'(w) f'(w) dw.$$

As in [26], the cell entropy inequality associated with the Godunov scheme will be used as a building block for studying such an inequality in conjunction with other TVD schemes. To this end, we shall make use of a modified flux function introduced below.

6. The Modified Flux Function. While studying total-variation diminishing schemes, only the modified flux *grid values*, g_ν , played a role; see the TVD requirement (1.8). Subsequently, a simple piecewise linear modified flux correction, $\tilde{g}_{\nu+1/2}(v)$, was introduced in (4.5a), in connection with the (first) recipe of SOR-TVD schemes.

Regarding the question of a cell entropy inequality which we now consider, a more sophisticated construction of a *modified flux function* is required: Let $g^{(\nu+1/2)}(v)$ denote its restriction to the $I_{\nu+1/2}$ -interval, $I_{\nu+1/2} = \{v \mid (v - v_\nu)(v - v_{\nu+1}) \leq 0\}$. The piecewise linear function $g^{(\nu+1/2)}(v)$ connects the modified flux grid values g_ν and $g_{\nu+1}$ on both ends of the interval, through a constant numerical flux value, $h_{\nu+1/2}$; it depends on two yet to be determined parameters (s^+, s^-) , $s^\pm > 0$, in the following fashion:

Setting the intermediate values

$$(6.1a) \quad v_{\nu+1/2}^+ \equiv v_{\nu+1/2}^+(s^+) = v_\nu + \frac{1}{s^+}(g_\nu - h_{\nu+1/2}),$$

$$(6.1b) \quad v_{\nu+1/2}^- \equiv v_{\nu+1/2}^-(s^-) = v_{\nu+1} - \frac{1}{s^-}(g_{\nu+1} - h_{\nu+1/2}),$$

³The maximum is taken over all values v varying between $v_{\nu-1}$, v_ν , and $v_{\nu+1}$.

we define, depending on whether $v_\nu \leq v_{\nu+1}$ or $v_{\nu+1} \leq v_\nu$,

$$(6.1c) \quad g^{(\nu+1/2)}(v; s^+, s^-) = \begin{bmatrix} g_\nu - s^+(v - v_\nu), & v_\nu \leq v \leq v_{\nu+1/2}^+ \\ h_{\nu+1/2}, & v_{\nu+1/2}^+ \leq v \leq v_{\nu+1/2}^- \\ g_{\nu+1} + s^-(v - v_{\nu+1}), & v_{\nu+1/2}^- \leq v \leq v_{\nu+1} \end{bmatrix},$$

$v_\nu \leq v_{\nu+1},$

$$(6.1d) \quad g^{(\nu+1/2)}(v; s^+, s^-) = \begin{bmatrix} g_{\nu+1} + s^-(v - v_{\nu+1}), & v_{\nu+1} \leq v \leq v_{\nu+1/2}^- \\ h_{\nu+1/2}, & v_{\nu+1/2}^- \leq v \leq v_{\nu+1/2}^+ \\ g_\nu - s^+(v - v_\nu), & v_{\nu+1/2}^+ \leq v \leq v_\nu \end{bmatrix},$$

$v_{\nu+1} \leq v_\nu.$

We have

LEMMA 6.1. *Assume the TVD condition, (1.8), holds:*

$$(6.2a) \quad \lambda \left| \frac{\Delta g_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \leq Q_{\nu+1/2} \leq 1.$$

Then $g^{(\nu+1/2)}(v; s^+, s^-)$, given in (6.1) and satisfying

$$(6.2b) \quad Q_{\nu+1/2} \leq \lambda s^\pm \leq 1,$$

is a well-defined piecewise linear function in the $I_{\nu+1/2}$ interval.

Verification. In view of (6.2a) we have, consult (1.10),

$$C_{\nu+1/2}^+ \geq 0, \quad C_{\nu+1/2}^- \geq 0, \quad Q_{\nu+1/2} = C_{\nu+1/2}^- + C_{\nu+1/2}^+ \leq 1,$$

while the definition of g_ν in (1.5b) yields

$$(6.3a) \quad v_{\nu+1/2}^+ = v_\nu + \frac{1}{s^+}(g_\nu - h_{\nu+1/2}) = v_\nu + \frac{1}{\lambda s^+} C_{\nu+1/2}^+ \Delta v_{\nu+1/2},$$

$$(6.3b) \quad v_{\nu+1/2}^- = v_{\nu+1} - \frac{1}{s^-}(g_{\nu+1} - h_{\nu+1/2}) = v_{\nu+1} - \frac{1}{\lambda s^-} C_{\nu+1/2}^- \Delta v_{\nu+1/2}.$$

Consider first the case where $v_\nu \leq v_{\nu+1}$: Then by (6.2b), $v_{\nu+1/2}^+$ is located to the left of $v_{\nu+1/2}^-$,

$$v_{\nu+1/2}^- - v_{\nu+1/2}^+ \geq \Delta v_{\nu+1/2} - \frac{1}{Q_{\nu+1/2}}(C_{\nu+1/2}^- + C_{\nu+1/2}^+) \Delta v_{\nu+1/2} = 0,$$

and to the right of v_ν ,

$$v_{\nu+1/2}^+ - v_\nu = \frac{1}{\lambda s^+} C_{\nu+1/2}^+ \Delta v_{\nu+1/2} \geq 0.$$

Taking into account the positivity of the incremental coefficients $C_{\nu+1/2}^\pm$, the following graph of $g^{(\nu+1/2)}(v; s^+, s^-)$ is obtained in this case; see Figure (6.4a). A completely analogous situation occurs when $v_{\nu+1} \leq v_\nu$, where the graph of $g^{(\nu+1/2)}(v)$ takes the form of Figure (6.4b).

Our next lemma is in the heart of the matter.

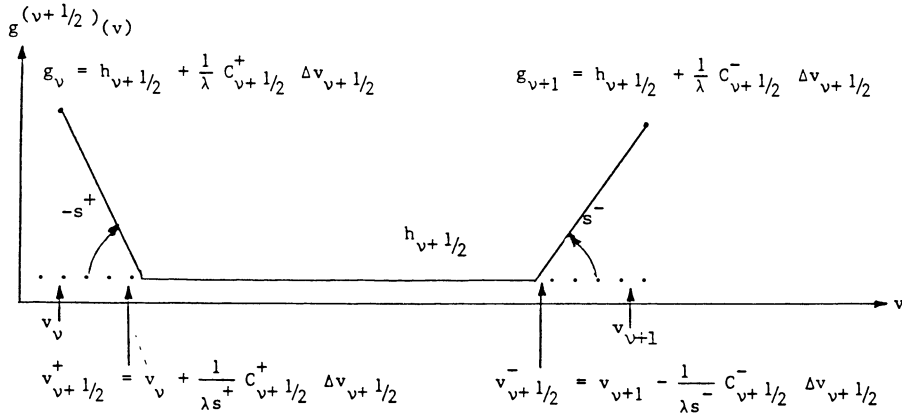


FIGURE 6.4a

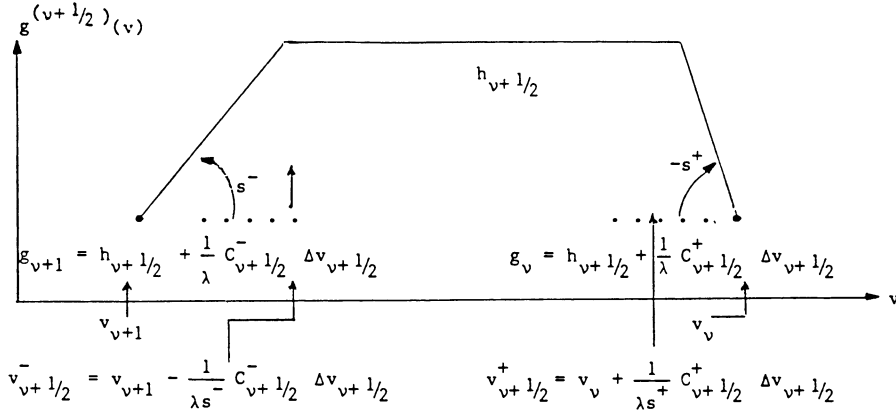


FIGURE 6.4b

LEMMA 6.2. Assume the TVD condition, compare (1.8),

$$(6.5)_{\pm} \quad \lambda \left| \frac{\Delta g_{\nu \pm 1/2}}{\Delta v_{\nu \pm 1/2}} \right| \leq Q_{\nu \pm 1/2} \leq \theta \leq 1$$

holds. Then we have

$$(6.6a) \quad \begin{aligned} H^G[v_{\nu}, v_{\nu}, v_{\nu+1}; g^{(\nu+1/2)}(s^+ = s^- = \theta \lambda^{-1}), \theta^{-1} \lambda] \\ = v_{\nu} - \theta^{-1} \lambda [h_{\nu+1/2} - g_{\nu}], \end{aligned}$$

$$(6.6b) \quad \begin{aligned} H^G[v_{\nu-1}, v_{\nu}, v_{\nu}; g^{(\nu-1/2)}(s^+ = s^- = \theta \lambda^{-1}), \theta^{-1} \lambda] \\ = v_{\nu} - \theta^{-1} \lambda [g_{\nu} - h_{\nu-1/2}]. \end{aligned}$$

Proof. We first note that according to Lemma 6.1, $g^{(\nu \pm 1/2)}(s^+ = s^- = \theta \lambda^{-1})$ are well defined: (6.5) $_{\pm}$ yields

$$\lambda \left| \frac{\Delta g_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \leq Q_{\nu+1/2} \leq \lambda s^{\pm} = \theta \leq 1.$$

By definition we have, see (5.4c),

$$(6.7a) \quad \begin{aligned} H^G[v_{\nu}, v_{\nu}, v_{\nu+1}; g^{(\nu+1/2)}(s^+ = s^- = \theta \lambda^{-1}), \theta^{-1} \lambda] \\ = v_{\nu} - \theta^{-1} \lambda [g^{(\nu+1/2)}(v_{\nu+1/2}^G) - g^{(\nu-1/2)}(v_{\nu})]; \end{aligned}$$

a straightforward computation gives, see (5.4a),

$$(6.7b) \quad g^{(\nu+1/2)}(v_{\nu+1/2}^G) = s_{\nu+1/2} \operatorname{Min}_{I_{\nu+1/2}} [s_{\nu+1/2} g^{(\nu+1/2)}(v)] = h_{\nu+1/2},$$

while consistency implies

$$(6.7c) \quad g^{(\nu+1/2)}(v_\nu) = g_\nu,$$

and (6.6a) follows. Similar arguments apply for (6.6b). $\square$

Remark 6.3. Godunov's numerical flux is determined by the differential one, computed at the intermediate value $v_{\nu+1/2}^G$; see (5.4a). In our case, (6.7b), *any* intermediate value between $v_{\nu+1/2}^+$ and $v_{\nu+1/2}^-$ can be taken: For later purposes, we shall choose in particular

$$(6.8) \quad \begin{aligned} v_{\nu+1/2}^G(\theta) &= \frac{1}{2} [v_{\nu+1/2}^+(s^+ = \theta\lambda^{-1}) + v_{\nu+1/2}^-(s^- = \theta\lambda^{-1})] \\ &= \frac{v_\nu + v_{\nu+1}}{2} - \frac{\lambda}{2\theta} (g_{\nu+1} - g_\nu). \end{aligned}$$

Taking $\theta = 1/2$ in Lemma 6.2 and averaging (6.6a) and (6.6b), we finally arrive at

THEOREM 6.4. *Consider the difference scheme (1.1) given in its viscous form (1.7a) and satisfying the TVD condition*

$$(6.9) \quad \lambda \left| \frac{\Delta g_{\nu\pm 1/2}}{\Delta v_{\nu\pm 1/2}} \right| \leq Q_{\nu\pm 1/2} \leq 1/2.$$

We then have

$$(6.10) \quad \begin{aligned} v_\nu(t+k) &= \frac{1}{2} \left[H^G \left[v_{\nu-1}, v_\nu, v_\nu; g^{(\nu-1/2)} \left(s^+ = s^- = \frac{1}{2\lambda} \right), 2\lambda \right] \right. \\ &\quad \left. + H^G \left[v_\nu, v_\nu, v_{\nu+1}; g^{(\nu+1/2)} \left(s^+ = s^- = \frac{1}{2\lambda} \right), 2\lambda \right] \right]. \end{aligned}$$

In other words, we have shown that any TVD satisfying difference scheme, (1.7a), (6.9), is given as an *average* of two Godunov solvers, thus refining a similar (convex) decomposition introduced in [26, Theorem 5.1]; the current decomposition differs, however, in its use of a modified flux function, rather than using the original differential one. In fact, away from the critical extremum values v_ν where $\Delta v_{\nu-1/2} \cdot \Delta v_{\nu+1/2} \leq 0$, we can do even better: Under appropriate TVD condition, any difference scheme coincides with *one* (rather than the average of two) Godunov solver. This is the content of our next

THEOREM 6.5. *Consider the difference scheme (1.1) given in its viscous form (1.7a) and satisfying the TVD condition*

$$(6.11) \quad \lambda \left| \frac{\Delta g_{\nu\pm 1/2}}{\Delta v_{\nu\pm 1/2}} \right| \leq Q_{\nu\pm 1/2} \leq 1.$$

Then, in the neighborhood of noncritical values v_ν such that $\Delta v_{\nu-1/2} \cdot \Delta v_{\nu+1/2} \geq 0$, we have

$$(6.12a) \quad v_\nu(t+k) = H^G \left[v_{\nu-1}, v_\nu, v_{\nu+1}; g \left(s^+ = s^- = \frac{1}{\lambda} \right), \lambda \right].$$

Remark 6.6. The modified flux function $g(v) = g(v; s^+, s^-)$ appearing in (6.12a) is *uniquely defined* in this case, composed of its restriction to the consecutive intervals $I_{\nu \pm 1/2}$; see Figures (6.13),

$$(6.12b) \quad g\left(v; s^+ = s^- = \frac{1}{\lambda}\right) = \begin{cases} g^{(\nu-1/2)}\left(v; s^+ = s^- = \frac{1}{\lambda}\right), & v \in I_{\nu-1/2}, \\ g^{(\nu+1/2)}\left(v; s^+ = s^- = \frac{1}{\lambda}\right), & v \in I_{\nu+1/2}. \end{cases}$$

Verification of Theorem 6.5 is straightforward: By (6.11) we may employ Lemma 6.2 with $\theta = 1$ in $I_{\nu \pm 1/2}$; using (5.4c) and (6.7b), the result (6.12) follows.

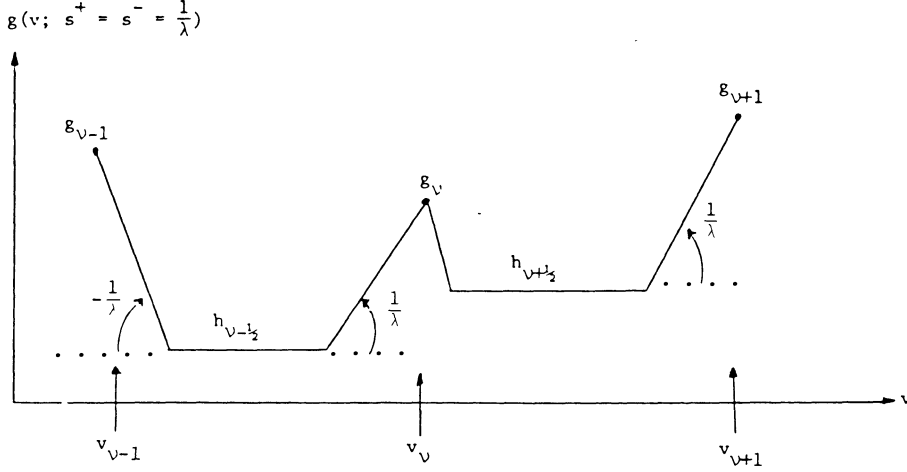


FIGURE 6.13a

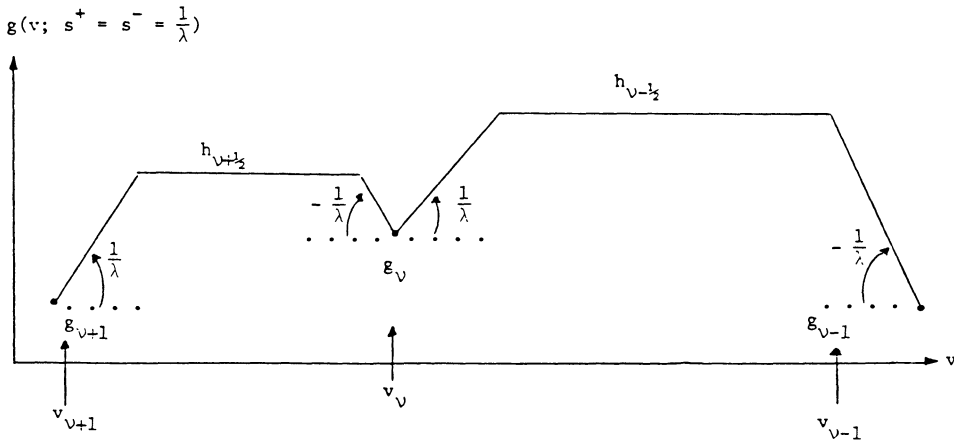


FIGURE 6.13b

We next use the cell entropy inequality associated with Godunov's scheme, see Lemma 5.1, to infer such an inequality for the difference schemes under considera-

tion. We start with

LEMMA 6.7. *Consider the TVD scheme (1.7a), (6.9). Then for all entropy pairs (U, F) we have*

$$(6.14a) \quad \begin{aligned} U[v_\nu(t+k)] &\leq U[v_\nu(t)] - \lambda(F_{\nu+1/2} - F_{\nu-1/2}) \\ &+ \lambda \int_{v_{\nu-1/2}^G}^{v_\nu} U''(v) \left[g^{(\nu-1/2)} \left(v; s^+ = s^- = \frac{1}{2\lambda} \right) - f(v) \right] dv \\ &+ \lambda \int_{v_\nu}^{v_{\nu+1/2}^G} U''(v) \left[g^{(\nu+1/2)} \left(v; s^+ = s^- = \frac{1}{2\lambda} \right) - f(v) \right] dv, \end{aligned}$$

with a consistent numerical entropy flux given by

$$(6.14b) \quad F_{\nu+1/2} = F(v_{\nu+1/2}^G) - U'(v_{\nu+1/2}^G) \cdot [h_{\nu+1/2} - f(v_{\nu+1/2}^G)].$$

Remark. Here, the intermediate Godunov value $v_{\nu+1/2}^G$ may be taken as, see (6.8),

$$(6.14c) \quad v_{\nu+1/2}^G = v_{\nu+1/2}^G(\theta = 1/2) = \frac{v_\nu + v_{\nu+1}}{2} - \lambda(g_{\nu+1} - g_\nu).$$

Proof. The maximal slope of $g^{(\nu+1/2)}(v; s^+ = s^- = 1/2\lambda)$ is $\frac{1}{2}\lambda$,

$$2\lambda \max_{v \in I_{\nu+1/2}} \left| g^{(\nu+1/2)} \left(v; s^+ = s^- = \frac{1}{2\lambda} \right) \right| \leq 1.$$

Therefore, we may apply Lemma 5.1, obtaining, see (5.5c),

$$(6.15a) \quad \begin{aligned} &U \left[H^G \left[v_{\nu-1}, v_\nu, v_\nu; g^{(\nu-1/2)} \left(v; s^+ = s^- = \frac{1}{2\lambda} \right), 2\lambda \right] \right] \\ &\leq U(v_\nu) - 2\lambda \int_{v_{\nu-1/2}^G}^{v_\nu} U'(v) g^{(\nu-1/2)'} \left(v; s^+ = s^- = \frac{1}{2\lambda} \right) dv, \end{aligned}$$

$$(6.15b) \quad \begin{aligned} &U \left[H^G \left[v_\nu, v_\nu, v_{\nu+1}; g^{(\nu+1/2)} \left(v; s^+ = s^- = \frac{1}{2\lambda} \right), 2\lambda \right] \right] \\ &\leq U(v_\nu) - 2\lambda \int_{v_\nu}^{v_{\nu+1/2}^G} U'(v) g^{(\nu+1/2)'} \left(v; s^+ = s^- = \frac{1}{2\lambda} \right) dv. \end{aligned}$$

Adding and subtracting $2\lambda[F(v_\nu) - F(v_{\nu-1/2}^G)] \equiv 2\lambda \int_{v_{\nu-1/2}^G}^{v_\nu} U' f'$ to the right-hand side of (6.15a), we find after integration by parts

$$(6.16a) \quad \begin{aligned} &U \left[H^G \left[v_{\nu-1}, v_\nu, v_\nu; g^{(\nu-1/2)} \left(v; s^+ = s^- = \frac{1}{2\lambda} \right), 2\lambda \right] \right] \\ &\leq U(v_\nu) - 2\lambda \int_{v_{\nu-1/2}^G}^{v_\nu} U' \left(g^{(\nu-1/2)} - f \right)' - 2\lambda [F(v_\nu) - F(v_{\nu+1/2}^G)] \\ &= U(v_\nu) - 2\lambda \cdot U'(v) \cdot \left(g^{(\nu-1/2)}(v) - f(v) \right) \Big|_{v=v_{\nu-1/2}^G}^{v=v_\nu} \\ &\quad - 2\lambda [F(v_\nu) - F(v_{\nu-1/2}^G)] + 2\lambda \int_{v_{\nu-1/2}^G}^{v_\nu} U'' [g^{(\nu-1/2)} - f]. \end{aligned}$$

In a similar manner, (6.15b) yields

$$\begin{aligned}
 (6.16b) \quad & U \left[H^G \left[v_\nu, v_\nu, v_{\nu+1}; g^{(\nu+1/2)} \left(v; s^+ = s^- = \frac{1}{2\lambda} \right), 2\lambda \right] \right] \\
 & \leq U(v_\nu) - 2\lambda \cdot U'(v) \cdot (g^{(\nu+1/2)}(v) - f(v)) \Big|_{v=v_\nu}^{v=v_{\nu+1/2}^G} \\
 & \quad - 2\lambda [F(v_{\nu+1/2}^G) - F(v_\nu)] + 2\lambda \int_{v_\nu}^{v_{\nu+1/2}^G} U''[g^{(\nu+1/2)} - f].
 \end{aligned}$$

By convexity, the average of the two terms on the left of (6.16a) and (6.16b) dominates the entropy value $U(\cdot)$, computed at the average of the corresponding two Godunov solvers; invoking Theorem 6.4, we arrive at

$$\begin{aligned}
 (6.17) \quad & U[v_\nu(t+k)] \leq U[v_\nu(t)] \\
 & - \lambda \Delta \left[F(v_{\nu-1/2}^G) + U'(v_{\nu-1/2}^G) \cdot \left[g^{(\nu-1/2)}(v_{\nu-1/2}^G) - f(v_{\nu-1/2}^G) \right] \right] \\
 & + \lambda \int_{v_{\nu-1/2}^G}^{v_\nu} U''(g^{(\nu-1/2)} - f) + \lambda \int_{v_\nu}^{v_{\nu+1/2}^G} U''(g^{(\nu+1/2)} - f),
 \end{aligned}$$

and the result (6.14) follows, noting that $g^{(\nu \pm 1/2)}(v_{\nu \pm 1/2}^G) = h_{\nu \pm 1/2}$; see (6.7b). $\square$

In the same way as the last lemma followed from Theorem 6.4, Theorem 6.5 gives us

LEMMA 6.8. *Consider the TVD scheme (1.7a), (6.11) in the neighborhood of a noncritical value v_ν where $\Delta v_{\nu-1/2} \cdot \Delta v_{\nu+1/2} \geq 0$. Then for all entropy pairs (U, F) we have*

$$\begin{aligned}
 (6.18) \quad & U[v_\nu(t+k)] \leq U[v_\nu(t)] - \lambda(F_{\nu+1/2} - F_{\nu-1/2}) \\
 & + \lambda \int_{v_{\nu-1/2}^G}^{v_\nu} U''(v) \left[g^{(\nu-1/2)} \left(v; s^+ = s^- = \frac{1}{\lambda} \right) - f(v) \right] dv \\
 & + \lambda \int_{v_\nu}^{v_{\nu+1/2}^G} U''(v) \left[g^{(\nu+1/2)} \left(v; s^+ = s^- = \frac{1}{\lambda} \right) - f(v) \right] dv.
 \end{aligned}$$

Remark 6.9. The numerical entropy flux, $F_{\nu+1/2}$, in (6.18) is determined by Godunov's intermediate value $v_{\nu+1/2}^G$. It coincides with the one given in (6.14b) and (6.14c): Indeed, $v_{\nu+1/2}^G = v_{\nu+1/2}^G(\theta = 1/2)$ is located in between $v_{\nu+1/2}^+(s^+ = \lambda^{-1})$ and $v_{\nu+1/2}^-(s^- = \lambda^{-1})$. The proof is omitted.

The only difference in Lemma 6.8 is in the slopes, $s^\pm$, involved in the last two integrals on the right of (6.18): At a noncritical value, they are twice than otherwise allowed by Lemma 6.7; see (6.14a). Making use of $|s_\nu|$ to distinguish between these two cases, see (4.3a), the corresponding two estimates can be unified as follows.

THEOREM 6.10. *Consider the difference scheme (1.7a) satisfying the TVD condition*

$$(6.19) \quad \lambda \left| \frac{\Delta g_{\nu \pm 1/2}}{\Delta v_{\nu \pm 1/2}} \right| \leq Q_{\nu \pm 1/2} \leq \frac{1 + |s_\nu|}{2}.$$

Then, for all entropy pairs (U, F) we have

$$\begin{aligned}
 (6.20a) \quad & U[v_\nu(t+k)] \\
 & \leq U[v_\nu(t)] - \lambda(F_{\nu+1/2} - F_{\nu-1/2}) \\
 & + \lambda \int_{v_{\nu-1/2}^G}^{v_\nu} U''(v) \left[g^{(\nu-1/2)} \left(v; s^+ = s^- = \frac{1+|s_\nu|}{2\lambda} \right) - f(v) \right] dv \\
 & + \lambda \int_{v_\nu}^{v_{\nu+1/2}^G} U''(v) \left[g^{(\nu+1/2)} \left(v; s^+ = s^- = \frac{1+|s_\nu|}{2\lambda} \right) - f(v) \right] dv,
 \end{aligned}$$

with a consistent numerical entropy flux

$$(6.20b) \quad F_{\nu+1/2} = F(v_{\nu+1/2}^G) - U'(v_{\nu+1/2}^G)[h_{\nu+1/2} - f(v_{\nu+1/2}^G)].$$

For later purposes, we shall prefer a slightly different version of inequality (6.20), where the last two integrals on the right are *shifted* over the $I_{\nu+1/2}$ interval. We state

THEOREM 6.11. *Consider the difference scheme (1.7a) satisfying the TVD condition*

$$(6.21) \quad \lambda \left| \frac{\Delta g_{\nu\pm 1/2}}{\Delta v_{\nu\pm 1/2}} \right| \leq Q_{\nu\pm 1/2} \leq \frac{1+|s_\nu|}{2}.$$

Then for all entropy pairs (U, F) we have

$$\begin{aligned}
 (6.22a) \quad & U[v_\nu(t+k)] \leq U[v_\nu(t)] - \lambda(F_{\nu+1/2} - F_{\nu-1/2}) \\
 & + \lambda \int_{v_\nu}^{v_{\nu+1}} U''(v) \left[g^{(\nu+1/2)} \left(v; s_\nu^+ = \frac{1+|s_\nu|}{2\lambda}, \right. \right. \\
 & \quad \left. \left. s_{\nu+1}^- = \frac{1+|s_{\nu+1}|}{2\lambda} \right) - f(v) \right] dv,
 \end{aligned}$$

with a consistent numerical entropy flux

$$\begin{aligned}
 (6.22b) \quad & F_{\nu+1/2} = F(v_{\nu+1/2}^G) - U'(v_{\nu+1/2}^G)[h_{\nu+1/2} - f(v_{\nu+1/2}^G)] \\
 & + \lambda \int_{v_{\nu+1/2}^G}^{v_{\nu+1}} U''(v) \left[g^{(\nu+1/2)} \left(v; s^+ = s^- = \frac{1+|s_{\nu+1}|}{2\lambda} \right) - f(v) \right] dv.
 \end{aligned}$$

Remark. Theorem 6.11 finally leads us to the *modified flux function*⁴ we were looking for; its restriction to the $I_{\nu+1/2}$ -interval is given, according to (6.22a), by

$$\begin{aligned}
 (6.23) \quad & g(v) |_{I_{\nu+1/2}} = g^{(\nu+1/2)}(v) \\
 & = g^{(\nu+1/2)} \left(v; s_\nu^+ = \frac{1+|s_\nu|}{2\lambda}, s_{\nu+1}^- = \frac{1+|s_{\nu+1}|}{2\lambda} \right).
 \end{aligned}$$

Proof. We first note the identity

$$\begin{aligned}
 (6.24) \quad & \int_{v_\nu}^{v_{\nu+1/2}^G} U''(v) g^{(\nu+1/2)}(v; s_1^+, s_1^-) dv + \int_{v_{\nu+1/2}^G}^{v_{\nu+1}} U''(v) g^{(\nu+1/2)}(v; s_2^+, s_2^-) dv \\
 & = \int_{v_\nu}^{v_{\nu+1}} U''(v) g^{(\nu+1/2)}(v; s_1^+, s_2^-) dv.
 \end{aligned}$$

⁴To be precise, $g(v)$ is double-valued in the critical neighborhood.

Indeed, of the two integrals on the left, the first depends only on s_1^+ , the second only on s_2^- . Adding and subtracting

$$\lambda \int_{v_{\nu+1/2}^G}^{v_{\nu+1}} U''(v) \left[g^{(\nu+1/2)} \left(v; s^+ = s^- = \frac{1 + |s_{\nu+1}|}{2\lambda} \right) - f(v) \right] dv$$

to the right-hand side of (6.20a) and using (6.24), the theorem follows. $\square$

Remark 6.12. The two integrals appearing on the right of (6.22) are responsible for the unnaturally shifted estimates with respect to the v_ν -grid value. If, in the previous proof, we instead add and subtract

$$\lambda \int_{v_{\nu-1}}^{v_{\nu-1/2}^G} U''(v) \left[g^{(\nu-1/2)} \left(v; s^+ = s^- = \frac{1 + |s_{\nu-1}|}{2} \right) - f(v) \right] dv$$

to the right-hand side of (6.20a) and use (6.24), we end up with a reversed shifted form,

$$(6.25a) \quad \begin{aligned} U[v_\nu(t+k)] &\leq U[v_\nu(t)] - \lambda(F_{\nu+1/2} - F_{\nu-1/2}) \\ &\quad + \lambda \int_{v_{\nu-1}}^{v_\nu} U''(v) \left[g^{(\nu-1/2)} \left(v; s_{\nu-1}^+ = \frac{1 + |s_{\nu-1}|}{2\lambda}, \right. \right. \\ &\quad \left. \left. s_\nu^- = \frac{1 + |s_\nu|}{2\lambda} \right) - f(v) \right] dv, \end{aligned}$$

where

$$(6.25b) \quad \begin{aligned} F_{\nu+1/2} &= F(v_{\nu+1/2}^G) - U'(v_{\nu+1/2}^G) [h_{\nu+1/2} - f(v_{\nu+1/2}^G)] \\ &\quad - \lambda \int_{v_\nu}^{v_{\nu+1/2}^G} U''(v) \left[g^{(\nu+1/2)} \left(v; s^+ = s^- = \frac{1 + |s_\nu|}{2\lambda} \right) - f(v) \right] dv. \end{aligned}$$

Averaging of (6.22) and (6.25) results in the symmetric entropy estimate of this type,

$$(6.26) \quad \begin{aligned} U[v_\nu(t+k)] &\leq U[v_\nu(t)] - \lambda(F_{\nu+1/2} - F_{\nu-1/2}) \\ &\quad + \lambda \int_{v_{\nu-1}}^{v_\nu} U''(v) [g(v) |_{I_{\nu-1/2}} - f(v)] dv \\ &\quad + \lambda \int_{v_\nu}^{v_{\nu+1}} U''(v) [g(v) |_{I_{\nu+1/2}} - f(v)] dv. \end{aligned}$$

7. Necessity and Sufficiency of E -Fluxes. Theorem 6.11 shows that a desired cell entropy inequality, (5.2a), follows, provided the last integral on the right of (6.22a) is negative,

$$(7.1) \quad \int_{v_\nu}^{v_{\nu+1}} U''(v) \left[g^{(\nu+1/2)} \left(v; s^+ = \frac{1 + |s_\nu|}{2\lambda}, s^- = \frac{1 + |s_{\nu+1}|}{2\lambda} \right) - f(v) \right] dv \leq 0.$$

The requirement of negativity for *all* convex entropy functions U , in this case, is equivalent to the requirement that $g(v) |_{I_{\nu+1/2}}$ will lie below or above $f(v)$, depending on whether $v_\nu \leq v_{\nu+1}$ or $v_{\nu+1} \leq v_\nu$; taking into account the piecewise linear form of $g(v) |_{I_{\nu+1/2}}$ in (6.23), see Figure 6.4, this requirement amounts to the

following inequalities:

$$\begin{aligned}
 (i) \quad & s_{\nu+1/2} \cdot [h_{\nu+1/2} - f(v)] \leq 0, \quad v \in I_{\nu+1/2},^5 \\
 (ii) \quad (a) \quad & \lambda \operatorname{Max}_{|v-v_\nu| \leq \varepsilon_\nu^+} \left| \frac{f(v) - g_\nu}{v - v_\nu} \right| \leq \frac{1 + |s_\nu|}{2}, \\
 & \varepsilon_\nu^+ = \frac{2}{1 + |s_\nu|} C_{\nu+1/2}^+ |\Delta v_{\nu+1/2}|, \\
 (b) \quad & \lambda \operatorname{Max}_{|v-v_{\nu+1}| \leq \varepsilon_{\nu+1}^-} \left| \frac{f(v) - g_{\nu+1}}{v - v_{\nu+1}} \right| \leq \frac{1 + |s_{\nu+1}|}{2}, \\
 & \varepsilon_{\nu+1}^- = \frac{2}{1 + |s_{\nu+1}|} C_{\nu+1/2}^- |\Delta v_{\nu+1/2}|.
 \end{aligned}$$

Numerical fluxes satisfying condition (i) are called *E-fluxes* after Osher [18]. Equivalently, the corresponding schemes are characterized as exactly those having more numerical viscosity than that of Godunov's scheme [26]. In terms of the modified viscosity coefficient $\mathbf{Q}_{\nu+1/2}$ in (3.3), the first inequality (i) reads, consult (2.5c),

$$(E) \quad Q_{\nu+1/2}^G \leq \mathbf{Q}_{\nu+1/2}.$$

Observe that two viscosity coefficients are comparable in the sense advocated in [26], provided that the *same* modified flux is used for the corresponding schemes; in the above E-condition, the original flux, f_ν , was used in both cases, see Remark 3.1. Regarding the CFL condition in the second inequality (ii), it can be shifted to the more symmetric form in the spirit of Remark 6.12,

$$(CFL) \quad \lambda \operatorname{Max}_{|v-v_\nu| \leq \varepsilon_\nu^\pm} \left| \frac{f(v) - g_\nu}{v - v_\nu} \right| \leq \frac{1 + |s_\nu|}{2}, \quad \varepsilon_\nu^\pm = \frac{2C_{\nu\pm 1/2}^\pm |\Delta v_{\nu+1/2}|}{1 + |s_\nu|}.$$

With a similar shifting of the TVD condition (6.21), we arrive at

THEOREM 7.1 (TADMOR [26, THEOREM 5.1]). *Consider an E-type TVD scheme (1.7a) such that*

$$(7.2) \quad Q_{\nu+1/2} \leq \operatorname{Min} \left(\frac{1 + |s_\nu|}{2}, \frac{1 + |s_{\nu+1}|}{2} \right),$$

and assume the following CFL condition holds:

$$\begin{aligned}
 (7.3) \quad & \lambda \operatorname{Max}_{|v-v_\nu| \leq \varepsilon_\nu^\pm} \left| \frac{f(v) - g_\nu}{v - v_\nu} \right| \leq \frac{1 + |s_\nu|}{2}, \\
 & \varepsilon_\nu^\pm = \frac{2C_{\nu\pm 1/2}^\pm |\Delta v_{\nu\pm 1/2}|}{1 + |s_\nu|}.
 \end{aligned}$$

Then, for all entropy pairs (U, F) the following cell entropy inequality is satisfied:

$$(7.4) \quad U[v_\nu(t+k)] \leq U[v_\nu(t)] - \lambda(F_{\nu+1/2} - F_{\nu-1/2}).$$

⁵ It is sufficient to consider v lying in between

$$v_{\nu+1/2}^+(s^+ = (1 + |s_\nu|)/2\lambda) \quad \text{and} \quad v_{\nu+1/2}^-(s^- = (1 + |s_{\nu+1}|)/2\lambda).$$

Remark. We have preferred to center the TVD condition (7.2) and CFL condition (7.3) around the v_ν -grid value. Therefore, no specific reference was given as to the consistent numerical entropy flux $F_{\nu+1/2}$ in (7.4): It may be chosen, *throughout the computational grid*, by (6.22b), (6.25b) or any convex combination of the two.

Proof. The only part which requires clarification is the validity of the TVD constraint, see (6.21),

$$(7.5) \quad \lambda \left| \frac{\Delta g_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \leq Q_{\nu+1/2}.$$

Indeed, (3.1) and (2.5) yield

$$(7.6) \quad \lambda \left| \frac{\Delta g_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \leq \lambda |a_{\nu+1/2}| + \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \leq Q_{\nu+1/2}^G + \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right|;$$

since $\tilde{g}_\nu/\Delta v_{\nu\pm 1/2}$ is positive by (3.7c), we find, as in (4.9b),

$$(7.7) \quad \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \leq \frac{\text{Max}(|\tilde{g}_\nu|, |\tilde{g}_{\nu+1}|)}{|\Delta v_{\nu+1/2}|} \leq \frac{\tilde{g}_\nu + \tilde{g}_{\nu+1}}{\Delta v_{\nu+1/2}}.$$

Augmented with the E-condition, (7.6) and (7.7) yield (7.5):

$$\lambda \left| \frac{\Delta g_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \leq Q_{\nu+1/2}^G + \frac{\tilde{g}_\nu + \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \leq Q_{\nu+1/2} + \frac{\tilde{g}_\nu + \tilde{g}_{\nu+1}}{\Delta v_{\nu+1/2}} \equiv Q_{\nu+1/2}. \quad \square$$

Remark 7.2. It is the CFL-like condition (7.3) which throws some further light on the first-order accuracy limitation encountered with difference schemes satisfying *all* cell-wise entropy inequalities (5.2a). Indeed, in case $\Delta v_{\nu+1/2} > 0$ (similarly $\Delta v_{\nu-1/2} < 0$), then necessarily $g_\nu = f_\nu + \tilde{g}_\nu/\lambda$ will lie *above* (similarly *below*) f_ν , in contrast to the requirement induced by (7.1); specifically, as v approaches v_ν , the CFL-like ratio on the left of (7.3),

$$(7.8) \quad \lambda \text{Max}_{|v-v_\nu| \leq \varepsilon_\nu^\pm} \left| \frac{f(v) - f(v_\nu)}{v - v_\nu} + \frac{1}{\lambda} \frac{\tilde{g}_\nu}{v - v_\nu} \right|,$$

will diverge unless the correction $\tilde{g}_\nu$ vanishes. In this case, the modified flux g_ν is reduced to the differential one, f_ν , and we are back in the standard first-order accurate E-schemes, subject to the familiar

$$(7.9) \quad \lambda \text{Max}_{|v-v_\nu| \leq \varepsilon_\nu^\pm} \left| \frac{f(v) - f(v_\nu)}{v - v_\nu} \right| \leq \frac{1}{2}.$$

In particular, according to (7.9), *the usual half CFL number is sufficient in this case*, in agreement with [26]:

$$(7.10) \quad \lambda \text{Max}_{|v-v_\nu| \leq |\Delta v_{\nu\pm 1/2}|} |f'(v)| \leq \frac{1}{2}.$$

Of course, (7.9) offers a more delicate alternative. Note that, unlike the TVD constraint (1.8), some further information regarding the flux behavior *in between the grid values* is always required for a cell entropy inequality to hold.

Next we show that the E-condition is, in a sense described below, *necessary* for a cell entropy inequality, thus complementing the sufficiency of Theorem 7.1. To

this end, we subtract $v_\nu(t)$ from both sides of (1.1) and (1.3), divide by $k = \Delta t$ and take the limit $\Delta t \downarrow 0$, obtaining the *semidiscrete approximation*

$$(7.11a) \quad \frac{dv_\nu}{dt} + \frac{1}{\Delta x}(h_{\nu+1/2} - h_{\nu-1/2}) = 0;$$

here, the time derivative is understood in the distribution sense, and the numerical flux is still allowed to depend on the mesh size, Δx ,

$$(7.11b) \quad h_{\nu+1/2} = h(v_{\nu-p+1}, \dots, v_{\nu+p}; f, \lambda \downarrow 0, \Delta x).$$

The corresponding cell entropy inequality will read

$$(7.12a) \quad \frac{dU}{dt}[v_\nu(t)] + \frac{1}{\Delta x}(F_{\nu+1/2} - F_{\nu-1/2}) \leq 0;$$

here, (U, F) is any consistent *numerical* entropy pair

$$(7.12b) \quad F(w, w, \dots, w; f, \lambda \downarrow 0, \Delta x) = F(w) \equiv \int^w U' f'.$$

In stating our next theorem, the terminology of *separated data* will be used: We shall say that the data are separated if all relevant grid values involved in the scheme's stencil are separated with respect to, say, the *separating interval* $I_{j+1/2}$; that is, we either have

$$(7.13a) \quad v_{j-k} \leq v_j < v_{j+1} \leq v_{j+k}, \quad k = 1, \dots, p,$$

or

$$(7.13b) \quad v_{j-k} \geq v_j > v_{j+1} \geq v_{j+k}, \quad k = 1, \dots, p.$$

A monotone profile is, of course, the canonical example of separated data: Any two consecutive grid values may serve as a separating interval in this case.

Equipped with this terminology, we now turn to

THEOREM 7.3. (*E-flux is necessary for cell entropy inequality.*) *Consider the semidiscrete approximation (7.11) and assume the cell entropy inequality (7.12) holds for all consistent numerical entropy pairs (U, F) .*

Then the numerical flux, h , is an E-flux in any separating interval.

Proof. Introducing the consistent numerical entropy flux

$$(7.14a) \quad \mathbf{F}_{\nu+1/2} = F(v_{\nu+1}) - U'(v_{\nu+1})[h_{\nu+1/2} - f(v_{\nu+1})],$$

the following *equality* [18, Section 3] holds,

$$(7.14b) \quad \frac{\partial U}{\partial t}[v_\nu(t)] + \frac{1}{\Delta x}[\mathbf{F}_{\nu+1/2} - \mathbf{F}_{\nu-1/2}] = \frac{1}{\Delta x} \int_{v_\nu}^{v_{\nu+1}} U''(w)[h_{\nu+1/2} - f(w)] dw.$$

Subtracting (7.14b) from (7.12a) we find, after multiplication by Δx ,

$$(7.15)_\nu \quad \int_{v_\nu}^{v_{\nu+1}} U''(w)[h_{\nu+1/2} - f(w)] dw + (H_{\nu+1/2} - H_{\nu-1/2}) \leq 0;$$

here, $H_{\nu+1/2}$ stands for the difference $F_{\nu+1/2} - \mathbf{F}_{\nu+1/2}$, which is consistent with zero,

$$(7.16a) \quad H(w, w, \dots, w; f, \lambda \downarrow 0, \Delta x) = F(w) - \mathbf{F}(w) = 0.$$

Consider a $(2p + 1)$ -grid-valued stencil separated by, say, the $I_{j+1/2}$ interval; we shall extend it to *all* grid points, defining

$$(7.16b) \quad v_{j+k} \equiv v_{j+p}, \quad v_{j-k} \equiv v_{j-p}, \quad k = p + 1, p + 2, \dots$$

Let v be *any* value in between v_j and v_{j+1} and apply $(7.15)_\nu$ with Kružkov's choice for entropy function, $U(w) = |w - v|$: For all but the $\nu = j$ index, the first term on the left of $(7.15)_\nu$ vanishes, and we end up with

$$(7.17a) \quad H_{j+1/2} \geq H_{j+3/2} \geq \dots \geq H_\infty = 0, \quad \nu = j + 1, j + 2, \dots,$$

$$(7.17b) \quad H_{j-1/2} \leq H_{j-3/2} \leq \dots \leq H_{-\infty} = 0, \quad \nu = j - 1, j - 2, \dots$$

Here, the equalities $H_{\pm\infty} = 0$ on the right of (7.17) follow from the consistency (7.16a) and the definition of our extension in (7.16b). Complemented with $(7.15)_j$, which in this case reads

$$(7.18) \quad \begin{aligned} 0 &\geq \int_{v_j}^{v_{j+1}} \delta(w - v) [h_{j+1/2} - f(w)] dw + (H_{j+1/2} - H_{j-1/2}) \\ &= s_{j+1/2} [h_{j+1/2} - f(v)] + (H_{j+1/2} - H_{j-1/2}), \end{aligned}$$

we find, for all v in $I_{j+1/2}$,

$$s_{j+1/2} [h_{j+1/2} - f(v)] \leq H_{j-1/2} - H_{j+1/2} \leq 0.$$

That is, h is an E-flux over the $I_{j+1/2}$ interval as asserted. $\square$

Next we have to confront the limitation of E-schemes being at most first-order accurate [18]: Indeed, according to the E-condition, their *modified* numerical viscosity is *bounded away* from that of the Lax-Wendroff scheme; hence (3.4) fails. Moreover, being an E-flux in separating intervals alone still meets the limitation of first-order accuracy. The proof follows along the lines of [18].

8. Second-Order Resolution and a Cell Entropy Inequality. In light of our discussion in the last section, we are seeking cell entropy inequalities for SOR-TVD schemes only for *special* entropy pairs: Specifically, we shall consider the pair $U(v) = v^2/2$, $F(v) = \int^v w f'(w) dw$. In the *genuinely nonlinear* case where f is, say, strictly convex, all (convex) entropy inequalities (and therefore convergence to the unique solution) then follow *in the limit*; see [14], [5].

We start with 3-point E-type *upwind* schemes. Upwinding simply means that we difference in the “streamwise” direction. In particular, considering the characteristic direction, we require either $C_{\nu+1/2}^+$ or $C_{\nu+1/2}^-$ to vanish, depending on whether $f'(v)$ is positive or negative *throughout the* $I_{\nu+1/2}$ -interval; in either of these two nonsonic cases, the sum of the incremental coefficients $C_{\nu+1/2}^\pm$ equals their difference in absolute value, i.e.,

$$(8.1a) \quad Q_{\nu+1/2}(f) = \lambda \left| \frac{\Delta f_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right|, \quad f'(v) |_{I_{\nu+1/2}} \neq 0.$$

This is, of course, nothing but a restatement of the well-known fact that away from sonic values, all upwind schemes coincide with Roe-Murman's (2.5d).

The SOR-TVD schemes to be considered are then constructed according to the piecewise constant viscosity modification of Theorem 4.7, where the viscosity coefficient used is given by

$$(8.1b) \quad Q_{\nu+1/2} \equiv Q_{\nu+1/2}^{\text{SOR}} = Q_{\nu+1/2}(f) + \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right|.$$

Here, $\tilde{g}$ is the flux correction, see (4.21b),

$$(8.1c) \quad \tilde{g}_\nu = \frac{s_\nu}{2} \text{Min}([Q_{\nu\pm 1/2}(f) - \lambda^2(a^2)_{\nu\pm 1/2}] \cdot |\Delta v_{\nu\pm 1/2}|)$$

with second-order resolution condition requiring

$$(8.1d) \quad \lambda^2(a^2)_{\nu+1/2} \equiv \lambda^2(a_{\nu+1/2})^2 + O(|\Delta v|)_{\nu+1/2} \leq Q_{\nu+1/2}(f),$$

and a first-order perturbation $O(|\Delta v|)_{\nu+1/2}$, which is yet to be precisely determined. Unless otherwise stated, the following CFL-like condition will be assumed,

$$(8.1e) \quad Q_{\nu+1/2}(f) \leq 1/3.$$

In particular, the above use of the piecewise viscosity modification from Theorem 4.7 is justified, since (4.15) holds. Furthermore, the essential estimate (4.9b),

$$\left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \leq \frac{1}{2} [Q_{\nu+1/2}(f) - \lambda^2(a^2)_{\nu+1/2}],$$

yields the TVD requirement (6.21),

$$\begin{aligned} \lambda \left| \frac{\Delta g_{\nu\pm 1/2}}{\Delta v_{\nu\pm 1/2}} \right| &\leq \lambda |a_{\nu\pm 1/2}| + \left| \frac{\Delta \tilde{g}_{\nu\pm 1/2}}{\Delta v_{\nu\pm 1/2}} \right| \\ &\leq Q_{\nu\pm 1/2}(f) + \left| \frac{\Delta \tilde{g}_{\nu\pm 1/2}}{\Delta v_{\nu\pm 1/2}} \right| \equiv Q_{\nu\pm 1/2}^{\text{SOR}} \leq \frac{3}{2} Q_{\nu\pm 1/2}(f) \leq \frac{1}{2}, \end{aligned}$$

so that Theorem 6.11 applies.⁶ According to this result, the desired cell entropy inequality holds for the above quadratic entropy, provided

$$(8.2) \quad \int_{v_\nu}^{v_{\nu+1}} \left[g^{(\nu+1/2)} \left(v; s_\nu^+ = \frac{1+|s_\nu|}{2\lambda}, s_{\nu+1}^- = \frac{1+|s_{\nu+1}|}{2\lambda} \right) - f(v) \right] dv$$

is nonpositive.

We first study the critical case, stating

LEMMA 8.1. *Consider the critical case $|s_\nu| = |s_{\nu+1}| = 0$ (that is, both v_ν and $v_{\nu+1}$ are critical grid values), and assume the following CFL condition holds*

$$(8.3) \quad \lambda \max_{|v-v_\nu| \leq |\Delta v_{\nu\pm 1/2}|} |f'(v)| \leq 1/2.$$

We then have

$$(8.4) \quad \int_{v_\nu}^{v_{\nu+1}} \left[g^{(\nu+1/2)} \left(v; s_\nu^+ = \frac{1+|s_\nu|}{2\lambda}, s_{\nu+1}^- = \frac{1+|s_{\nu+1}|}{2\lambda} \right) - f(v) \right] dv \leq 0.$$

Proof. We appeal to the *conservative* form of the modified scheme, quoted in (4.22): By assumption, s_ν and therefore $\tilde{g}_\nu$ vanish, and hence

$$\tilde{h}_{\nu+1/2} = \frac{1}{2} [\tilde{g}_{\nu+1} - s_{\nu+1/2} |\tilde{g}_{\nu+1}|];$$

since $s_{\nu+1/2}$ agrees with the sign of $\tilde{g}_{\nu+1}$, we conclude that the numerical flux correction $\tilde{h}_{\nu+1/2}$ vanishes in this case. Thus, *we are left with the original E-flux we have started with.* By Theorem 7.1 therefore, all entropy weighted integrals in (7.1), and in particular (8.4), are nonpositive, provided the CFL condition (7.3)

⁶In fact, twice the CFL limit (8.1e) can be used away from critical neighborhoods where $|s_\nu| \cdot |s_{\nu+1}| \neq 0$.

holds: The latter follows in view of (8.3), noting that the modified flux corrections $\tilde{g}_\nu$ and $\tilde{g}_{\nu+1}$ vanish—consult Remark 7.2. $\square$

Next we turn to the noncritical case $|s_\nu| + |s_{\nu+1}| \geq 1$: Calculating the area below the piecewise linear “double-wing” form of

$$g^{(\nu+1/2)} \left(v; s_\nu^+ = \frac{1 + |s_\nu|}{2\lambda}, s_{\nu+1}^- = \frac{1 + |s_{\nu+1}|}{2\lambda} \right),$$

see Figure 6.4, we obtain

$$\begin{aligned} (8.5a) \quad & \int_{v_\nu}^{v_{\nu+1}} \left[g^{(\nu+1/2)} \left(v; s_\nu^+ = \frac{1 + |s_\nu|}{2\lambda}, s_{\nu+1}^- = \frac{1 + |s_{\nu+1}|}{2\lambda} \right) - f(v) \right] dv \\ &= \Delta v_{\nu+1/2} h_{\nu+1/2} + \frac{1}{2\lambda} \left[\frac{2}{1 + |s_\nu|} (C_{\nu+1/2}^+)^2 \right. \\ & \quad \left. + \frac{2}{1 + |s_{\nu+1}|} (C_{\nu+1/2}^-)^2 \right] (\Delta v_{\nu+1/2})^2 \\ & \quad - \int_{v_\nu}^{v_{\nu+1}} f(v) dv; \end{aligned}$$

inserting the value of $h_{\nu+1/2}$ from (1.7b) and (3.1),

$$h_{\nu+1/2} = \frac{f_\nu + f_{\nu+1}}{2} + \frac{1}{2\lambda} (\tilde{g}_\nu + \tilde{g}_{\nu+1}) - \frac{1}{2\lambda} Q_{\nu+1/2}^{\text{SOR}} \Delta v_{\nu+1/2},$$

we find

$$\begin{aligned} (8.5b) \quad & \int_{v_\nu}^{v_{\nu+1}} \left[g^{(\nu+1/2)} \left(v; s_\nu^+ = \frac{1 + |s_\nu|}{2\lambda}, s_{\nu+1}^- = \frac{1 + |s_{\nu+1}|}{2\lambda} \right) - f(v) \right] dv \\ &= \Delta v_{\nu+1/2} \frac{f_\nu + f_{\nu+1}}{2} - \int_{v_\nu}^{v_{\nu+1}} f(v) dv \\ & \quad + \frac{(\Delta v_{\nu+1/2})^2}{2\lambda} \left[\frac{\tilde{g}_\nu + \tilde{g}_{\nu+1}}{\Delta v_{\nu+1/2}} - Q_{\nu+1/2}^{\text{SOR}} + \frac{2}{1 + |s_\nu|} (C_{\nu+1/2}^+)^2 \right. \\ & \quad \left. + \frac{2}{1 + |s_{\nu+1}|} (C_{\nu+1/2}^-)^2 \right]. \end{aligned}$$

In the next proposition we estimate from above the last brackets on the right of (8.5b). The somewhat technical proof is postponed to the end of this section.

PROPOSITION 8.2. *The following estimate holds:*

$$\begin{aligned} (8.6) \quad & \frac{\tilde{g}_\nu + \tilde{g}_{\nu+1}}{\Delta v_{\nu+1/2}} - Q_{\nu+1/2}^{\text{SOR}} + \frac{2}{1 + |s_\nu|} (C_{\nu+1/2}^+)^2 + \frac{2}{1 + |s_{\nu+1}|} (C_{\nu+1/2}^-)^2 \\ & \leq \frac{|s_\nu| + |s_{\nu+1}|}{2} [Q_{\nu+1/2}^2(f) - \lambda^2 (a^2)_{\nu+1/2}]. \end{aligned}$$

In view of (8.5b) and Proposition 8.2, we can now complement Lemma 8.1 with

LEMMA 8.3. *Consider the noncritical case $|s_\nu| + |s_{\nu+1}| \geq 1$ (that is, either v_ν or $v_{\nu+1}$ is a noncritical value). We then have*

$$\begin{aligned} (8.7a) \quad & \int_{v_\nu}^{v_{\nu+1}} \left[g^{(\nu+1/2)} \left(v; s_\nu^+ = \frac{1 + |s_\nu|}{2\lambda}, s_{\nu+1}^- = \frac{1 + |s_{\nu+1}|}{2\lambda} \right) - f(v) \right] dv \\ & \leq \frac{(\Delta v_{\nu+1/2})^2}{2\lambda} \left[\frac{2\lambda}{(\Delta v_{\nu+1/2})^2} \text{ET}_{\nu+1/2} \right. \\ & \quad \left. + \frac{|s_\nu| + |s_{\nu+1}|}{2} (Q_{\nu+1/2}^2(f) - \lambda^2 (a^2)_{\nu+1/2}) \right]. \end{aligned}$$

Here, $\text{ET}_{\nu+1/2}$ stands for the error in the trapezoidal rule, applied over the $I_{\nu+1/2}$ -interval,

$$(8.7b) \quad \text{ET}_{\nu+1/2} = \Delta v_{\nu+1/2} \left[\frac{f(v_\nu) + f(v_{\nu+1})}{2} \right] - \int_{v_\nu}^{v_{\nu+1}} f(v) dv.$$

In order to guarantee a cell entropy inequality in the noncritical case, the integral on the left of (8.7a) should be nonpositive. This can be achieved by choosing $\lambda^2(a^2)_{\nu+1/2}$ so as to make the expression in brackets on the right of (8.7a) nonpositive.

We claim that such a choice admissible by the second-order resolution condition (8.1d) is indeed available in all nonsonic intervals: According to the upwinding property (8.1a),

$$(8.8a) \quad Q_{\nu+1/2}^2(f) = \lambda^2 |a_{\nu+1/2}|^2,$$

and by a standard error estimate [4, Section 2],

$$(8.8b) \quad \text{ET}_{\nu+1/2} = \frac{(\Delta v_{\nu+1/2})^3}{12} f''(\zeta), \quad \zeta \in I_{\nu+1/2}.$$

It is here that we use the first-order flexibility in defining $\lambda^2(a^2)_{\nu+1/2}$ previously noticed in Theorem 4.2. We choose

$$\lambda^2(a^2)_{\nu+1/2} = Q_{\nu+1/2}^2(f) + \lambda \frac{|\Delta v_{\nu+1/2}|}{3} \text{Max}_{I_{\nu+1/2}} |f''(v)|.$$

Unifying the critical and noncritical cases, we arrive at

$$(8.9) \quad \lambda^2(a^2)_{\nu+1/2} = Q_{\nu+1/2}^2(f) + \frac{\lambda}{3} |\Delta v_{\nu+1/2}| \text{Max}_{I_{\nu+1/2}} |f''(v)|.$$

In order to comply with the inequality on the right of (8.1d), we set in (8.1c)

$$(8.10a) \quad \tilde{g}_\nu = \frac{s_\nu}{2} [\text{Min}[Q_{\nu\pm 1/2}(f) - \lambda^2(a^2)_{\nu\pm 1/2}]^+ \cdot |\Delta v_{\nu\pm 1/2}|].$$

In terms of the numerical flux modification

$$(8.10b) \quad \tilde{g}_{\nu+1/2} = \frac{1}{2} [\tilde{g}_\nu + \tilde{g}_{\nu+1} - s_{\nu+1/2} |\tilde{g}_{\nu+1} - \tilde{g}_\nu|],$$

the resulting scheme then reads:

$$(8.10c) \quad v_\nu(t+k) = v_\nu(t) - \lambda(f_{\nu+1} - f_{\nu-1}) + \frac{1}{2} [\Delta(Q_{\nu+1/2}(f)\Delta v_{\nu+1/2}(t))] - (\tilde{h}_{\nu+1/2} - \tilde{h}_{\nu-1/2}),$$

and we summarize what we have shown in

THEOREM 8.4. *Let $Q_{\nu+1/2}(f)$ be an upwind E-type viscosity coefficient, and assume the following CFL condition holds:*

$$(8.11) \quad \lambda \cdot \text{Max}|f'(v)| \leq 1/3.$$

Then the difference scheme (8.9), (8.10) satisfies

- (i) *total-variation diminishing;*
- (ii) *second-order resolution;*

(iii) *the cell entropy inequality*⁷

$$\frac{1}{2}v_\nu^2(t+k) \leq \frac{1}{2}v_\nu^2(t) - \lambda(F_{\nu+1/2} - F_{\nu+1/2}).$$

Remark 8.5. As argued in Theorem 4.2, the scheme has formal second-order accuracy in all intervals such that $\tilde{g}_\nu$ and $\tilde{g}_{\nu+1}$ do not vanish. According to (8.10a), therefore, all nonsonic noncritical intervals such that

$$\begin{aligned} Q_{\nu+1/2}(f) - \lambda^2(a^2)_{\nu+1/2} &= \lambda|a_{\nu+1/2}| - \lambda^2|a_{\nu+1/2}|^2 \\ &\quad - \lambda \frac{|\Delta v_{\nu+1/2}|}{3} \text{Max}_{I_{\nu+1/2}} |f''(v)| \geq 0 \end{aligned}$$

are included; in view of the CFL condition (8.11), the first difference on the right exceeds $2\lambda|a_{\nu+1/2}|/3$, and the last inequality is valid provided that

$$(8.12) \quad \frac{1}{2}|\Delta v_{\nu+1/2}| \cdot \text{Max}_{I_{\nu+1/2}} |f''(v)| \leq |a_{\nu+1/2}|.$$

In other words, *scheme (8.9), (8.10) is second-order accurate at all noncritical neighborhoods satisfying (8.12)*. Observe that increasing $\lambda^2(a^2)_{\nu+1/2}$ results in a similar increase of the *modified* viscosity coefficient $Q_{\nu+1/2}$ in (3.4a).

Remark 8.6. The additional viscosity added is directly related to the amount of $u^2/2$ entropy loss across shock discontinuities which is precisely $|\text{ET}_{\nu+1/2}|$. This should indicate the possible generalizations to other entropy pairs; moreover, it seems to imply that second-order resolution is the maximum possible for cell entropy satisfying schemes, the entropy loss being cubic [14].

As a special case, let us consider now the *genuinely nonlinear model* where f is, say, strictly convex. We distinguish between two possibilities.

(i) *The shock case:* $\Delta v_{\nu+1/2} \leq 0$. Since the trapezoidal error in (8.8b) is negative in this case, the admissible choice $\lambda^2(a^2)_{\nu+1/2} = \lambda^2|a_{\nu+1/2}|^2 = Q_{\nu+1/2}^2(f)$ will make the integral (8.7a) nonpositive, as required.

(ii) *The rarefaction case:* $\Delta v_{\nu+1/2} > 0$. As expected, this is the more delicate case because the sign of the trapezoidal error is now reversed to be positive, which will be compensated by additional dissipation as before. To summarize, we choose

$$(8.13) \quad \lambda^2(a^2)_{\nu+1/2} = Q_{\nu+1/2}^2(f) + \frac{\lambda}{3}(\Delta v_{\nu+1/2})^+ \text{Max}_{I_{\nu+1/2}} [f''(v)],$$

in which case we have

THEOREM 8.5. (*The convergence of SOR-TVD schemes.*) *Assume f is strictly convex, $Q_{\nu+1/2}(f)$ is an upwind E-type viscosity and the CFL condition (8.11) holds. Then the difference scheme (8.13), (8.10) satisfies*

- (i) *total-variation diminishing;*
- (ii) *second-order resolution, where*

$$\frac{1}{2}(\Delta v_{\nu+1/2})^+ \text{Max}_{I_{\nu+1/2}} [f''(v)] \leq |a_{\nu+1/2}|;$$

- (iii) *a consistent quadratic cell entropy inequality;*
- and, as a consequence of (i) and (iii),*
- (iv) *convergence to the unique physically relevant solution.*

⁷The numerical entropy flux is given in (6.22b).

Remark 8.6. The above-mentioned quadratic cell entropy inequality is due to the related inequality of Theorem 6.11. If (the nonshifted form of) Theorem 6.10 is used instead, the same cell entropy inequality follows, this time with added dissipation compensating the *rectangular rule* error rather than the trapezoidal one. Details are omitted. Recently, Osher [19] has considered SOR-TVD *semidiscrete* approximations to the genuinely nonlinear model (1.2). Convergence was then guaranteed with more dissipation added at *shocks*, rather than at the expected rarefactions. Indeed, the estimates used in [19] correspond to those quoted in our Theorem 6.10, and the dissipation unnaturally, yet necessarily, added at shocks can be therefore attributed to the *reversed sign* of the rectangular rule error. We also note that in the cases studied in [19], [20] and in the last two theorems, formal second-order accuracy had to be given up at strong jumps, which were measured with respect to the amount of convexity, $\text{Max}_{I_{\nu+1/2}}[f''(v)]$.

Remark 8.7. One can obtain similar results by starting with 3-point TVD schemes which are not necessarily of upwind type. For example, choosing the modified Lax-Friedrichs schemes where $Q_{\nu+1/2}(f) \equiv 1/2$, see [26], one can repeat the above arguments and conclude similar convergence results. Details are omitted.

We close this section with the

Proof of Proposition 8.2. The flux correction $\tilde{g}_\nu$ in (8.1c) satisfies

$$(8.14)_\nu \quad \frac{\tilde{g}_\nu}{\Delta v_{\nu\pm 1/2}} \leq \frac{|s_\nu|}{2} [Q_{\nu\pm 1/2}(f) - \lambda^2(a^2)_{\nu\pm 1/2}];$$

summing (8.14) $_\nu$ and (8.14) $_{\nu+1}$ we find

$$(8.15) \quad \frac{\tilde{g}_\nu + \tilde{g}_{\nu+1}}{\Delta v_{\nu+1/2}} \leq \frac{|s_\nu| + |s_{\nu+1}|}{2} [Q_{\nu+1/2}(f) - \lambda^2(a^2)_{\nu+1/2}].$$

We recall the definition of the viscosity coefficient in (8.16),

$$(8.16) \quad -Q_{\nu+1/2}^{\text{SOR}} = -Q_{\nu+1/2}(f) - \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right|.$$

Finally, we examine the incremental coefficients $C_{\nu+1/2}^\pm$ in (1.9): Utilizing (3.1), we have

$$(8.17a) \quad C_{\nu+1/2}^\pm = \frac{1}{2} \left[Q_{\nu+1/2}(f) + \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \mp \lambda a_{\nu+1/2} \mp \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right].$$

Consequently, the following equality holds,

$$(8.17b) \quad \frac{1 - |s_\nu|}{1 + |s_\nu|} (C_{\nu+1/2}^+)^2 = \frac{1 - |s_\nu|}{4} [Q_{\nu+1/2}(f) - \lambda a_{\nu+1/2}]^2.$$

Indeed, if $|s_\nu| = 1$, then both sides vanish; if otherwise s_ν and therefore $\tilde{g}_\nu$ vanish, then

$$\left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| - \left(\frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right) = \left| \frac{\tilde{g}_{\nu+1}}{\Delta v_{\nu+1/2}} \right| - \frac{\tilde{g}_{\nu+1}}{\Delta v_{\nu+1/2}}$$

vanishes as well, so that in view of (8.17a) equality holds in (8.17b). Similar arguments yield

$$(8.17c) \quad \frac{1 - |s_{\nu+1}|}{1 + |s_{\nu+1}|} (C_{\nu+1/2}^-)^2 = \frac{1 - |s_{\nu+1}|}{4} [Q_{\nu+1/2}(f) + \lambda a_{\nu+1/2}]^2.$$

In view of the TVD constraint (1.8) we also have

$$(8.17d) \quad \begin{aligned} (C_{\nu+1/2}^+)^2 + (C_{\nu+1/2}^-)^2 &= \frac{1}{2} \left[(Q_{\nu+1/2}^{\text{SOR}})^2 + \lambda^2 \left(\frac{\Delta g_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right)^2 \right] \\ &\leq \left[Q_{\nu+1/2}(f) + \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \right]^2. \end{aligned}$$

Adding (8.17b), (8.17c), and (8.17d) we arrive at

$$(8.18) \quad \begin{aligned} \frac{2}{1+|s_\nu|} (C_{\nu+1/2}^+)^2 + \frac{2}{1+|s_{\nu+1}|} (C_{\nu+1/2}^-)^2 \\ \leq \left[Q_{\nu+1/2}(f) + \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \right]^2 + (2 - |s_\nu| - |s_{\nu+1}|) Q_{\nu+1/2}^2(f). \end{aligned}$$

The inequalities (8.15), (8.16) together with (8.18) amount to the desired inequality (8.6): If $|s_\nu| = |s_{\nu+1}| = 1$, the derived upper bound reads

$$Q_{\nu+1/2}^2(f) - \lambda^2 (a^2)_{\nu+1/2} + \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \cdot \left[-1 + 2Q_{\nu+1/2}(f) + \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \right],$$

where in view of (8.14)_ν, together with the CFL condition (8.1e), the last brackets on the right are indeed nonpositive:

$$-1 + 2Q_{\nu+1/2}(f) + \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \leq 1 + \frac{5}{2} Q_{\nu+1/2}(f) < 0;$$

if $|s_\nu| + |s_{\nu+1}| = 1$, we find

$$\begin{aligned} -\frac{1}{2} Q_{\nu+1/2}(f) - \frac{1}{2} \lambda^2 (a^2)_{\nu+1/2} + 2Q_{\nu+1/2}^2(f) \\ + \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \cdot \left[-1 + 2Q_{\nu+1/2}(f) + \left| \frac{\Delta \tilde{g}_{\nu+1/2}}{\Delta v_{\nu+1/2}} \right| \right]. \end{aligned}$$

As before, the last brackets on the right are nonpositive, while the sum of the preceding three terms does not exceed the asserted value,

$$\begin{aligned} -\frac{1}{2} Q_{\nu+1/2}(f) + \frac{3}{2} Q_{\nu+1/2}^2(f) + \frac{1}{2} \left[Q_{\nu+1/2}^2(f) - \lambda^2 (a^2)_{\nu+1/2} \right] \\ \leq \frac{1}{2} [Q_{\nu+1/2}^2(f) - \lambda^2 (a^2)_{\nu+1/2}]; \end{aligned}$$

finally, although not specifically referred to, $\tilde{g}_\nu = \tilde{g}_{\nu+1} = 0$ in the critical case, $|s_\nu| = |s_{\nu+1}| = 0$, yielding

$$-Q_{\nu+1/2}(f) + 3Q_{\nu+1/2}^2(f) \leq 0. \quad \square$$

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